

Explanation of the ATLAS Z-peaked excess in the NMSSM

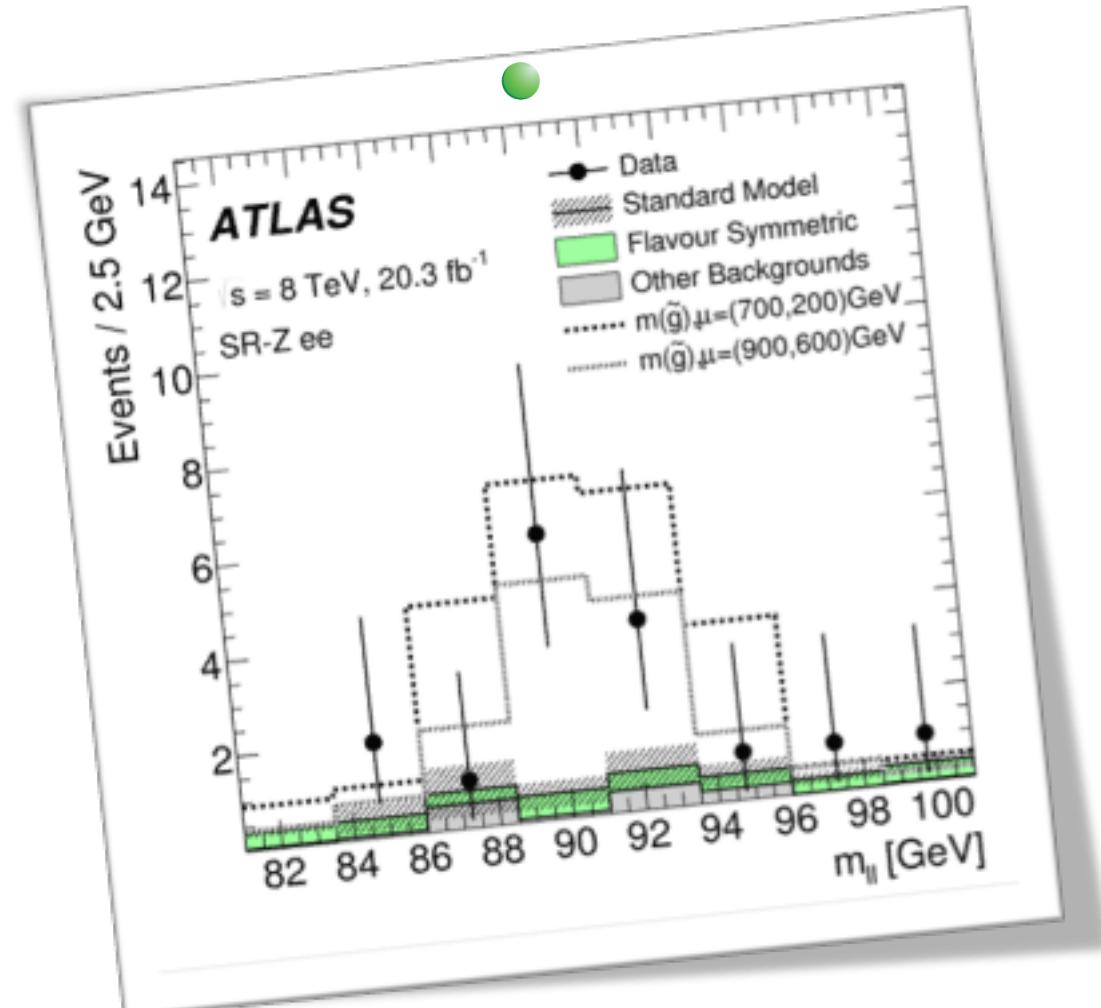
Yang Zhang, ITP, CAS

Beijing University of Technology

2015.05.18

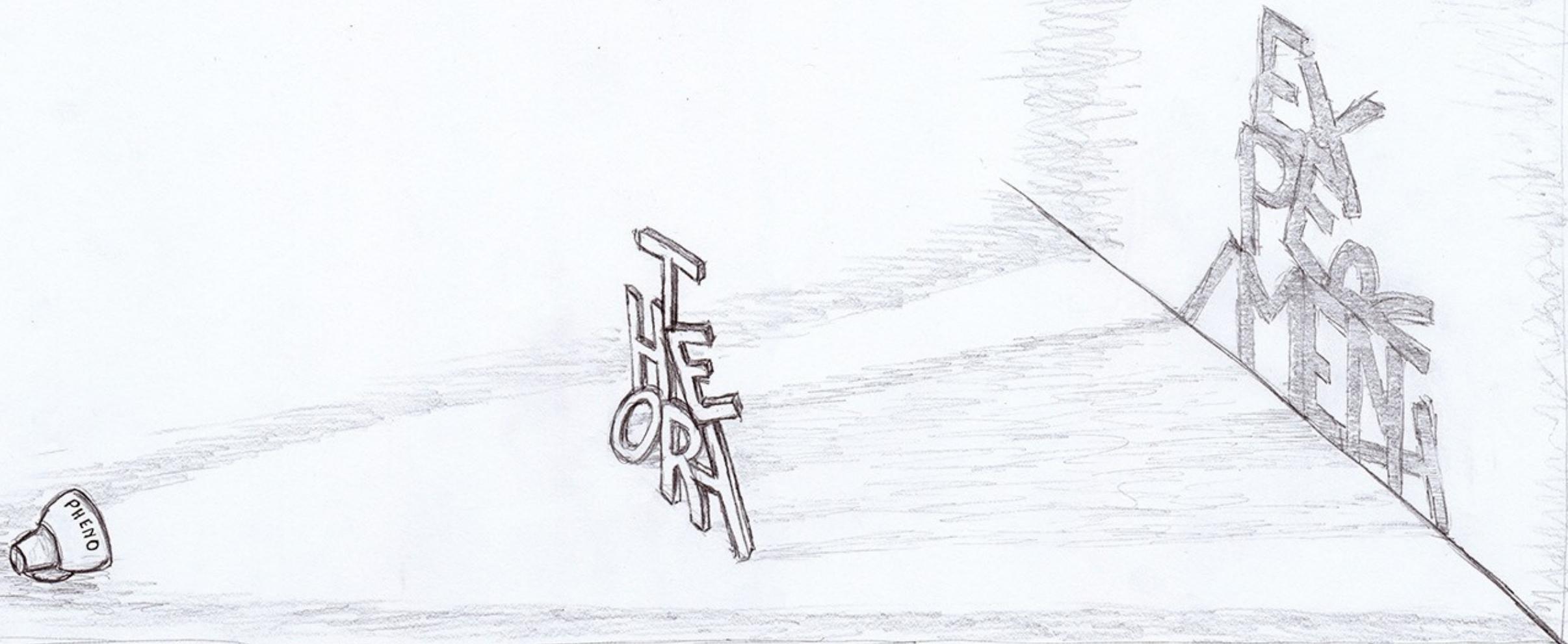
Cooperate with

Junjie Cao, Jinmin Yang,
Liangliang Shang



- I Monte Carlo simulation
 - 1. Why do we need to do MC simulation?
 - 2. How to do MC simulation in collider physics?
 - 3. Which kind of work can MC simulation do?
- II Example: Explanation of the ATLAS Z-peaked excess in the NMSSM
 - 1. Brief introduction of SUSY
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I Monte Carlo simulation



1.Why do we need to do MC simulation?

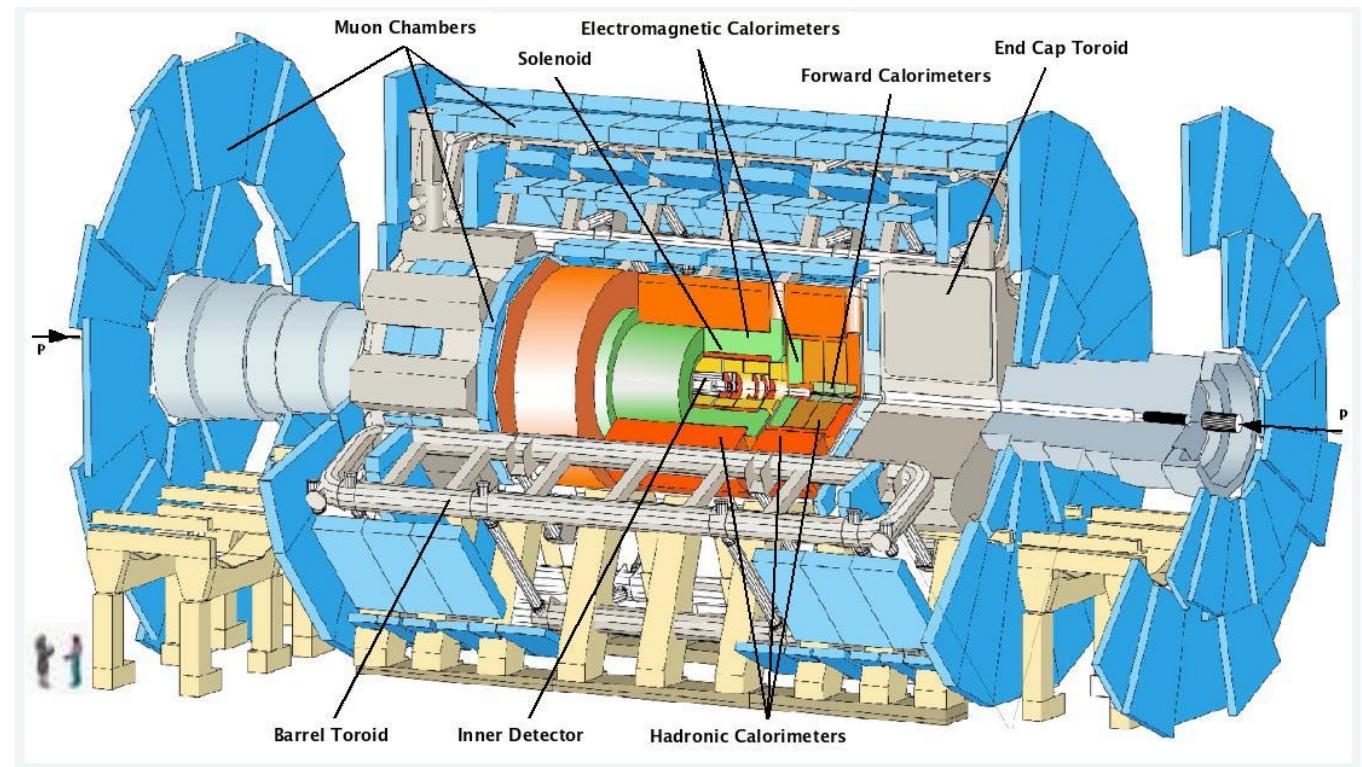
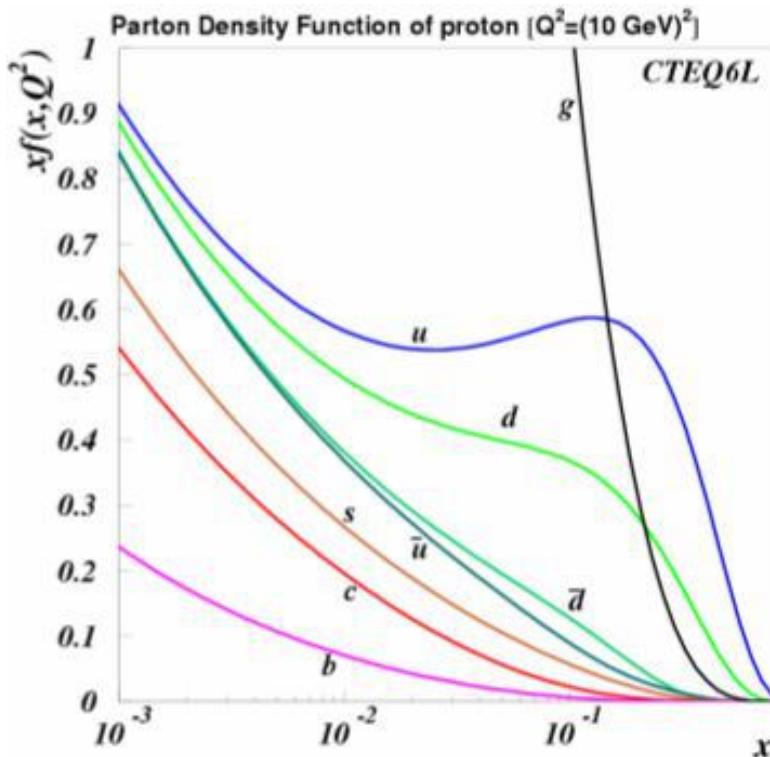
Particle physics **phenomenology** is the part of theoretical particle physics that deals with the application of theoretical physics to high-energy particle physics experiments.

Some examples [\[edit\]](#)

- Monte Carlo simulation studies of physics processes at colliders.
- Next-to-leading order calculations of particle production rates and distributions.
- Extraction of parton distribution functions from data.
- Application of heavy quark effective field theory to extract CKM matrix elements.
- Using lattice QCD to extract quark masses and CKM matrix elements from experiment.
- "Phenomenological analyses," in which one studies the experimental consequences of adding the most general set of beyond Standard Model effects in a given sector of the Standard Model, usually parameterized in terms of anomalous couplings and higher operators. In this case, the term "phenomenological" is being used more in its philosophy of science sense.

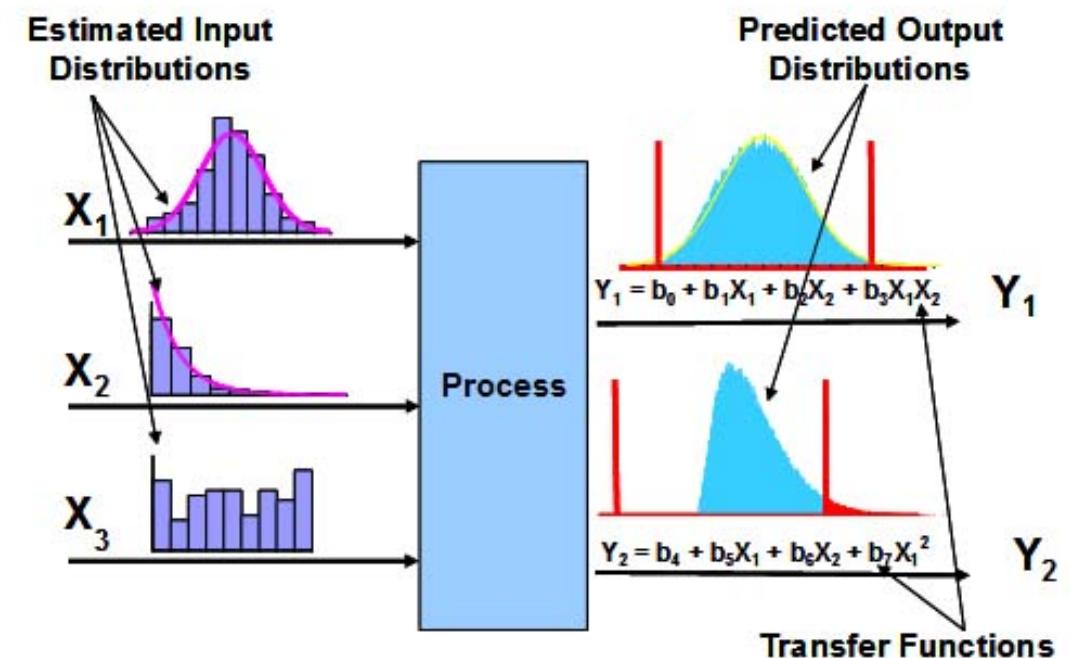
1. Why do we need to do MC simulation?

- Theory $\longrightarrow \dot{N} = \sigma L \longrightarrow$ Experiment

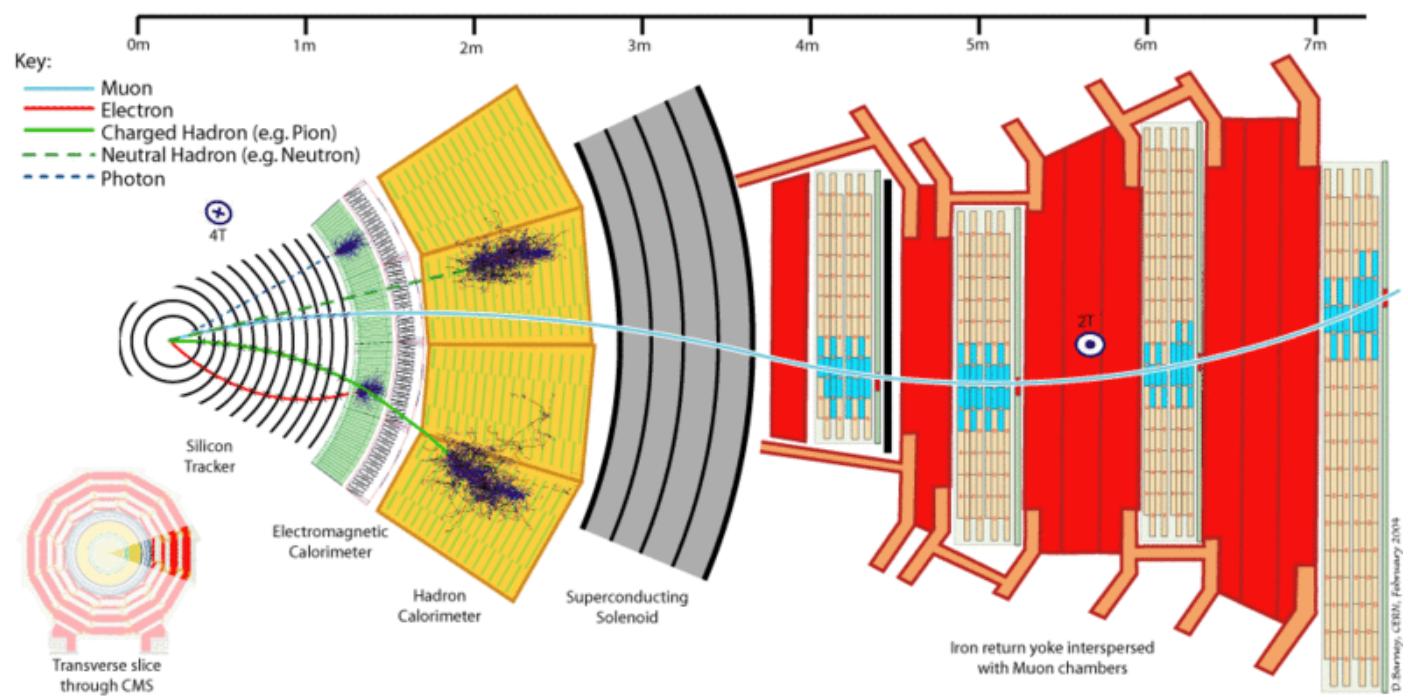
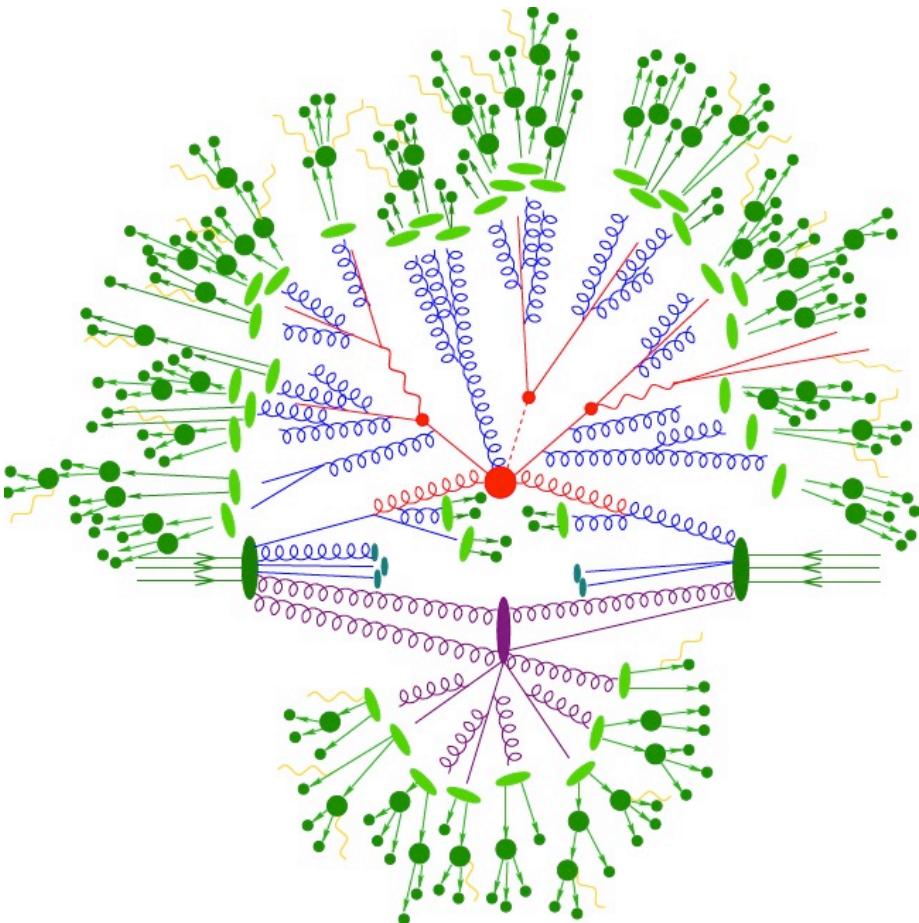


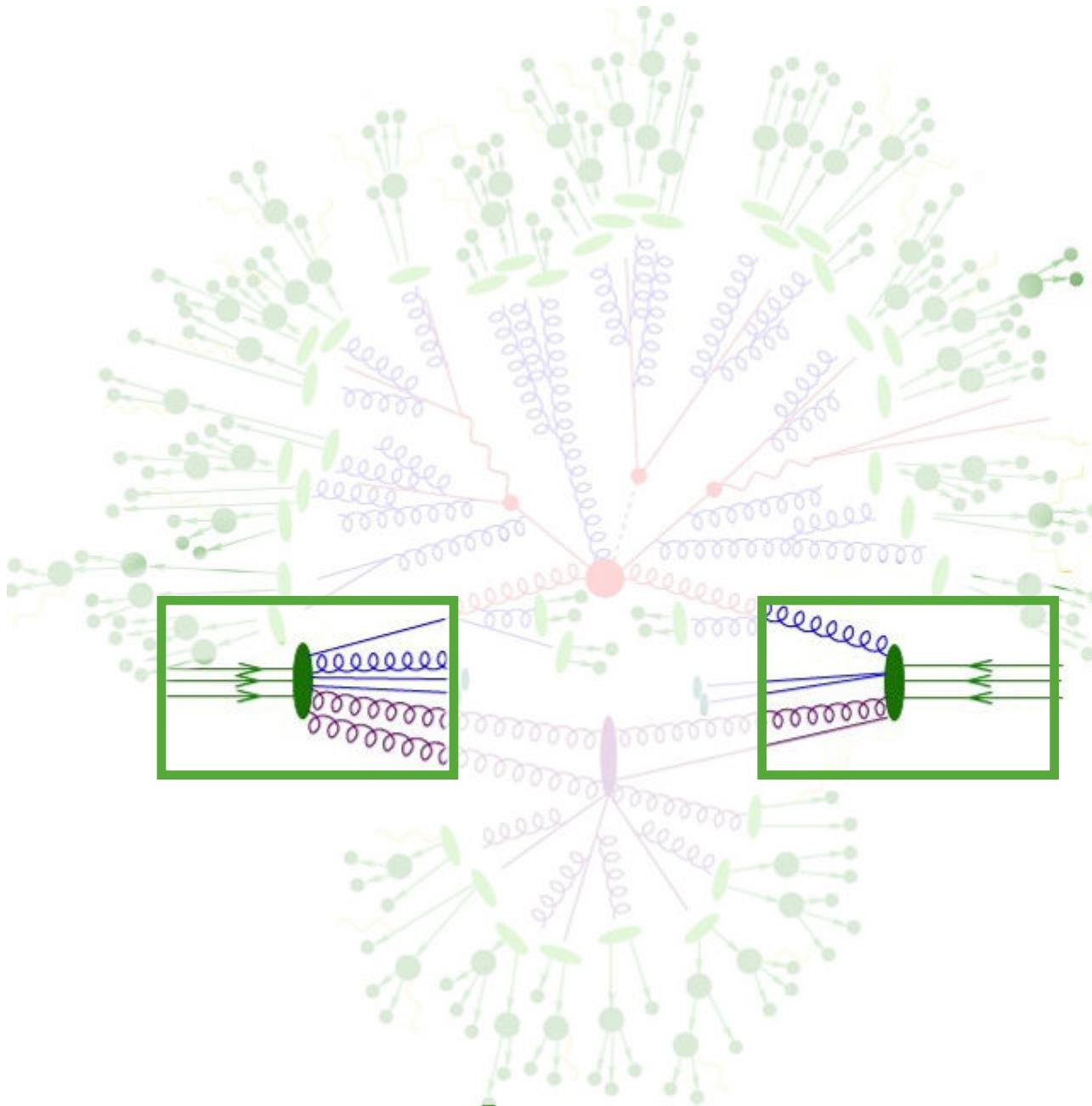
1.Why do we need to do MC simulation?

- Monte Carlo methods (or Monte Carlo experiments) are a broad class of computational algorithms that rely on repeated random sampling to obtain numerical results. Monte Carlo methods are mainly used in three distinct problem classes: optimization, numerical integration, and generating draws from a probability distribution.

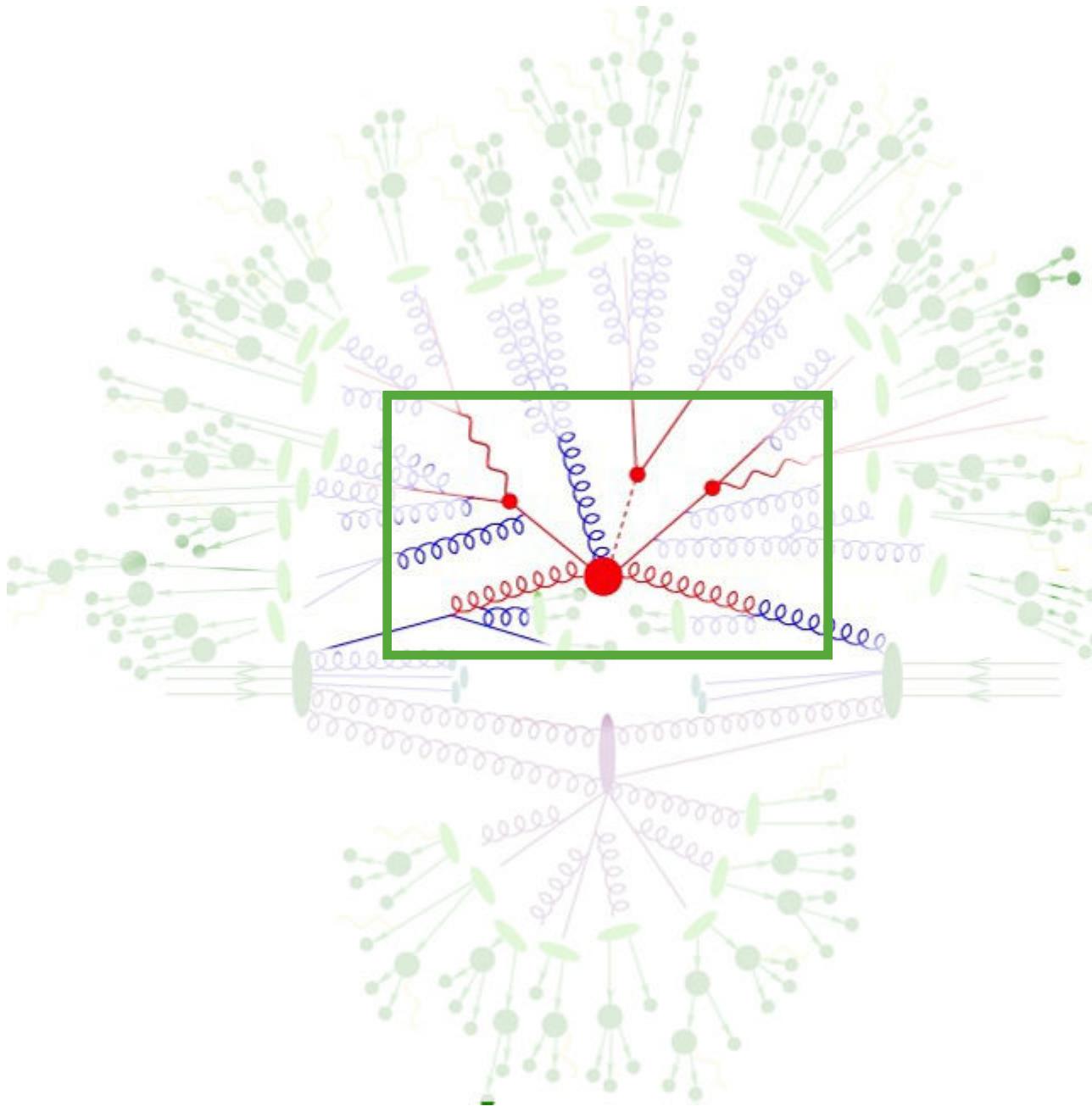


2. How to do MC simulation in collider physics?

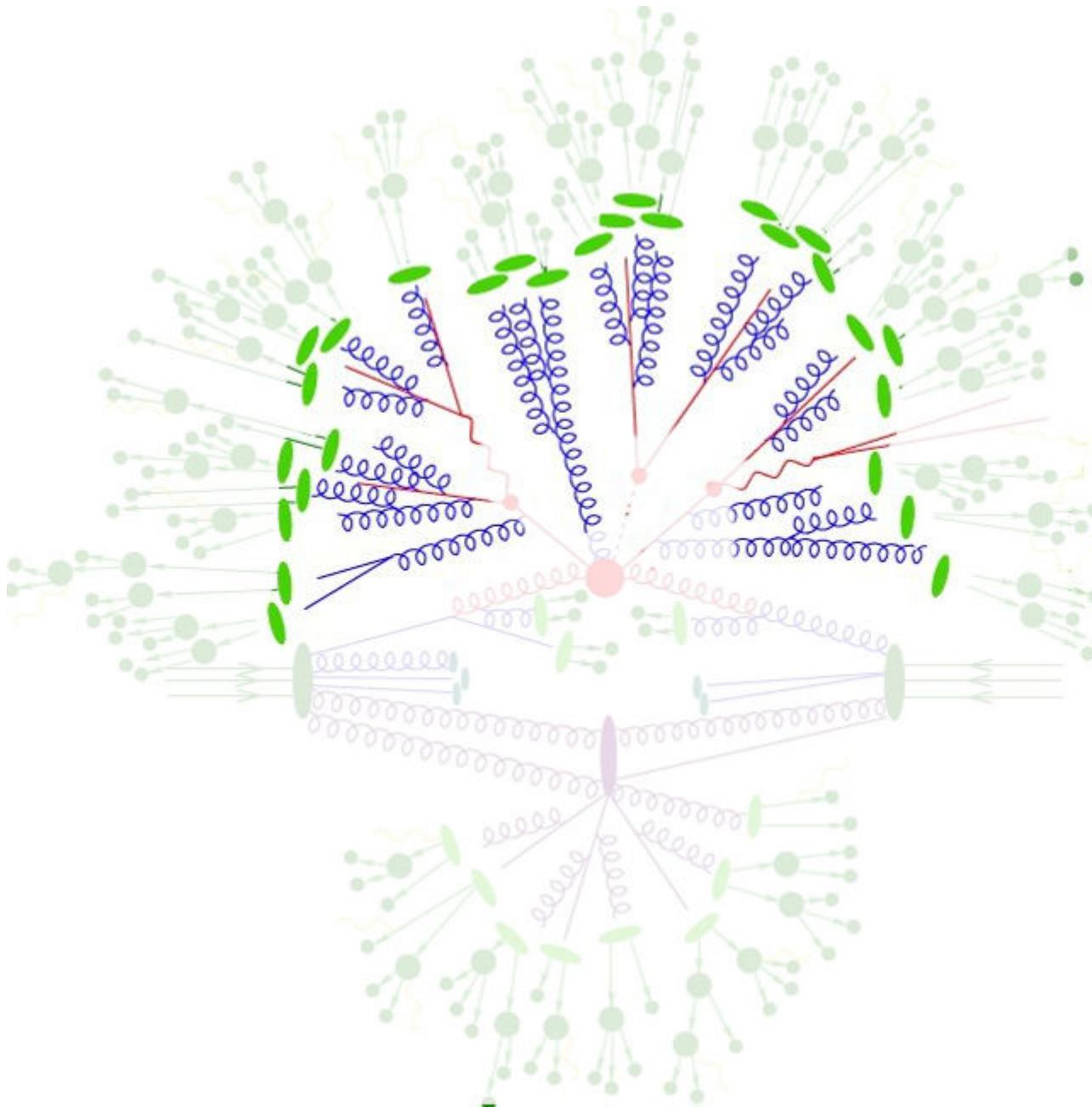




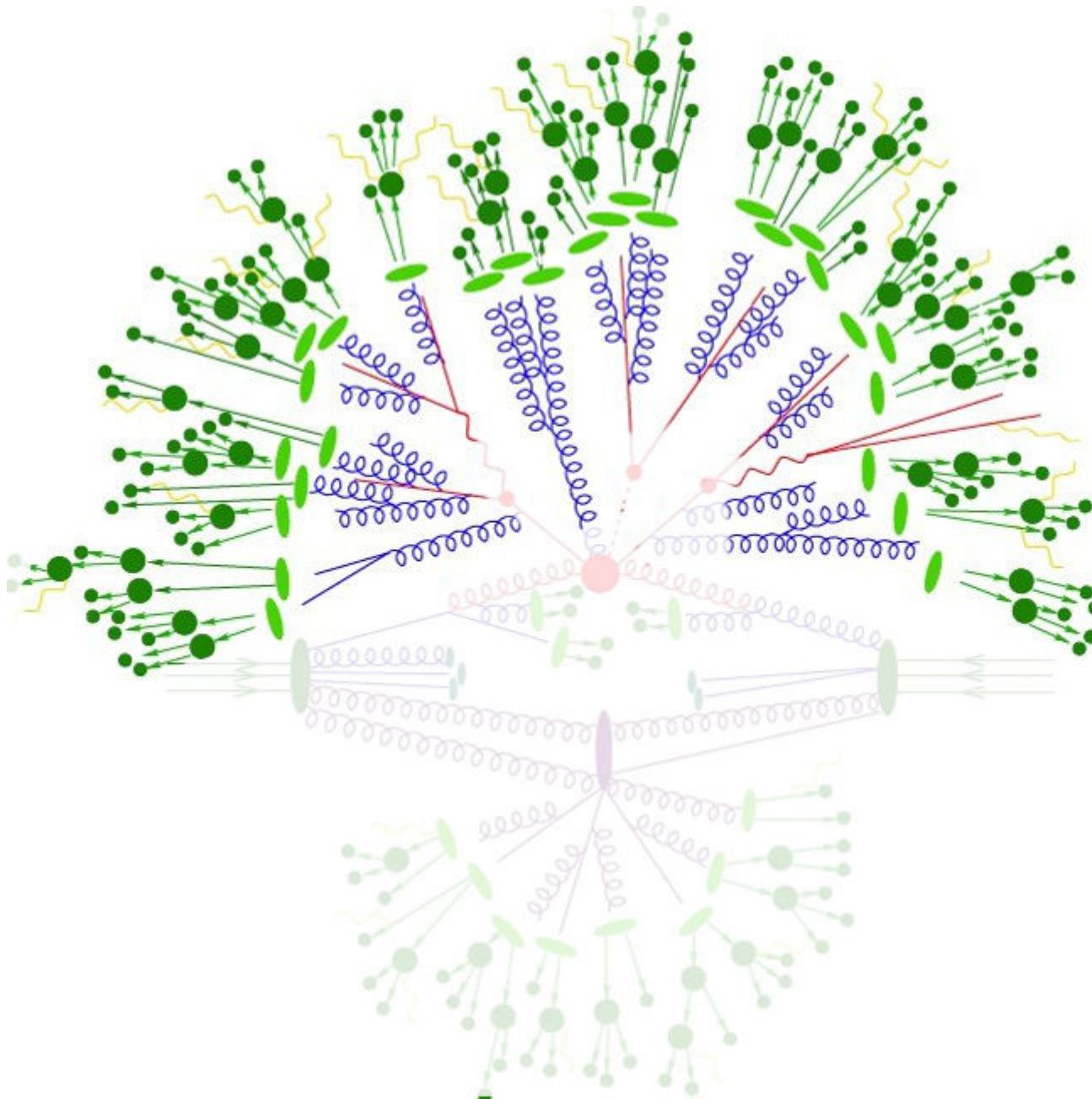
- a. PDF
 - CTEQ6L1
- b. Hard scattering
 - MadGraph/MadEvent, SHERPA, ...
- c. Parton Shower
- Hadronization
 - PYTHIA
- d. Detector
 - Delphes, PGS, GEANT4, ...
 - FastJet



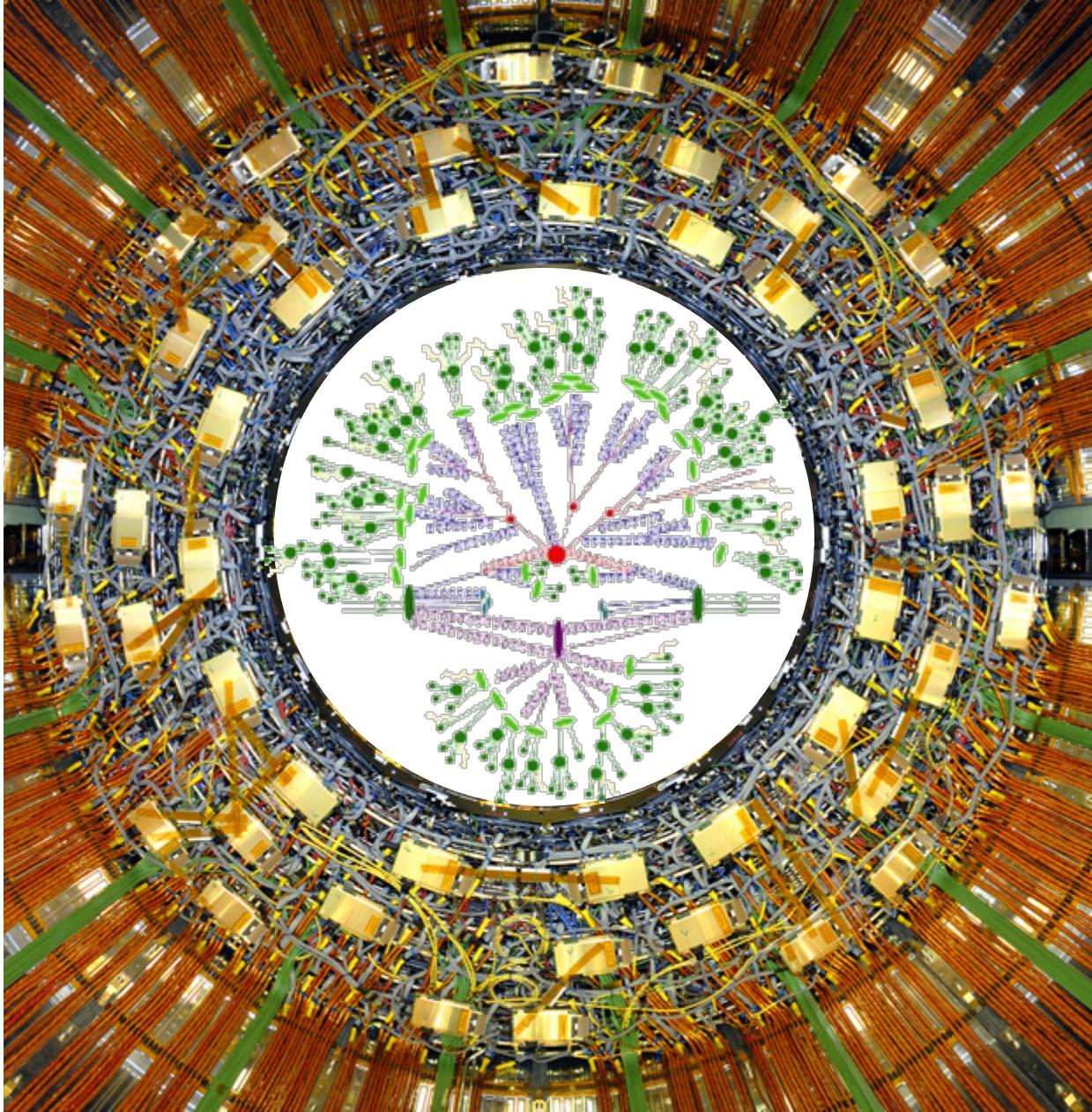
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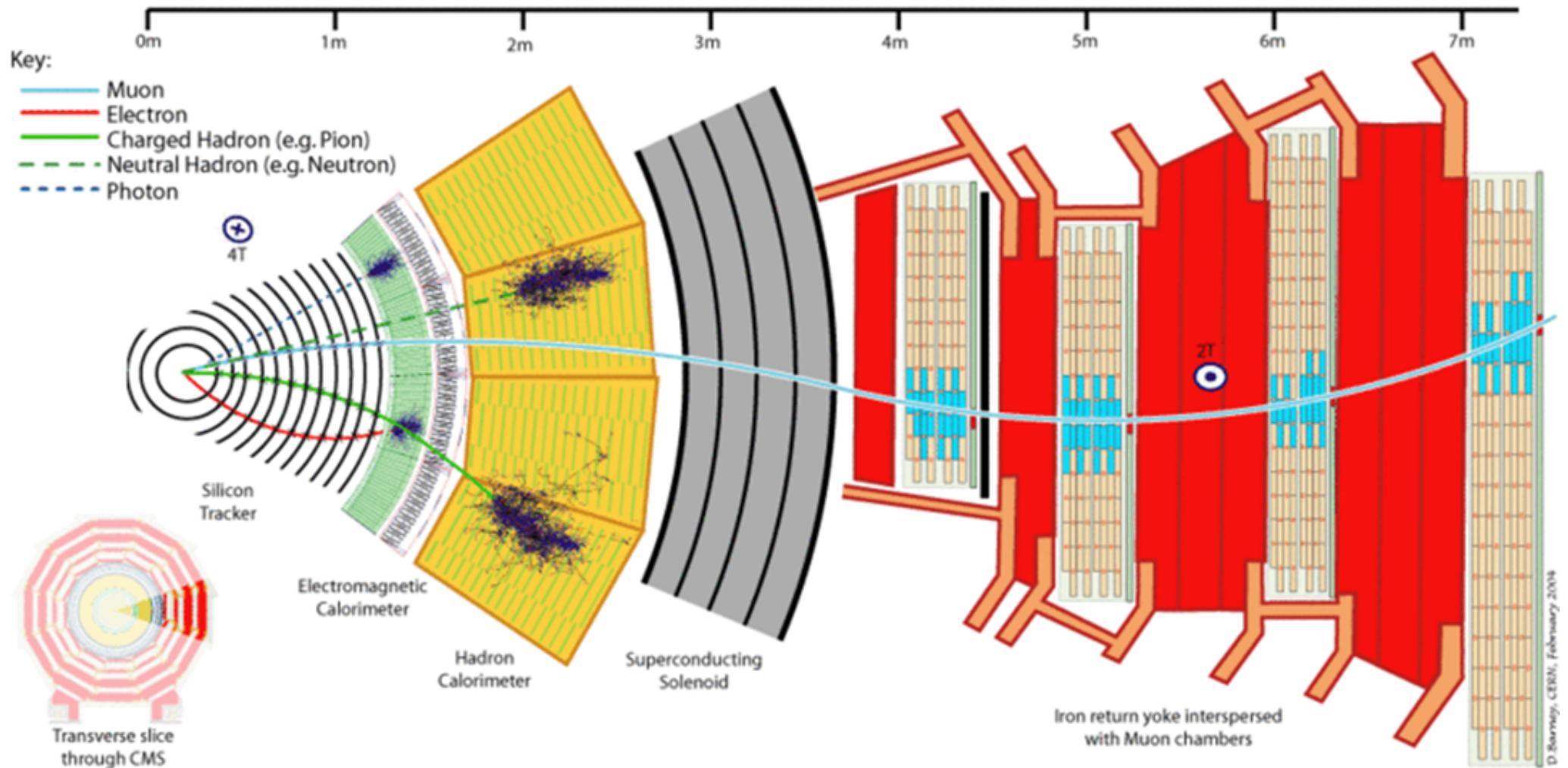


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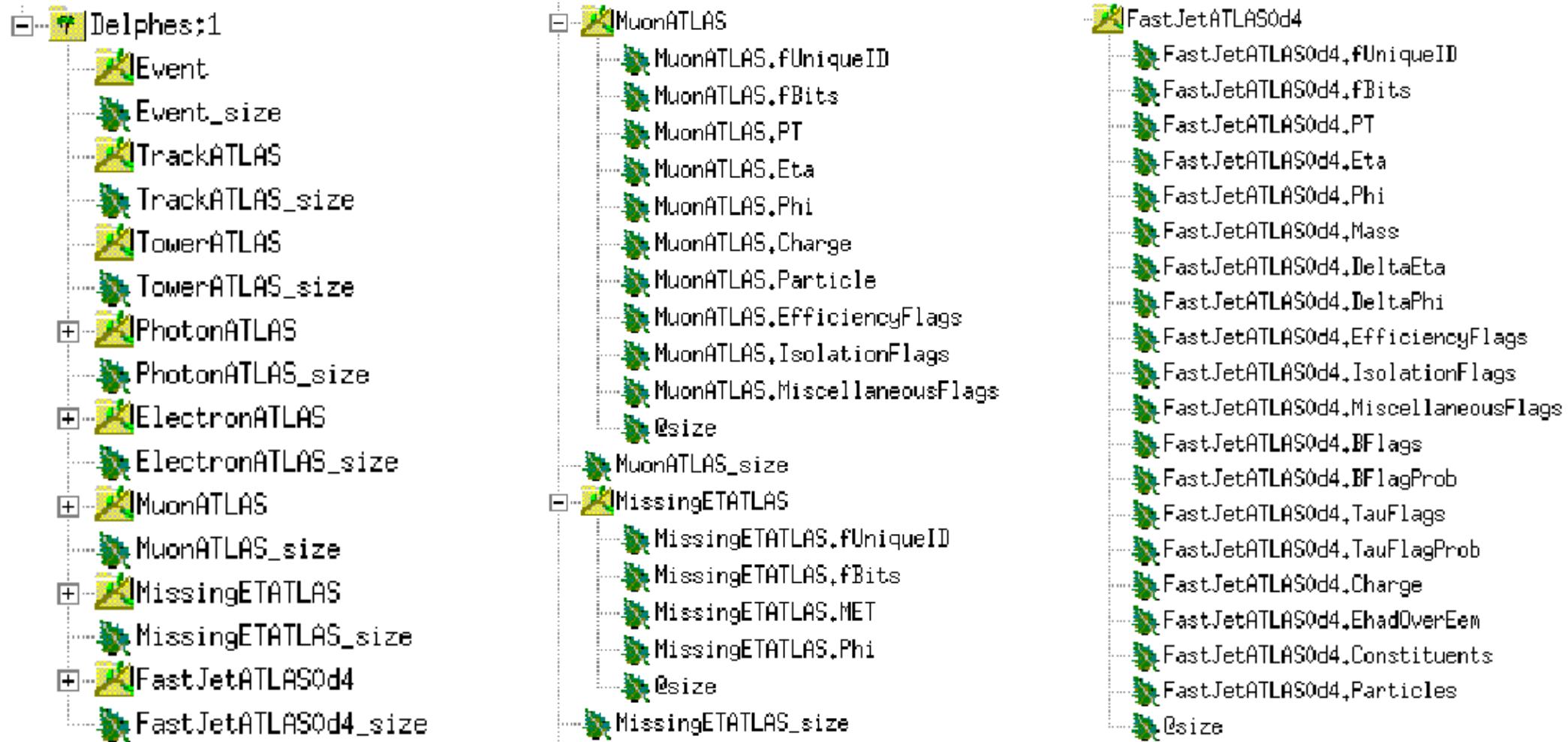


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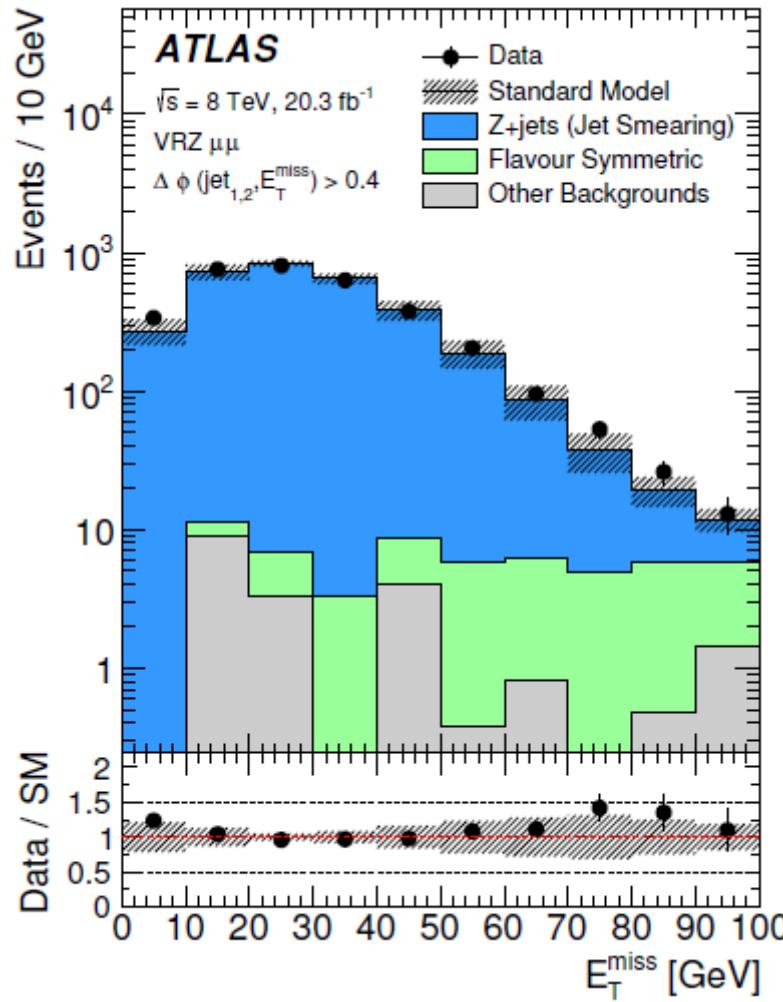
2. How to do MC simulation in collider physics?



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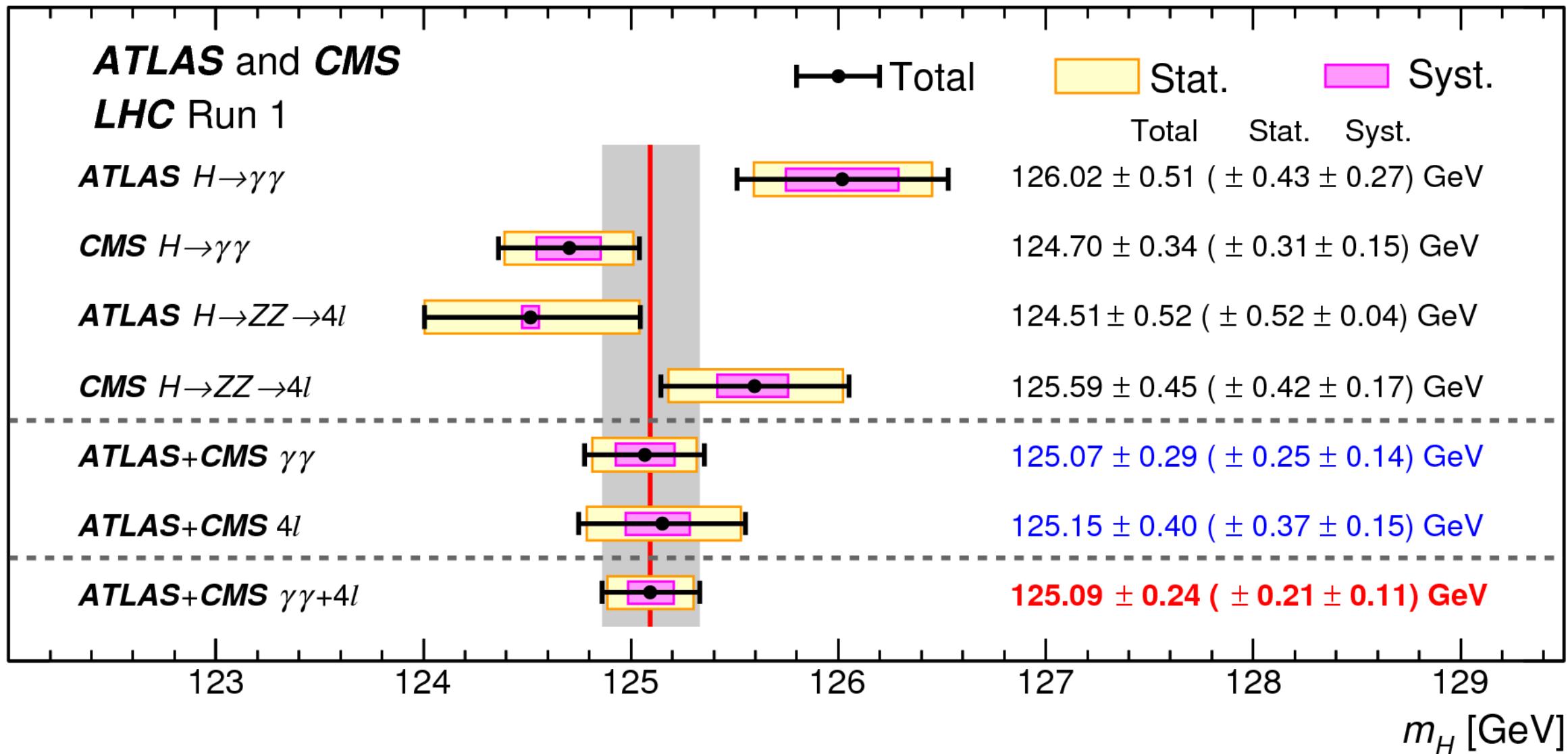


2. How to do MC simulation in collider physics?



Below-Z ($20 < m_{\ell\ell} < 80 \text{ GeV}$)	SR-2j-bveto ee
Observed events	30
Expected background events	$26 \pm 4 \pm 3$
Flavour-symmetric backgrounds	$24 \pm 4 \pm 3$
$Z/\gamma^* + \text{jets}$	$0.6 \pm 0.3 \pm 0.7$
Rare top	< 0.1
WZ/ZZ diboson	$0.6 \pm 0.2 \pm 0.1$
Fake leptons	$0.6 \pm 0.9 \pm 0.1$

2. How to do MC simulation in collider physics?



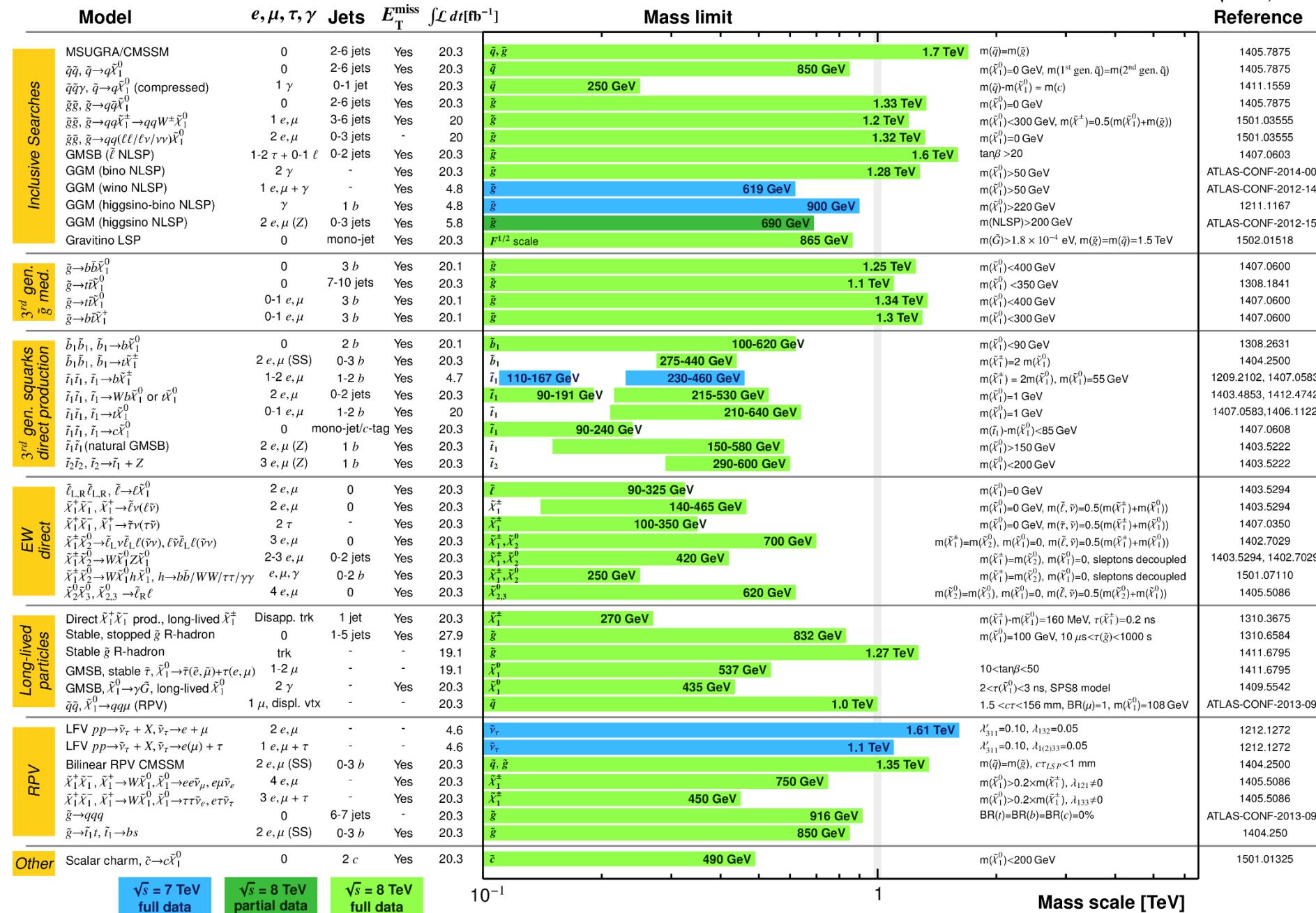
ATLAS SUSY Searches* - 95% CL Lower Limits

Status: Feb 2015

ATLAS Preliminary

$\sqrt{s} = 7, 8 \text{ TeV}$

Reference



$\sqrt{s} = 7 \text{ TeV}$
full data

$\sqrt{s} = 8 \text{ TeV}$
partial data

$\sqrt{s} = 8 \text{ TeV}$
full data

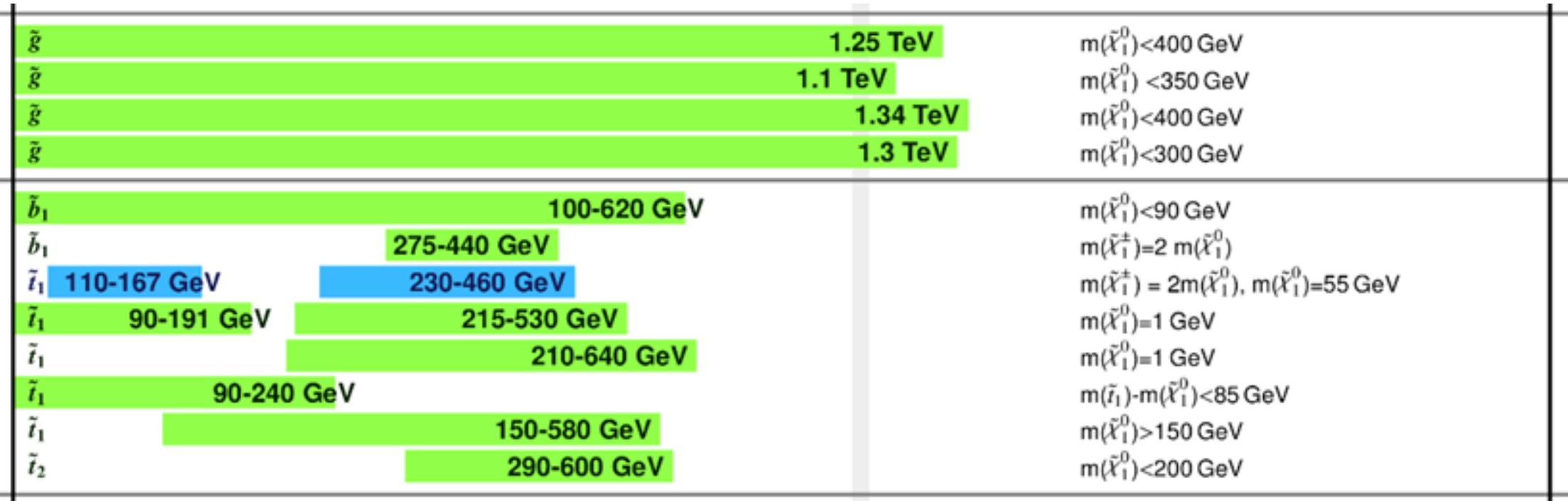
10⁻¹

1

Mass scale [TeV]

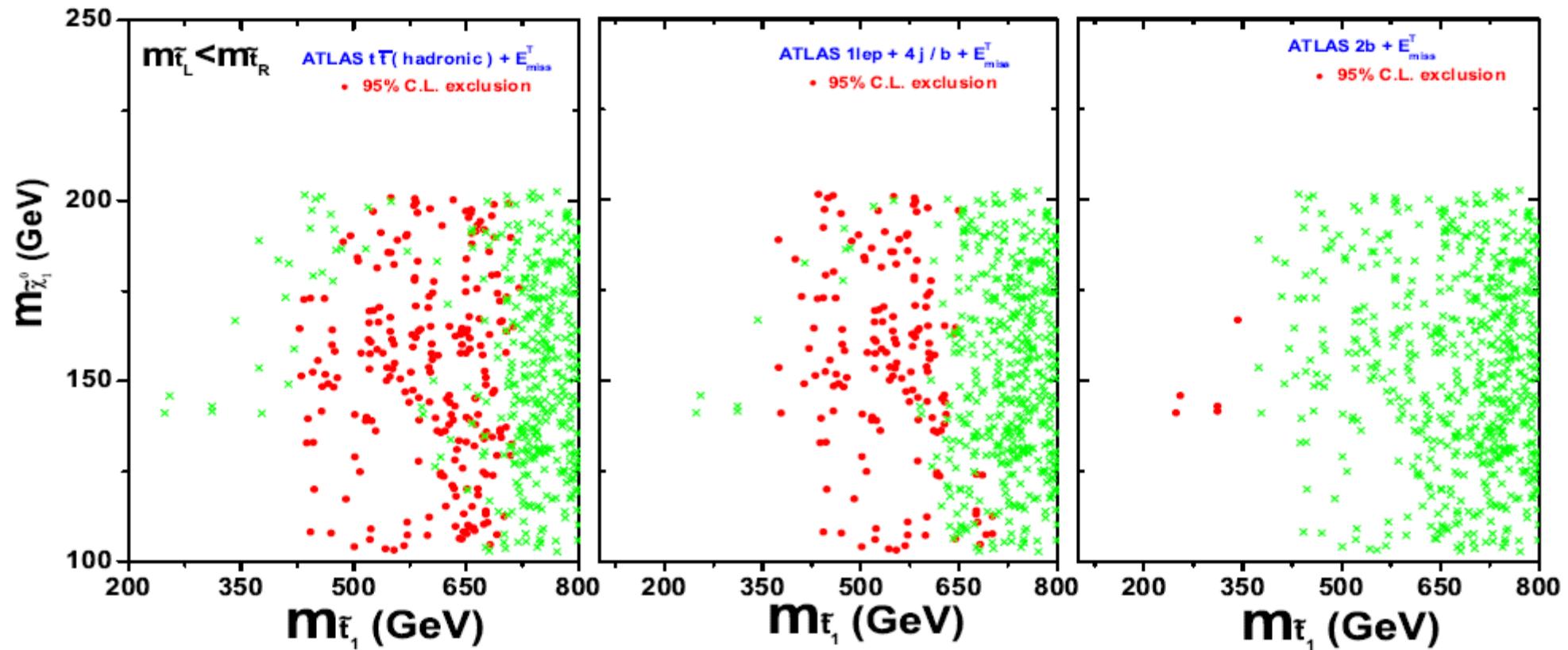
*Only a selection of the available mass limits on new states or phenomena is shown. All limits quoted are observed minus 1 σ theoretical signal cross section uncertainty.

Why do **we** need to do MC simulation?



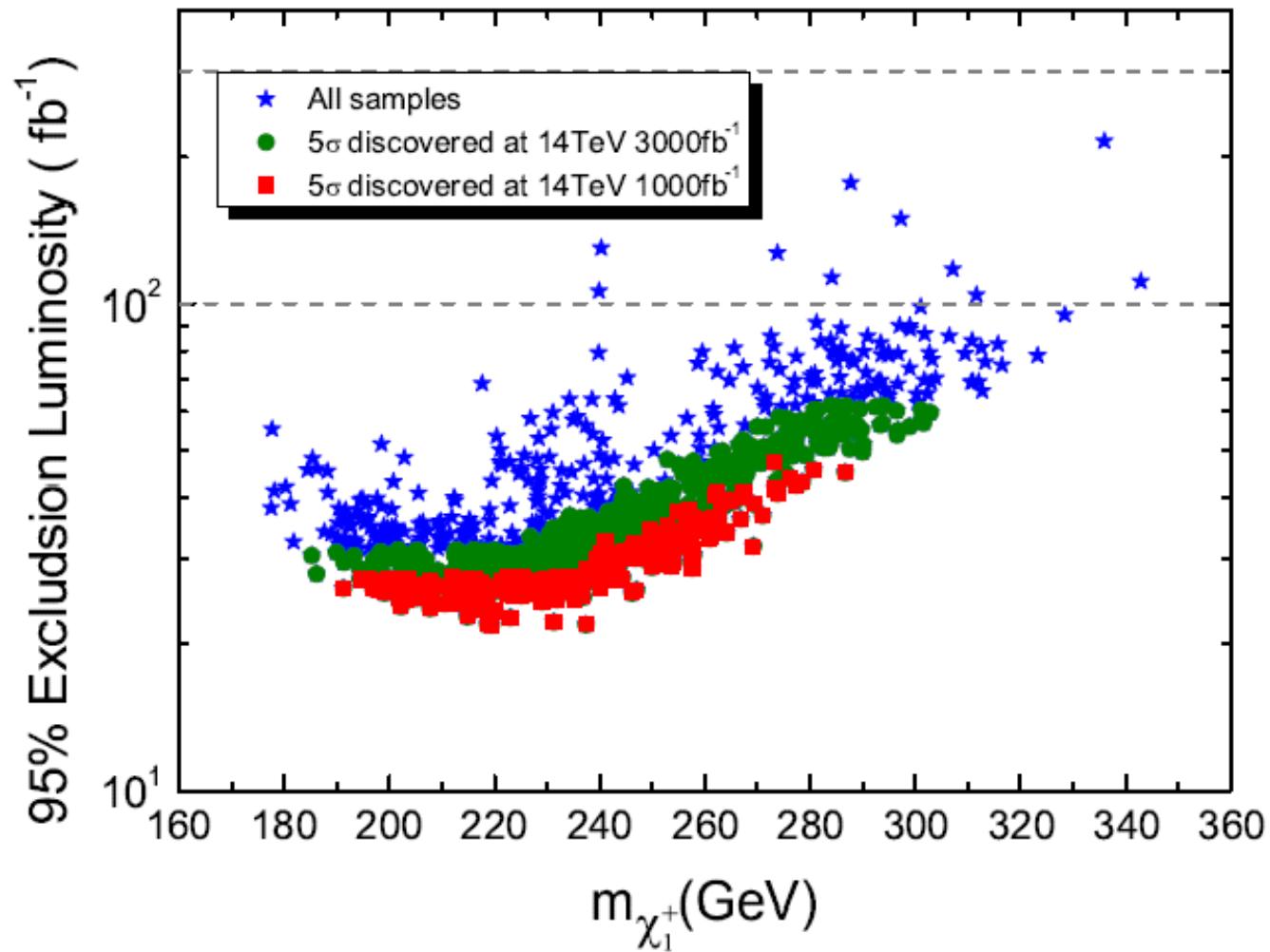
3. Which kind of work can MC simulation do?

arXiv:1308.5307

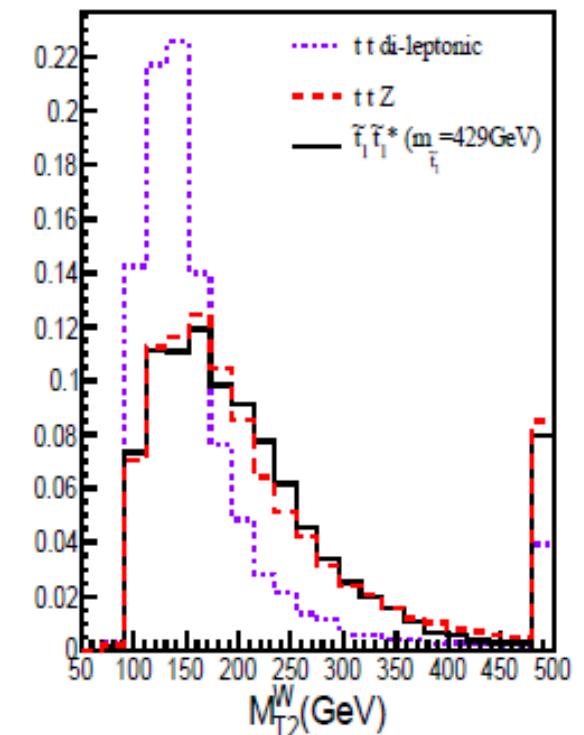
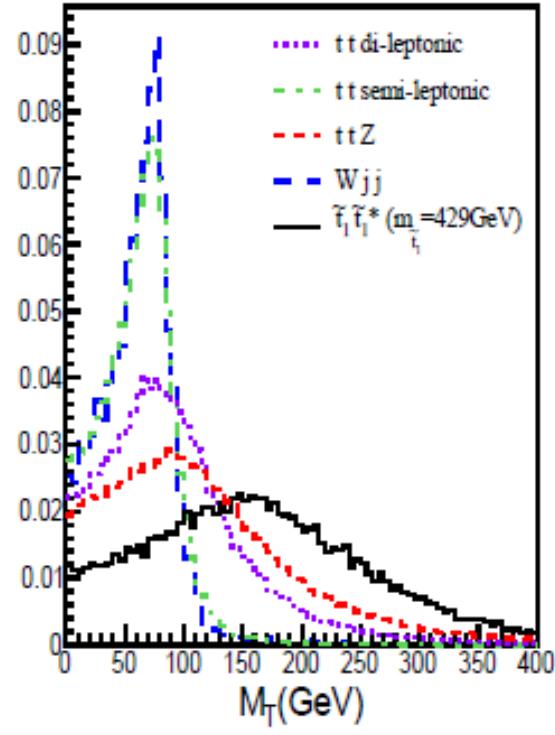
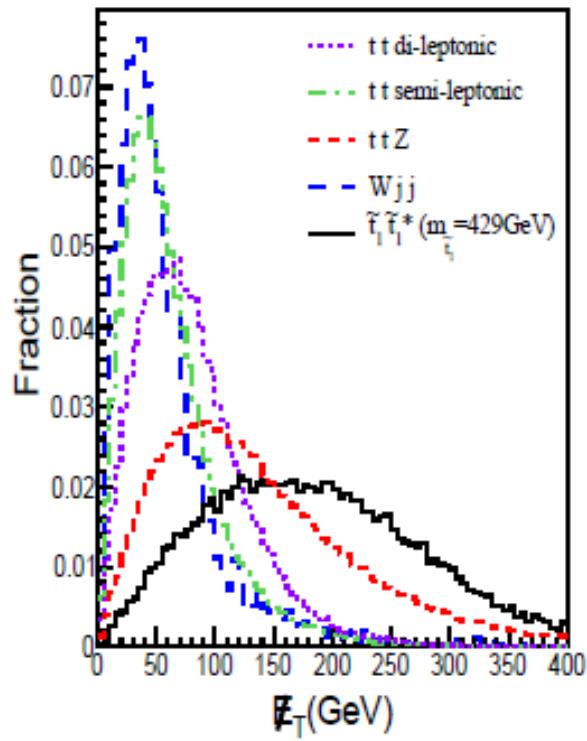


3. Which kind of work can MC simulation do?

arXiv:1410.3239

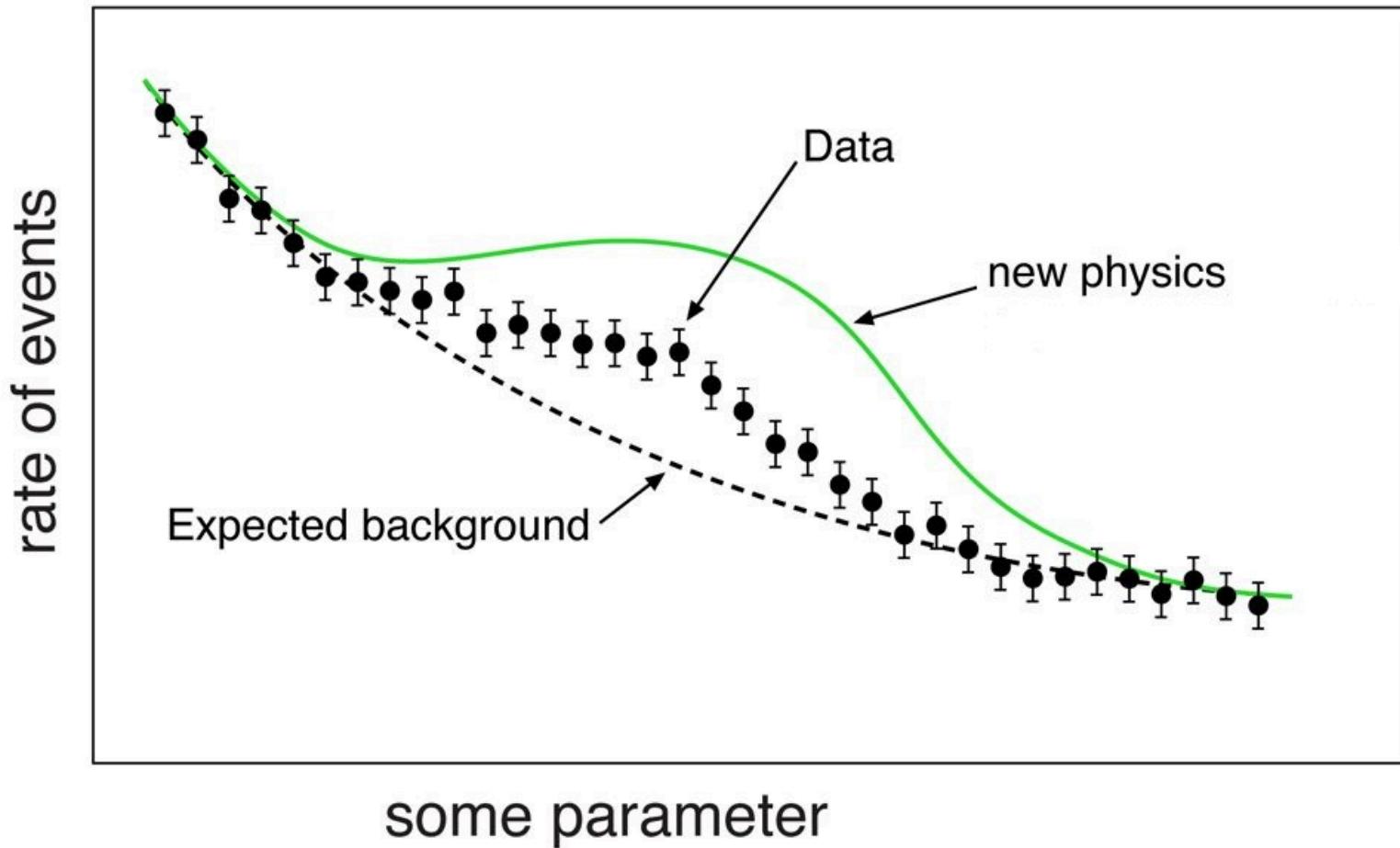


3. Which kind of work can MC simulation do?



arXiv:1206.3865

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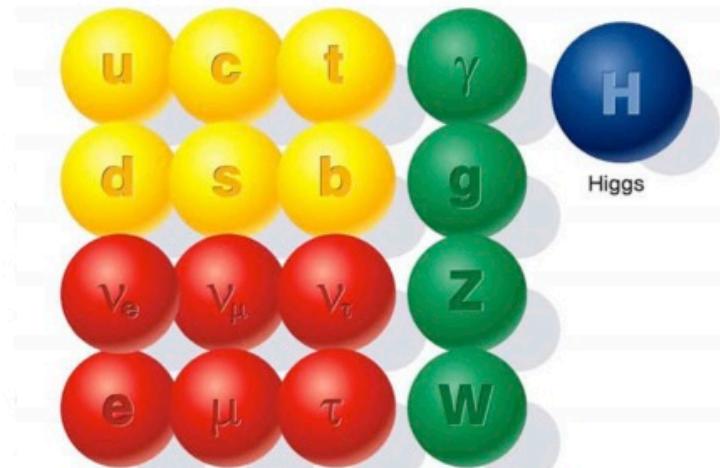
II Example



**Explanation of the
ATLAS Z-peaked
excess in the NM-**

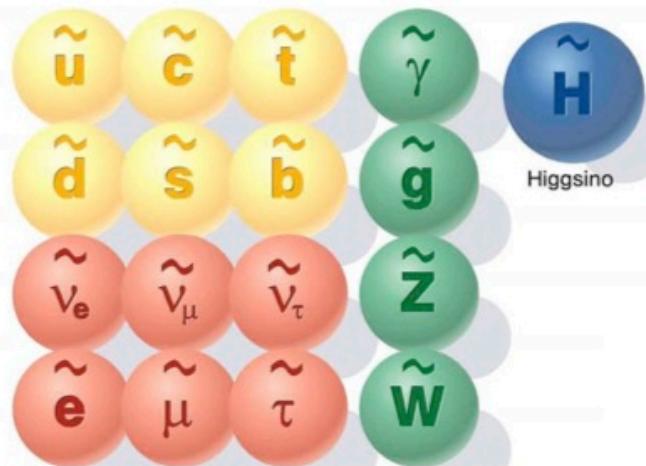
1. Brief introduction of SUSY

The known world of
Standard Model particles



- yellow circle: quarks
- red circle: leptons
- green circle: force carriers

The hypothetical world of
SUSY particles



- yellow circle: squarks
- red circle: sleptons
- green circle: SUSY force carriers

($\tilde{B}, \tilde{W}, \tilde{H}_u, \tilde{H}_d$)



($\tilde{\chi}_1^0, \tilde{\chi}_2^0, \tilde{\chi}_3^0, \tilde{\chi}_4^0$)

NMSSM

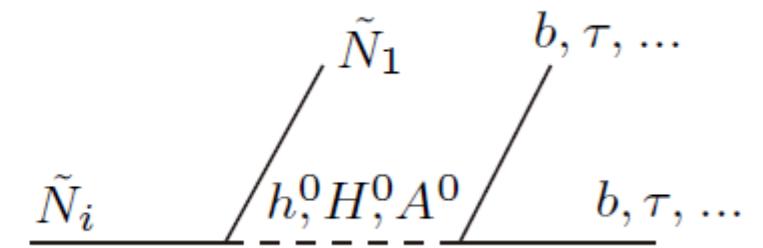
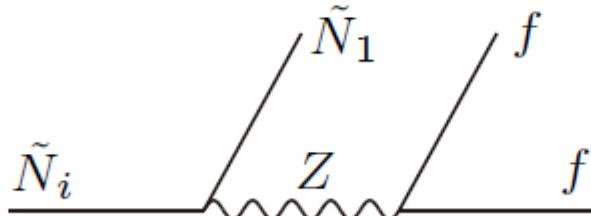
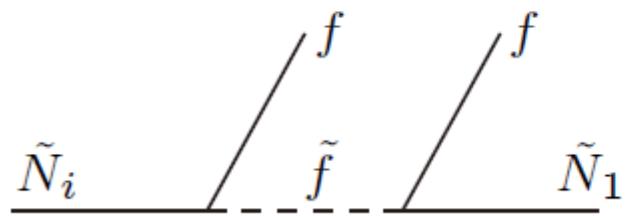
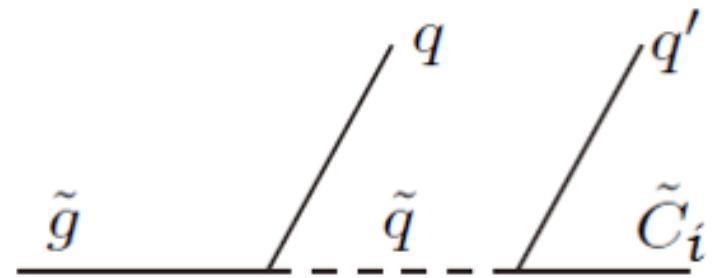
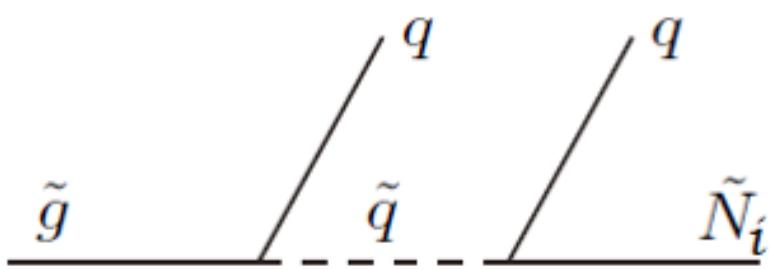
$$W_{NMSSM} = W_{MSSM} + \lambda \hat{S} \hat{H}_u \cdot \hat{H}_d + \xi_F \hat{S} + \frac{1}{2} \mu' \hat{S}^2 + \frac{\kappa}{3} \hat{S}^3$$

$$\mathcal{M} = \begin{pmatrix} M_1 & 0 & \frac{ev_u}{\sqrt{2}c_w} & -\frac{ev_d}{\sqrt{2}c_w} & 0 \\ 0 & M_2 & -\frac{ev_u}{\sqrt{2}s_w} & \frac{ev_d}{\sqrt{2}s_w} & 0 \\ \frac{ev_u}{\sqrt{2}c_w} & -\frac{ev_u}{\sqrt{2}s_w} & 0 & -\mu_{eff} & -\lambda v_d \\ -\frac{ev_d}{\sqrt{2}c_w} & \frac{ev_d}{\sqrt{2}s_w} & -\mu_{eff} & 0 & -\lambda v_u \\ 0 & 0 & -\lambda v_d & -\lambda v_u & 2\kappa s + \mu' \end{pmatrix}$$

$(\tilde{B}, \tilde{W}, \tilde{H}_u, \tilde{H}_d, \tilde{S})$

 $(\tilde{\chi}_1^0, \tilde{\chi}_2^0, \tilde{\chi}_3^0, \tilde{\chi}_4^0, \tilde{\chi}_5^0)$

Sparticle Decay

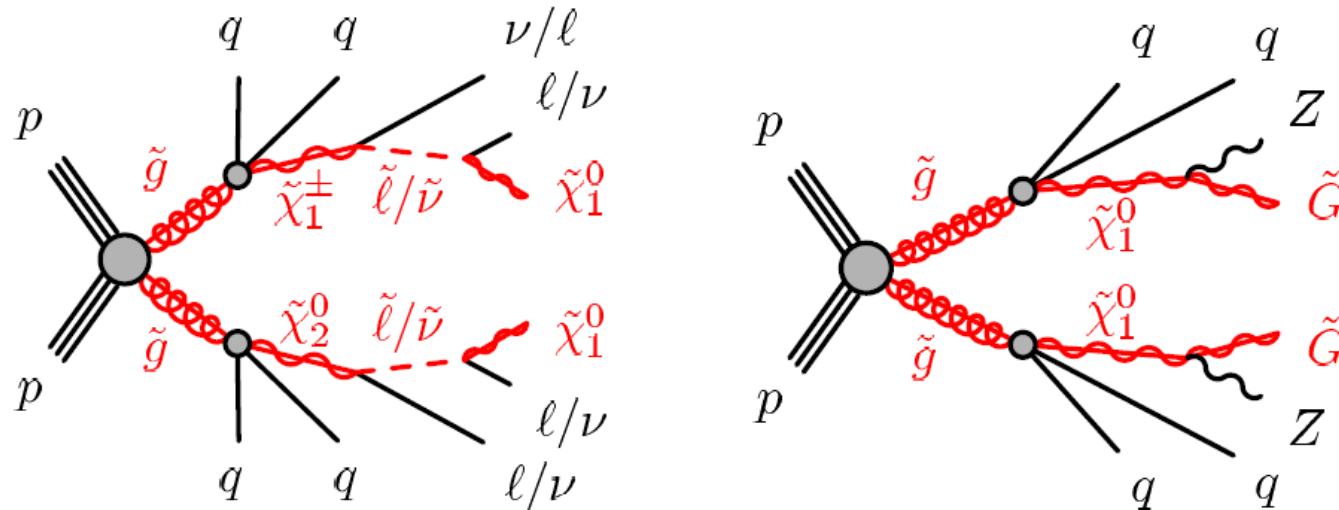


Search for supersymmetry in events containing a same-flavour opposite-sign dilepton pair, jets, and large missing transverse momentum in $\sqrt{s} = 8$ TeV pp collisions with the ATLAS detector

Two searches for supersymmetric particles in final states containing a same-flavour opposite-sign lepton pair, jets and large missing transverse momentum are presented. The proton–proton collision data used in these searches were collected at a centre-of-mass energy $\sqrt{s} = 8$ TeV by the ATLAS detector at the Large Hadron Collider and corresponds to an integrated luminosity of 20.3 fb^{-1} . Two leptonic production mechanisms are considered: decays of squarks and gluinos with Z bosons in the final state, resulting in a peak in the dilepton invariant mass distribution around the Z -boson mass; and decays of neutralinos (e.g. $\tilde{\chi}_2^0 \rightarrow \ell^+ \ell^- \tilde{\chi}_1^0$), resulting in a kinematic endpoint in the dilepton invariant mass distribution. For the former, an excess of events above the expected Standard Model background is observed, with a significance of 3 standard deviations. In the latter case, the data are well-described by the expected Standard Model background. The results from each channel are interpreted in the context of several supersymmetric models involving the production of squarks and gluinos.

1. The ATLAS report

general gauge mediation models



Physics process	Generator	Parton Shower	Cross section	Tune	PDF set
$Z/\gamma^*(\rightarrow \ell\ell) + \text{jets}$	SHERPA 1.4.1	SHERPA 1.4.1	NNLO [29, 30]	SHERPA default	NLO CT10 [31]
$t\bar{t}$	POWHEG-BOX r2129	PyTHIA 6.426	NNLO+NNLL [32, 33]	PERUGIA2011C	NLO CT10
Single-top (Wt)	POWHEG-BOX r1556	PyTHIA 6.426	Approx. NNLO [34, 35]	PERUGIA2011C	NLO CT10
$t + Z$	MADGRAPH5 1.3.28	PyTHIA 6.426	LO	AUET2	CTEQ6L1 [36]
$t\bar{t} + W$ and $t\bar{t} + Z$	MADGRAPH5 1.3.28	PyTHIA 6.426	NLO [37, 38]	AUET2	CTEQ6L1
$t\bar{t} + WW$	MADGRAPH5 1.3.28	PyTHIA 8.165	LO	AUET2	CTEQ6L1
WW , WZ and ZZ	POWHEG-BOX r1508	PyTHIA 8.163	NLO [39, 40]	AUET2	NLO CT10

Electron candidates are reconstructed using energy clusters in the electromagnetic calorimeter matched to ID tracks. Electrons used in this analysis are assigned either “baseline” or “signal” status. Baseline electrons are required to have transverse energy $E_T > 10$ GeV, satisfy the “medium” criteria described in Ref. [64] and reside within $|\eta| < 2.47$ and not in the range $1.37 < |\eta| < 1.52$. Signal electrons are further

Baseline muons are reconstructed from either ID tracks matched to muon segments or combined tracks formed from the ID and muon spectrometer [65]. They are required to be of good quality, as described in Ref. [66], and to satisfy $p_T > 10$ GeV and $|\eta| < 2.4$. Signal muons are further required to be isolated,

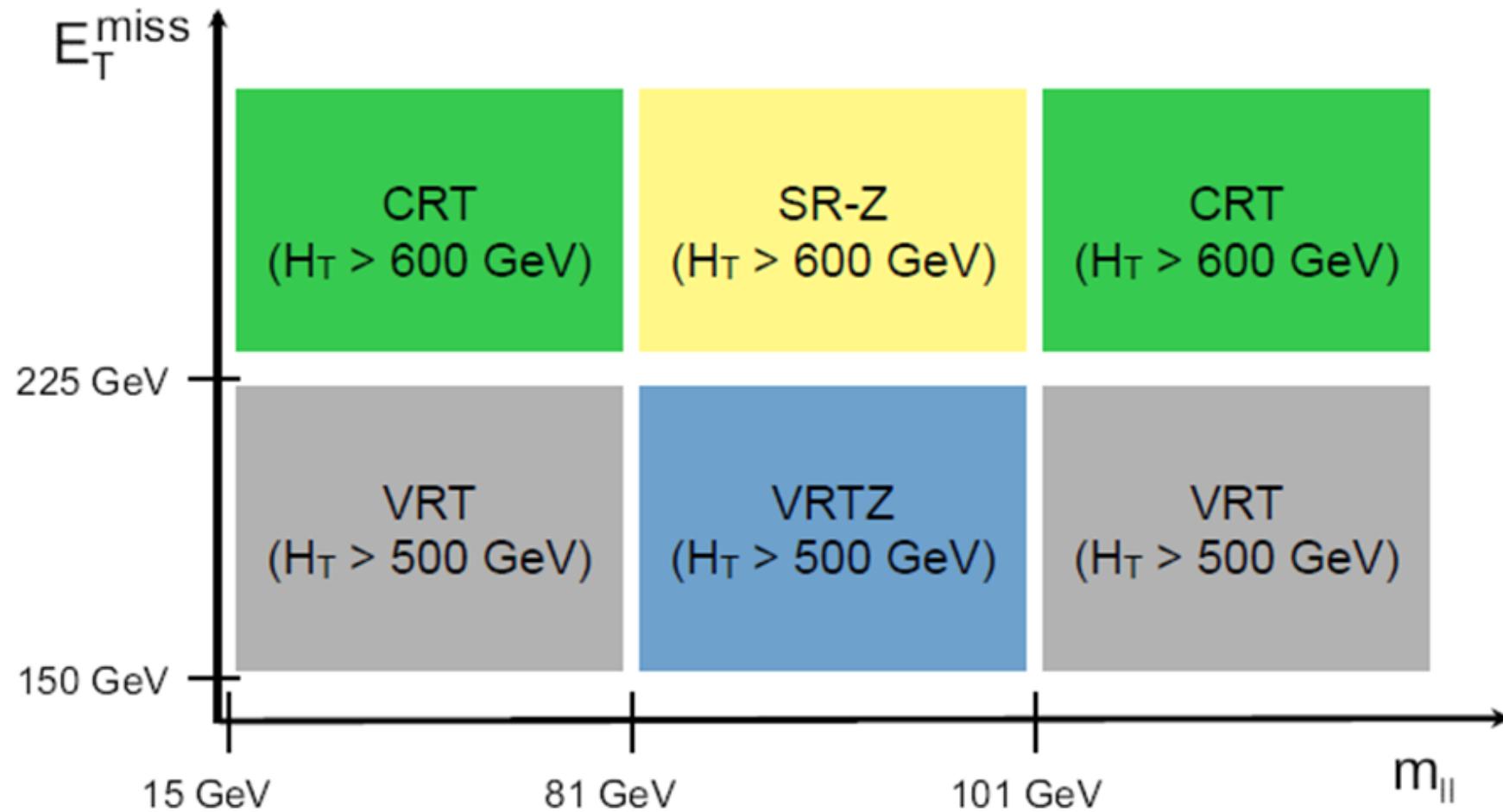
Jets are reconstructed from topological clusters in the calorimeter using the anti- k_t algorithm [67] with a distance parameter of 0.4. Each cluster is categorised as being electromagnetic or hadronic in origin between response in data and MC simulation. Baseline jets are selected with $p_T > 20$ GeV. Events in which these jets do not pass specific jet quality requirements are rejected so as to remove events affected by detector noise and non-collision backgrounds [72, 73]. Signal jets are required to satisfy $p_T > 35$ GeV and $|\eta| < 2.5$. To avoid the inclusion of jets resulting from pile-up, jets with $p_T < 50$ GeV within $|\eta| < 2.4$

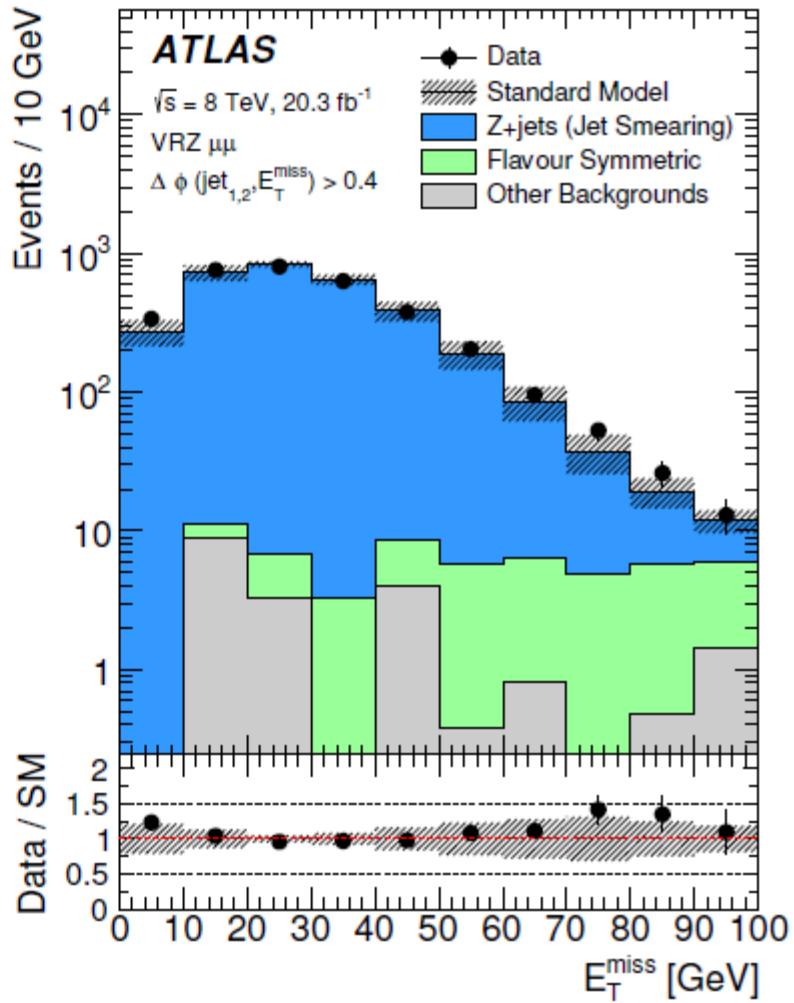
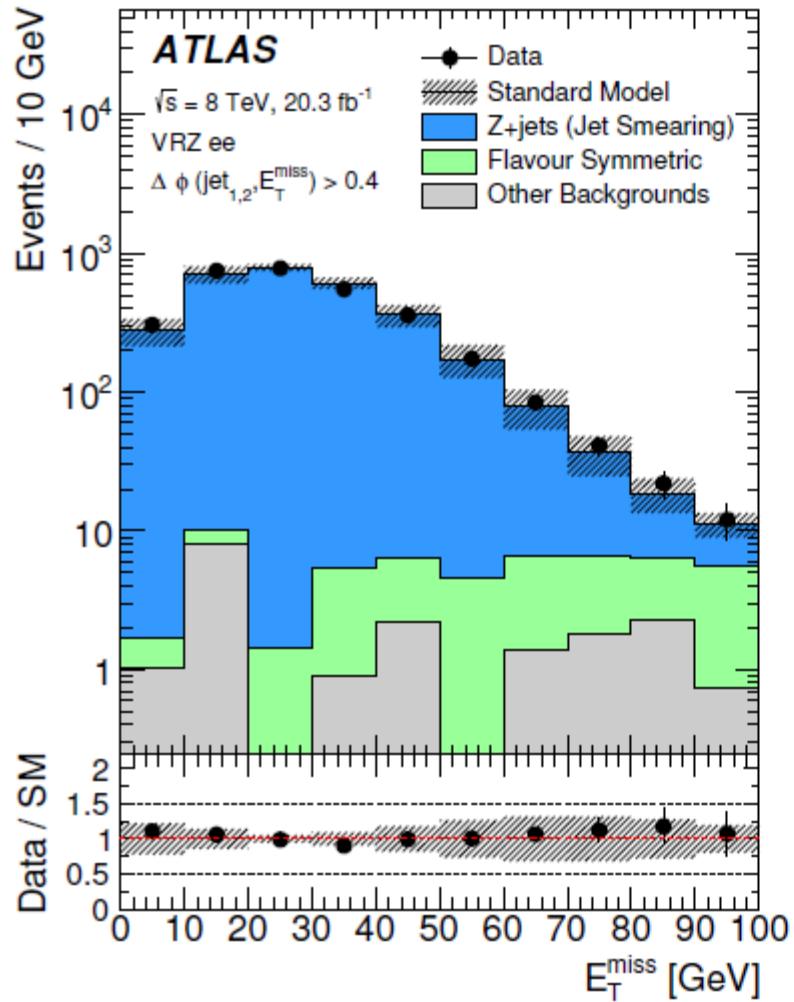
of associated tracks and any reconstructed secondary vertices. For this analysis, the working point corresponding to a 60 % efficiency for tagging b -jets in simulated $t\bar{t}$ events is used, resulting in a charm quark

Event selection

On-Z Region	E_T^{miss} [GeV]	H_T [GeV]	n_{jets}	$m_{\ell\ell}$ [GeV]	SF/DF	$E_T^{\text{miss}} \text{ sig.}$ [$\sqrt{\text{GeV}}$]	f_{ST}	$\Delta\phi(\text{jet}_{12}, E_T^{\text{miss}})$
Signal regions								
SR-Z	> 225	> 600	≥ 2	$81 < m_{\ell\ell} < 101$	SF	-	-	> 0.4
Control regions								
Seed region	-	> 600	≥ 2	$81 < m_{\ell\ell} < 101$	SF	< 0.9	< 0.6	-
CRe μ	> 225	> 600	≥ 2	$81 < m_{\ell\ell} < 101$	DF	-	-	> 0.4
CRT	> 225	> 600	≥ 2	$m_{\ell\ell} \notin [81, 101]$	SF	-	-	> 0.4
Validation regions								
VRZ	< 150	> 600	≥ 2	$81 < m_{\ell\ell} < 101$	SF	-	-	-
VRT	150–225	> 500	≥ 2	$m_{\ell\ell} \notin [81, 101]$	SF	-	-	> 0.4
VRTZ	150–225	> 500	≥ 2	$81 < m_{\ell\ell} < 101$	SF	-	-	> 0.4

Event selection

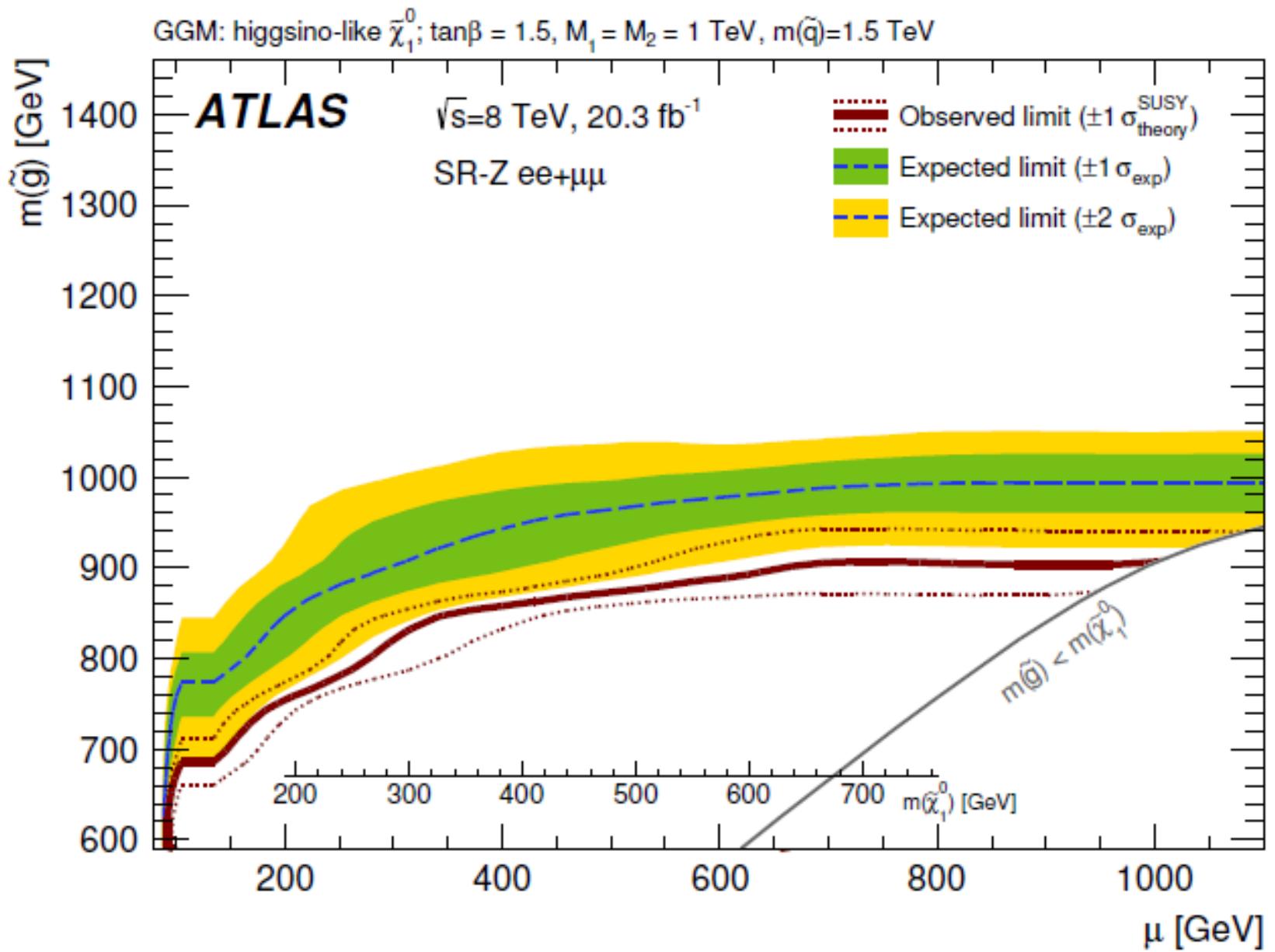




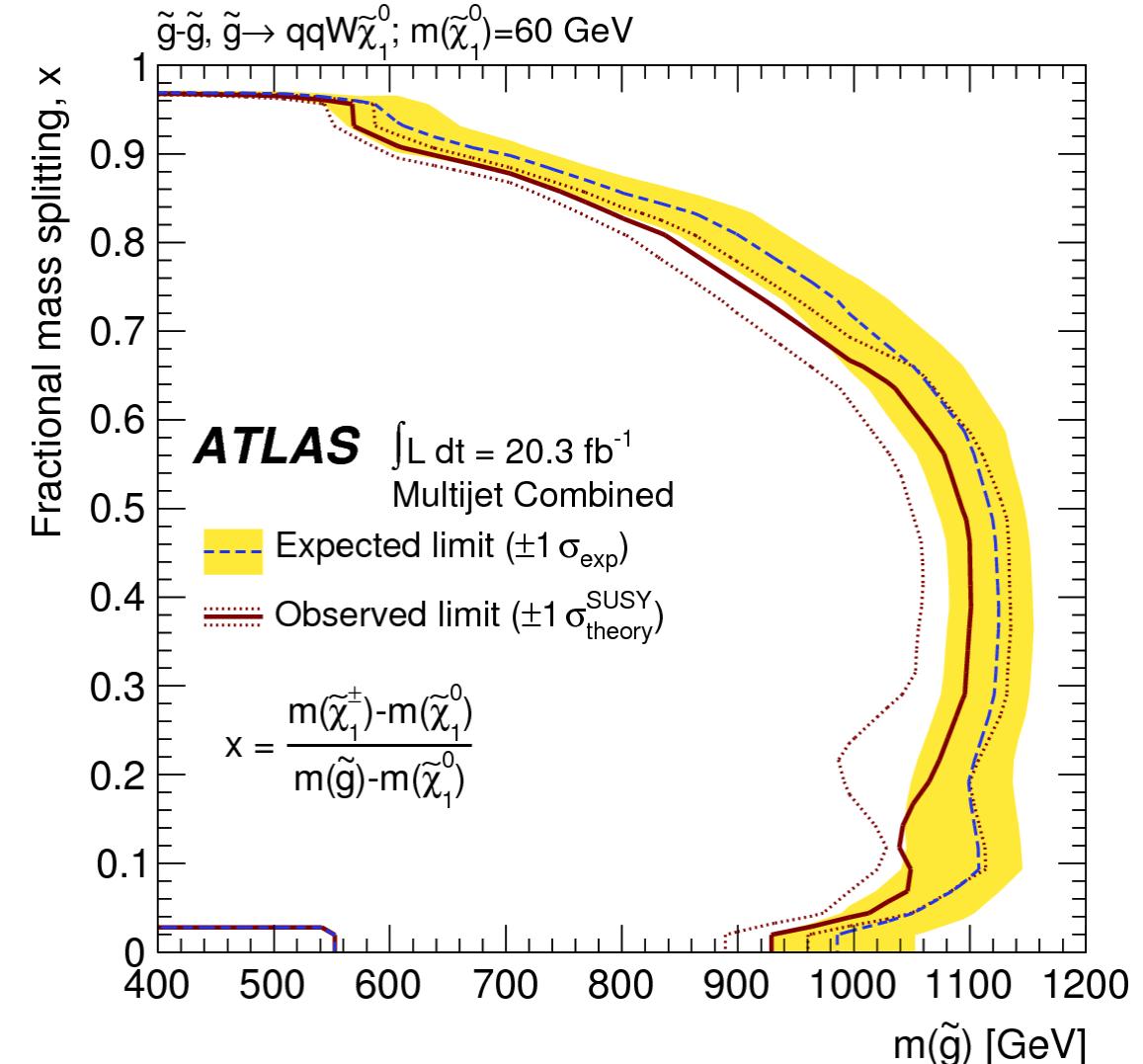
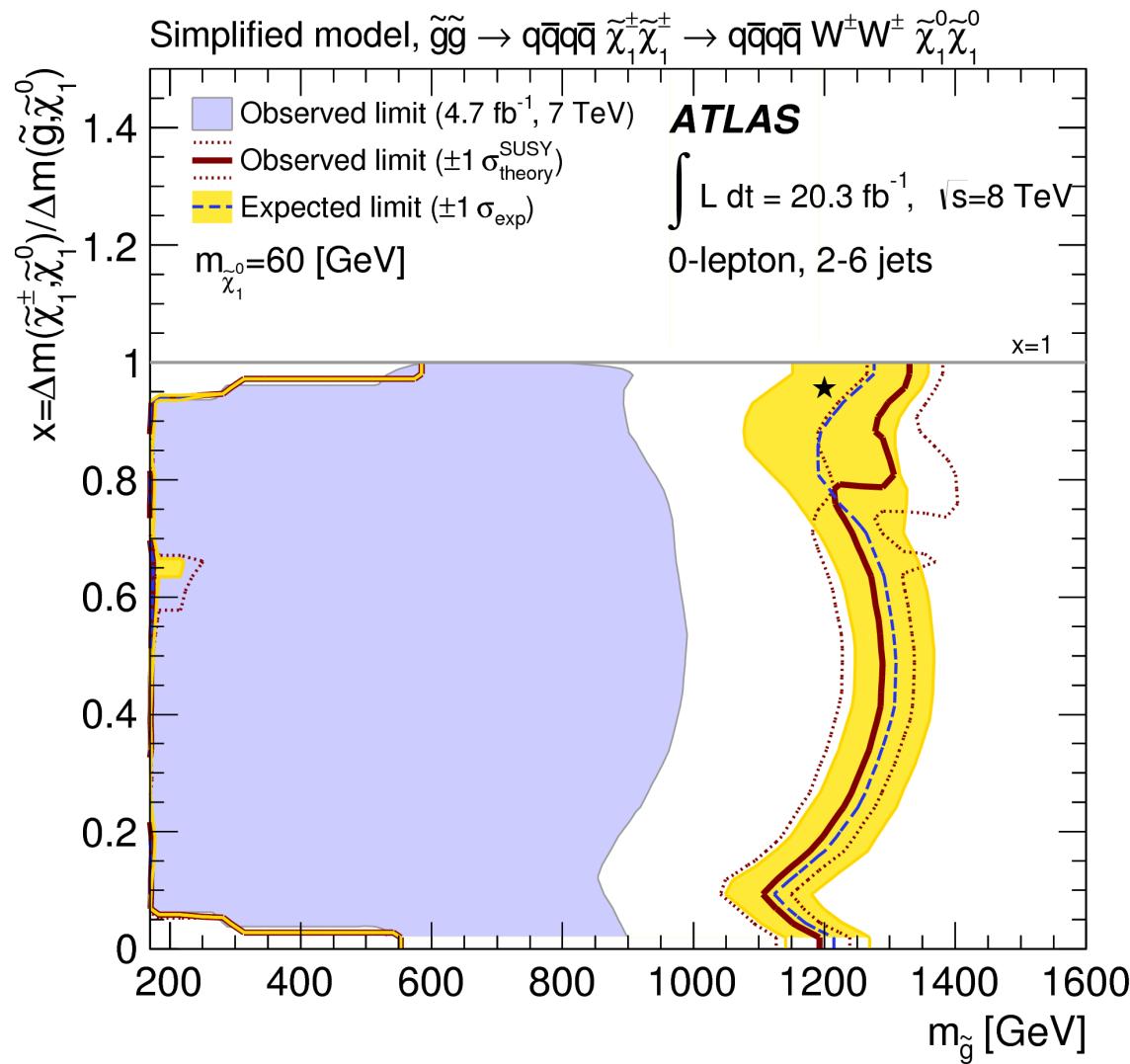
The jet smearing method using the data-corrected jet response function is validated in VRZ, comparing smeared pseudo-data to data. The resulting E_T^{miss} distributions show agreement within uncertainties

Result

Channel	SR-Z ee	SR-Z $\mu\mu$	SR-Z same-flavour combined				
Observed events	16	13	29				
Expected background events	4.2 ± 1.6	6.4 ± 2.2	10.6 ± 3.2				
Flavour-symmetric backgrounds	2.8 ± 1.4	3.3 ± 1.6	6.0 ± 2.6				
Z/γ^* + jets (jet-smearing)	0.05 ± 0.04	$0.02^{+0.03}_{-0.02}$	0.07 ± 0.05				
Rare top	0.18 ± 0.06	0.17 ± 0.06	0.35 ± 0.12				
WZ/ZZ diboson	1.2 ± 0.5	1.7 ± 0.6	2.9 ± 1.0				
Fake leptons	$0.1^{+0.7}_{-0.1}$	$1.2^{+1.3}_{-1.2}$	$1.3^{+1.7}_{-1.3}$				
<hr/>							
Signal region	Channel	$\langle \epsilon \sigma \rangle_{\text{obs}}^{95}$ [fb]	S_{obs}^{95}	S_{exp}^{95}	CL_B	$p(s=0)$	Gaussian significance
SR-Z	$ee + \mu\mu$	1.46	29.6	12^{+5}_{-2}	0.998	0.0013	3.0
	ee	1.00	20.2	8^{+4}_{-2}	0.998	0.0013	3.0
	$\mu\mu$	0.72	14.7	9^{+4}_{-2}	0.951	0.0430	1.7



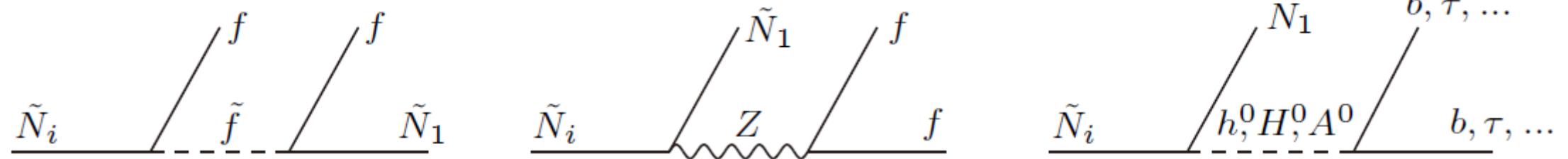
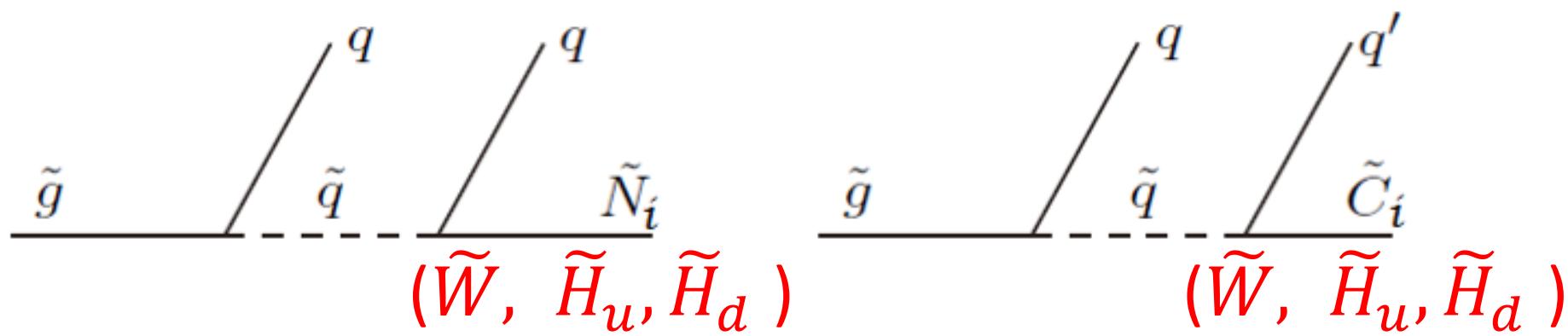
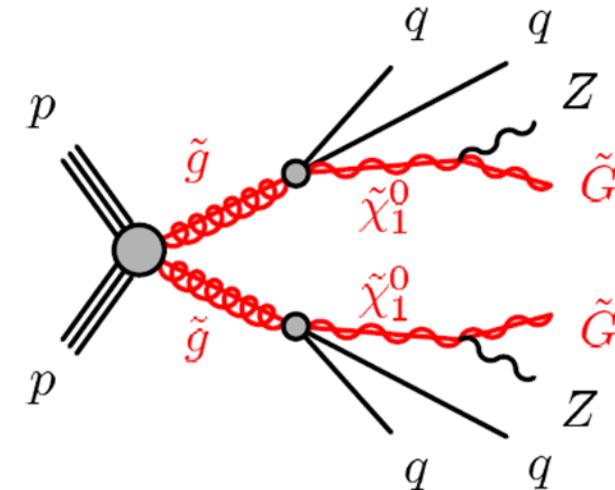
3. Explanation



0 lepton, 2~6 jets

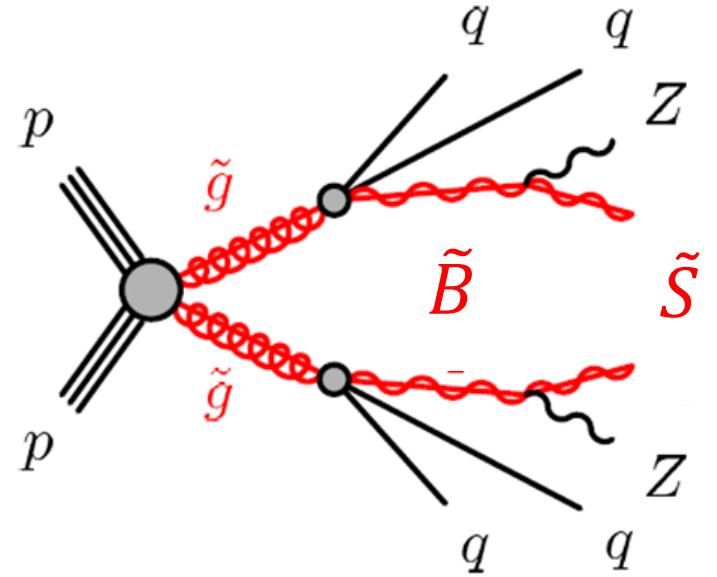
Requirement	Signal Region							
	2jl	2jm	2jt	2jW	3j	4jW		
$E_T^{\text{miss}} [\text{GeV}] >$	160							
$p_T(j_1) [\text{GeV}] >$	130							
$p_T(j_2) [\text{GeV}] >$	60							
$p_T(j_3) [\text{GeV}] >$	–		60		40			
$p_T(j_4) [\text{GeV}] >$	–			40				
$\Delta\phi(\text{jet}_{1,2,(3)}, E_T^{\text{miss}})_{\min} >$	0.4							
$\Delta\phi(\text{jet}_{i>3}, E_T^{\text{miss}})_{\min} >$	–			0.2				
W candidates	–		$2(W \rightarrow j)$	–	$(W \rightarrow j) + (W \rightarrow jj)$			
$E_T^{\text{miss}} / \sqrt{H_T} [\text{GeV}^{1/2}] >$	8	15	–					
$E_T^{\text{miss}} / m_{\text{eff}}(N_j) >$	–		0.25	0.3	0.35			
$m_{\text{eff}}(\text{incl.}) [\text{GeV}] >$	800	1200	1600	1800	2200	1100		

3. Explanation



NMSSM

$$\begin{aligned}
\mathcal{L}_{\tilde{\chi}^0} = & \tilde{u}_L^* \tilde{\chi}_j^0 \left[\frac{-e}{\sqrt{2} s_w c_w} \left(\frac{1}{3} N_{1j} s_w + N_{2j} c_w \right) P_L - y_u N_{4j}^* P_R \right] u \\
& + \tilde{d}_L^* \tilde{\chi}_j^0 \left[\frac{-e}{\sqrt{2} s_w c_w} \left(\frac{1}{3} N_{1j} s_w - N_{2j} c_w \right) P_L + y_d N_{3j}^* P_R \right] d \\
& + \tilde{u}_R^* \tilde{\chi}_j^0 \left[\frac{2\sqrt{2}e}{3c_w} N_{1j}^* P_R - y_u N_{4j} P_L \right] u + \tilde{d}_R^* \tilde{\chi}_j^0 \left[\frac{-\sqrt{2}e}{3c_w} N_{1j}^* P_R + y_d N_{3j} P_L \right] d \\
& + \frac{e}{s_w c_w} Z_\mu \tilde{\chi}_i^0 \gamma^\mu (\mathcal{O}_{ij}^L P_L + \mathcal{O}_{ij}^R P_R) \tilde{\chi}_j^0 + h_u \tilde{\chi}_i^0 \left(\frac{\lambda}{\sqrt{2}} \Pi_{ij}^{45} - \frac{g_1}{2} \Pi_{ij}^{13} + \frac{g_2}{2} \Pi_{ij}^{23} \right) \tilde{\chi}_j^0 + \cdot \quad (2.3)
\end{aligned}$$



$$N_{ij} \ (i, j = 1, \dots, 5), \mathcal{O}_{ij}^L = -\mathcal{O}_{ij}^{R*} = -\tfrac{1}{2} N_{i3} N_{j3}^* + \tfrac{1}{2} N_{i4} N_{j4}^*,$$

CheckMATE

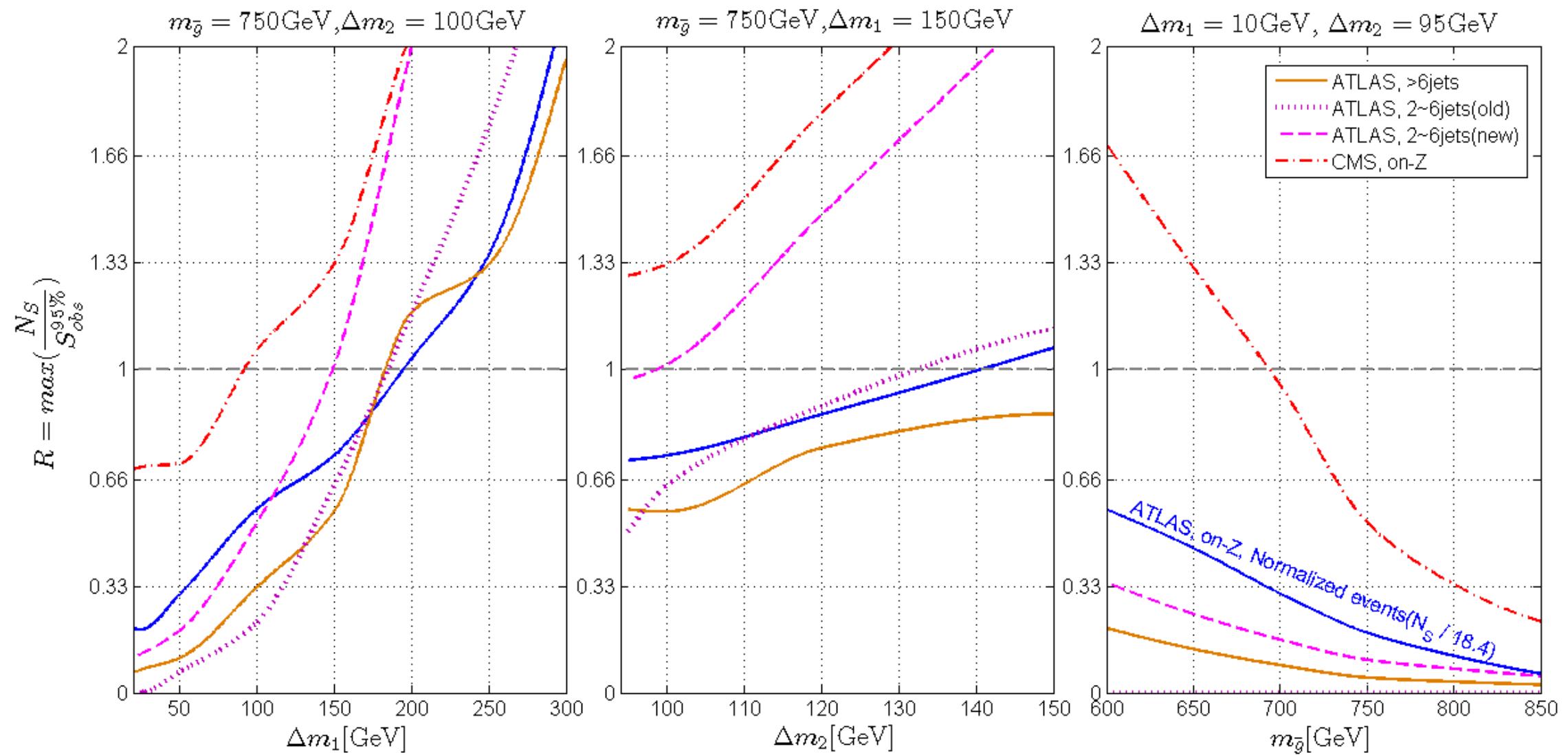
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45 "jets_btagging_n": 1,
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47 "jets_ptmin": "20.0",
48 "jets_second": "n",
49 "jets_tautagging": "n",
50 "long_info": [
51   "ATLAS",
52   "ATLAS_1503_03290",
53   "2 leptons + jets + missingET"
54   "sqrt(s) = 8 TeV",
55   "int(L) = 20.3 fb^-1"
56 ],
57 "lumi": "20.3",
58 "muon_iso_absorrel": [
59   "r",
60   "r",
61   "a"
62 ],
63 "muon_iso_dR": [
64   "0.05",
65   "0.3",
66   "0.2"
67 ],
68 "muon_iso_ptmin": [
69   "0.5",
45 void Atlas_1503_03290::analyze() {
46   missingET->addMuons(muonsCombined);
47   countCutflowEvent("CR00_noCuts");
48
49   histmissET->Fill(missingET->P4().Et());
50
51   //for baseline particles and removing overlap between them
52   electronsMedium = filterPhaseSpace(electronsMedium, 10., -2.47, 2.47, true);
53   electronsTight = filterPhaseSpace(electronsTight, 10., -2.47, 2.47, true);
54   muonsCombined = filterPhaseSpace(muonsCombined, 10., -2.4, 2.4);
55   jets = filterPhaseSpace(jets, 20., -2.5, 2.5);
56
57   electronsMedium = overlapRemoval(electronsMedium, 0.05);
58   electronsTight = overlapRemoval(electronsTight, 0.05);
59   jets = overlapRemoval(jets, electronsMedium, 0.2);
60   electronsMedium = overlapRemoval(electronsMedium, jets, 0.4);
61   electronsTight = overlapRemoval(electronsTight, jets, 0.4);
62   muonsCombined = overlapRemoval(muonsCombined, jets, 0.4);
63   electronsMedium = overlapRemoval(electronsMedium, muonsCombined, 0.01);
64   electronsTight = overlapRemoval(electronsTight, muonsCombined, 0.01);
```

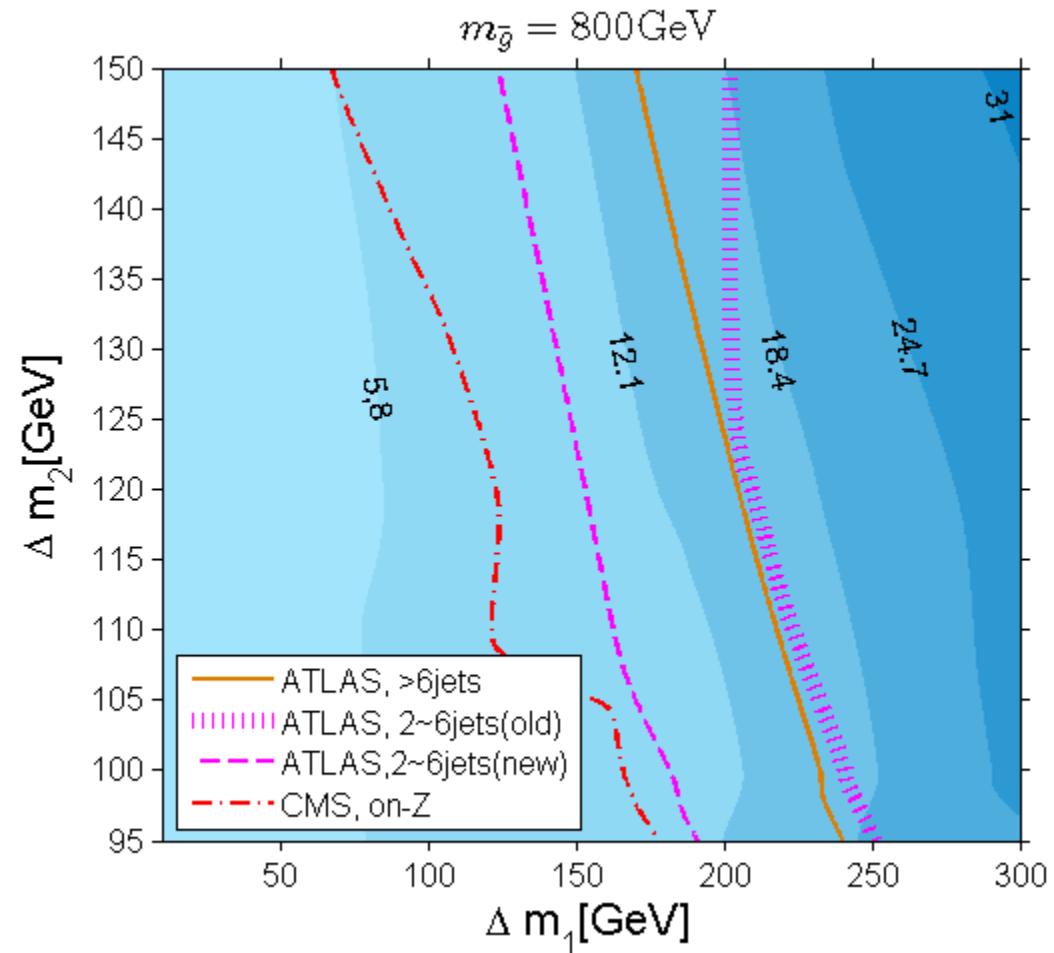
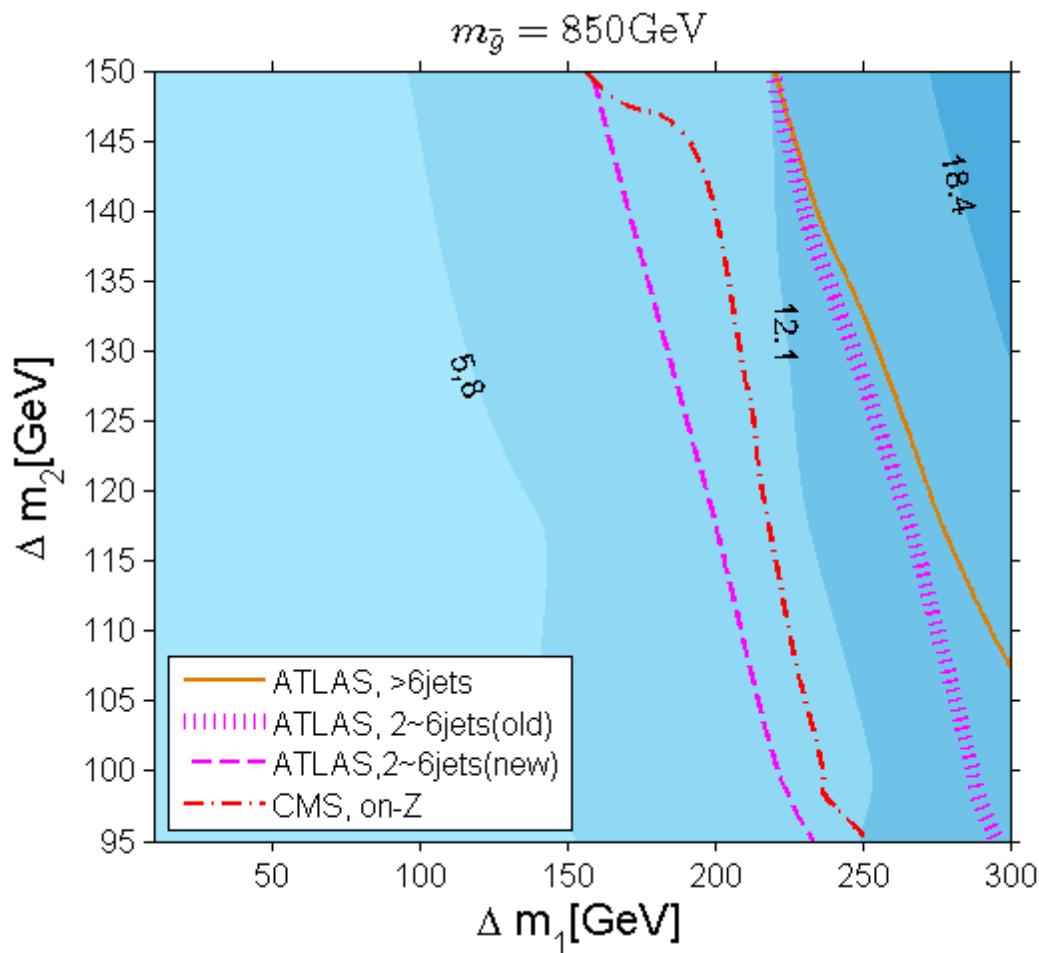
Validation

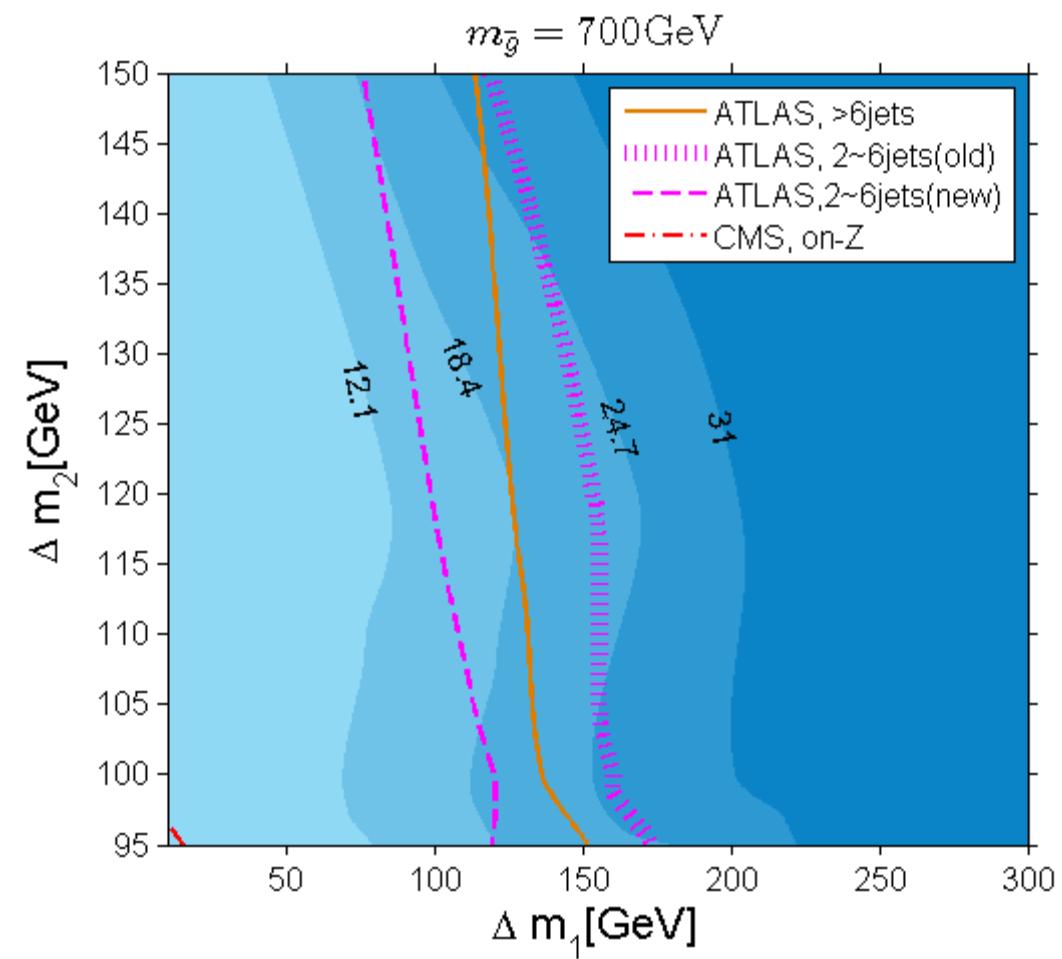
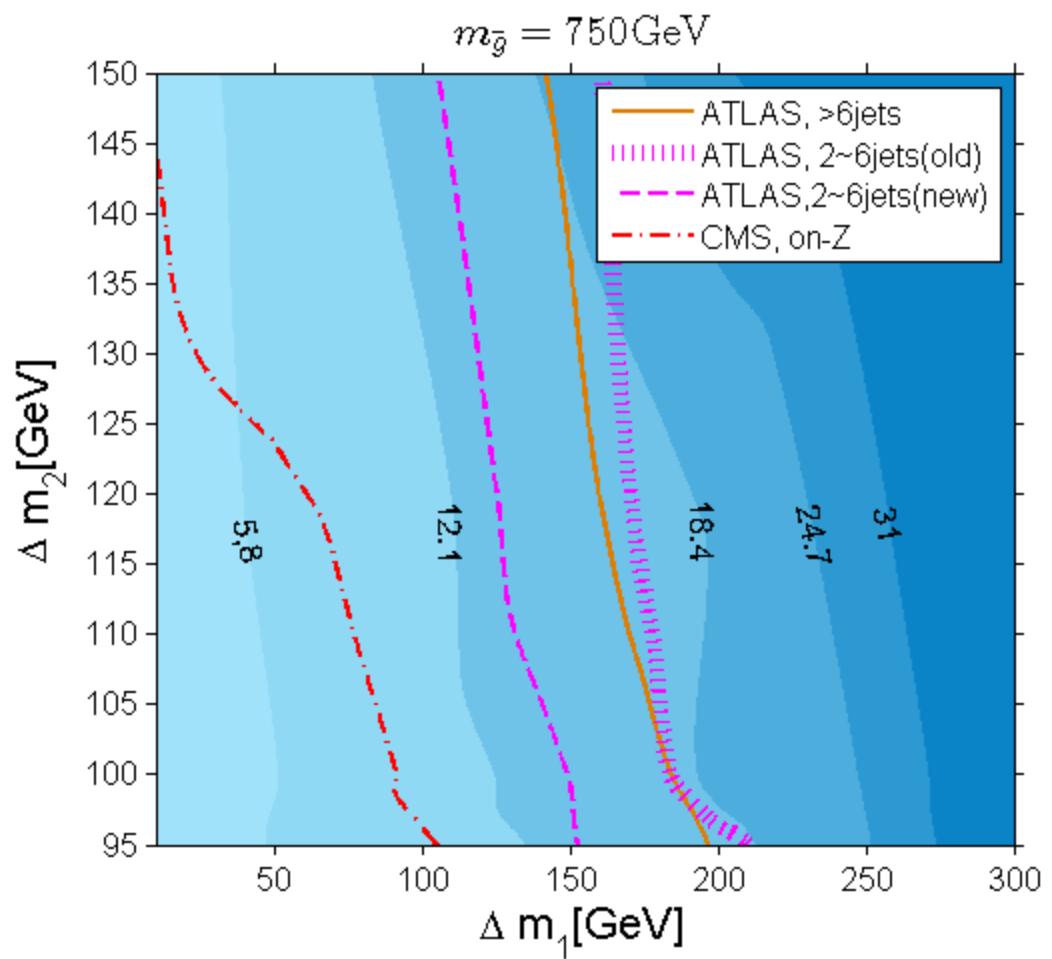
$\tilde{g} \tilde{g}$ one step, $m_{\tilde{g}} = 1200 GeV$, $m_{\tilde{\chi}_1^\pm} = 1150 GeV$, $m_{\tilde{\chi}_1^0} = 60 GeV$			
SR:2jW	EXP	OUR	DIFF
$E_T^{miss} > 160 GeV$, $P_T(j_{1,2}) > 130(60) GeV$	52.70	56.35	-7%
$\Delta\phi(j_{1,2,3}, E_T^{miss}) > 0.4$	46.30	49.01	-6%
$N(W)$ unresolved ≥ 2	9.20	8.70	5%
$E_T^{miss}/m_{eff}(N_j) > 0.25$	7.00	6.69	4%
$m_{eff}(\text{incl.}) > 1800 GeV$	5.30	4.86	8%
$\tilde{g} \tilde{g}$ one step, $m_{\tilde{g}} = 1265 GeV$, $m_{\tilde{\chi}_1^\pm} = 945 GeV$, $m_{\tilde{\chi}_1^0} = 625 GeV$			
SR:6jt	EXP	OUR	DIFF
$E_T^{miss} > 160 GeV$, $P_T(j_1, j_2) > 130(60) GeV$	53.30	54.16	-2%
$P_T(j_3) > 60 GeV$	53.00	53.83	-2%
$P_T(j_4) > 60 GeV$	50.50	51.50	-2%
$P_T(j_5) > 60 GeV$	41.40	43.49	-5%
$P_T(j_6) > 60 GeV$	26.70	29.78	-12%
$\Delta\phi(j_{1,2,3}, E_T^{miss}) > 0.4$	22.40	25.38	-13%
$\Delta\phi(j_{i>3}, E_T^{miss}) > 0.2$	18.20	20.54	-13%
$E_T^{miss}/m_{eff}(N_j) > 0.25$	10.90	11.64	-7%
$m_{eff}(\text{incl.}) > 1500 GeV$	4.20	4.81	-15%

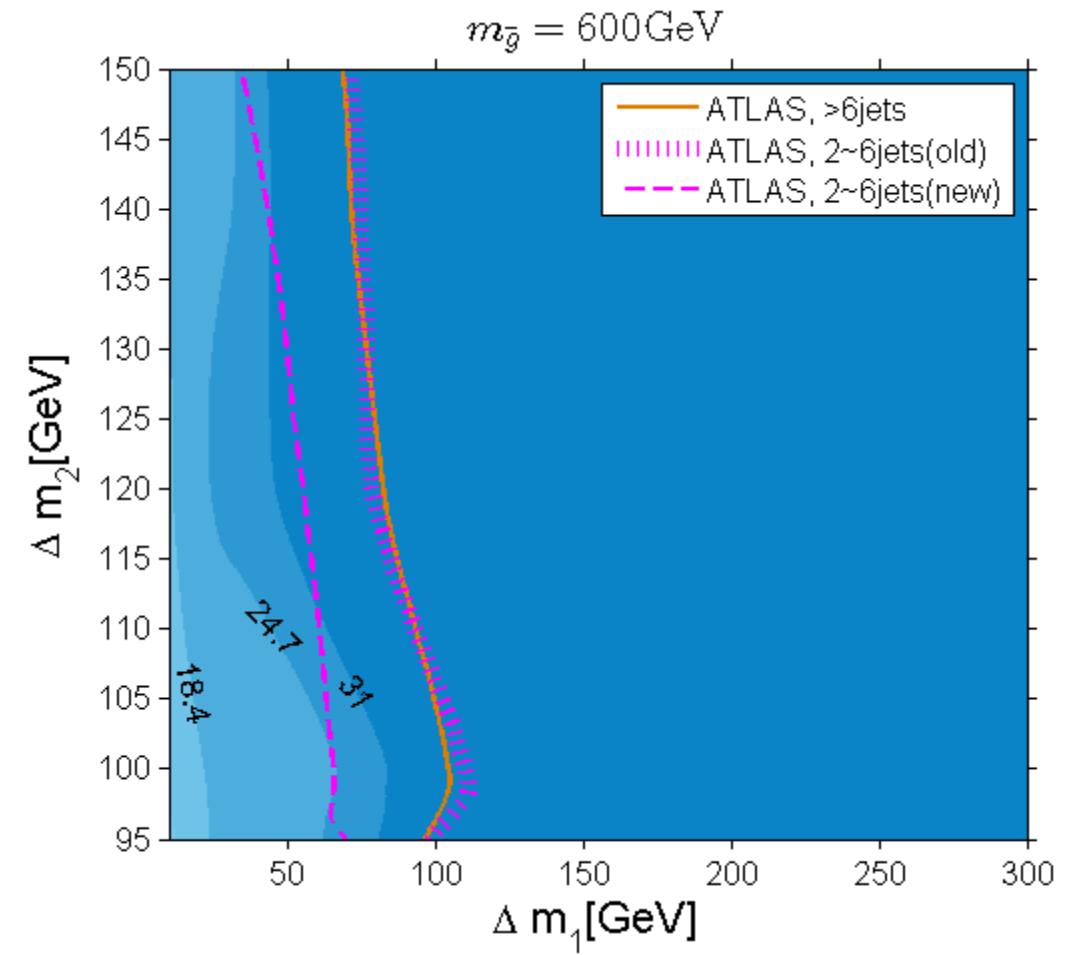
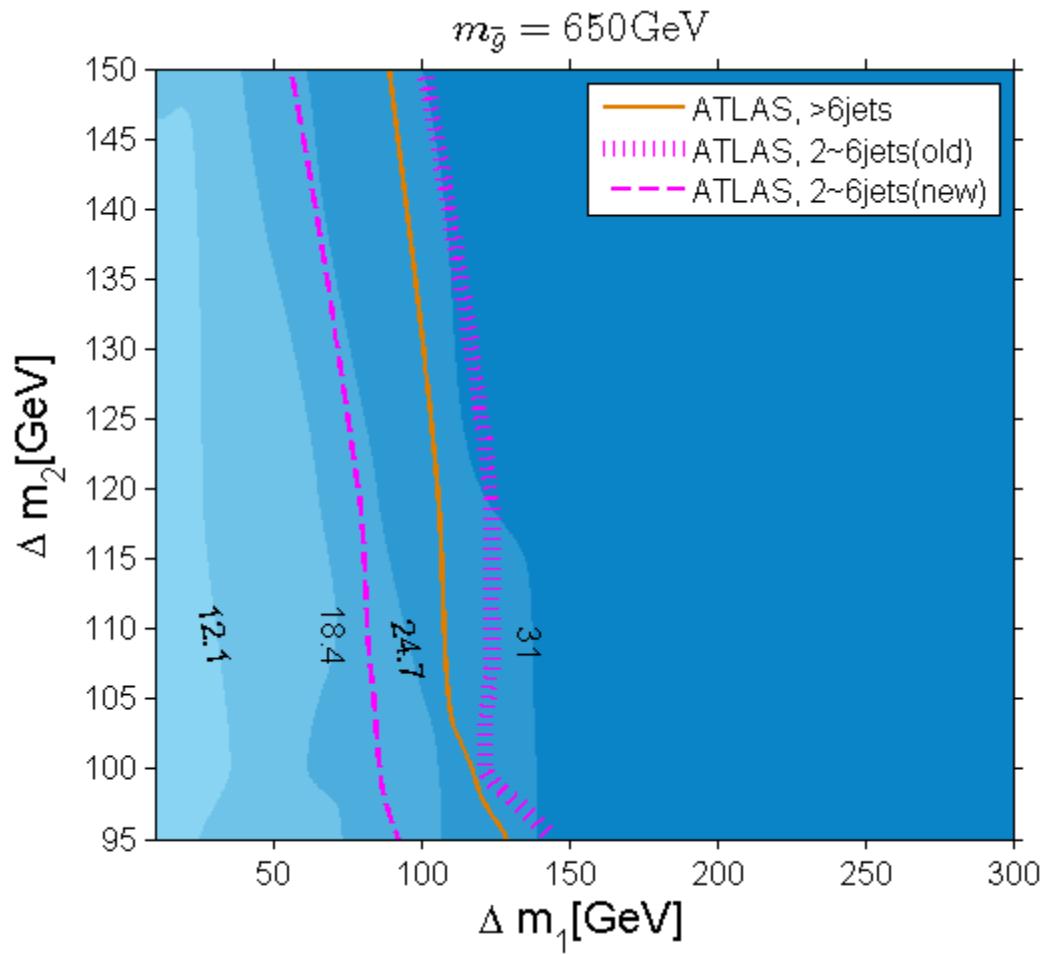
Validation

	$m_{\tilde{g}} = 900GeV, \mu = 600GeV$					
	EXP	OUR	DIFF	EXP	OUR	DIFF
No cuts	189	189				
At least 2 leptons	88.8	73.41	-17%			
	ee			$\mu\mu$		
	EXP	OUR	DIFF	EXP	OUR	DIFF
Lepton flavour	36.1	30.77	-15%	25.7	28.5	11%
PromptLeptons	35.3	30.73	-13%	25.6	28.43	11%
Opposite charged leptons	33.6	30.05	-11%	24.2	28.01	16%
> 1jet	32.2	27.75	-14%	23.1	25.97	12%
$m_{ll} > 15$	30.0	27.69	-8%	23.0	25.93	13%
$\Delta\phi(j_1, E_T^{miss}) > 0.4$	28.3	25.91	-8%	21.9	24.48	12%
$\Delta\phi(j_2, E_T^{miss}) > 0.4$	25.7	23.27	-9%	19.9	22.13	11%
$81GeV < m_{ll} < 101GeV$	22.1	21.38	-3%	16.6	18.77	13%
$H_T > 600GeV$	20.5	18.01	-12%	15.1	15.54	3%
$E_T^{miss} > 225GeV$	15.0	13.74	-8%	11.1	11.47	3%







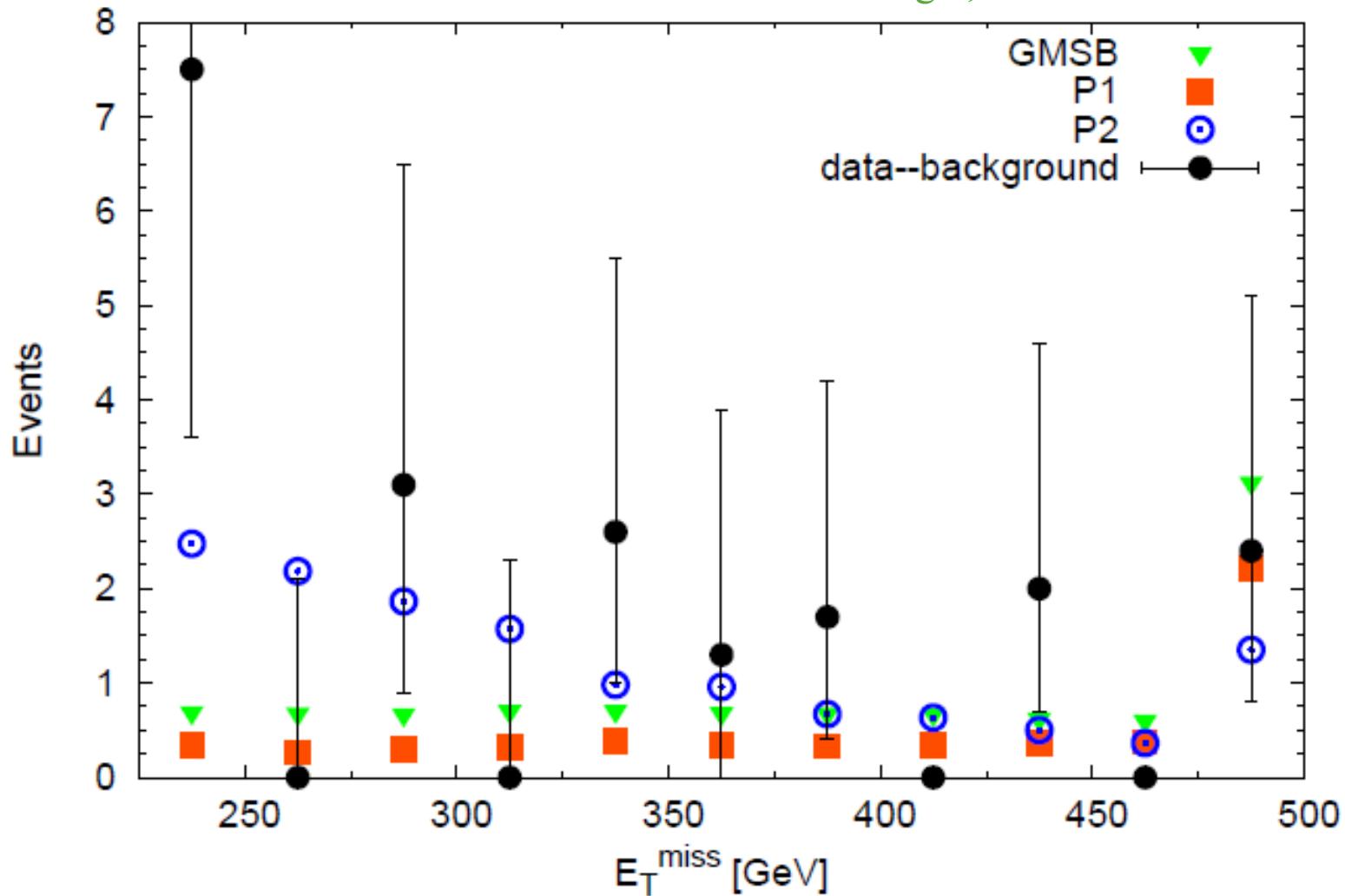


Result

1. With out CMS limit, NMSSM can explain the ATLAS Z-peaked excess at 1σ within in optimal cases, such as $m_{\tilde{g}} = 650GeV$, $m_{\tilde{\chi}_2^0} = 565GeV$, $m_{\tilde{\chi}_1^0} = 465GeV$.
2. With CMS limit, NMSSM can explain the ATLAS Z-peaked excess at 1.2σ away(11 Events).

Distribution

Ulrich Ellwanger, arXiv:1504.02244



Conclusion

- I Monte Carlo simulation
 - 1. Why do we need to do MC simulation?
 - An important way in phenomenology
 - 2. How to do MC simulation in collider physics?
 - MadGraph/MadEvent, Pythia, Delphes, Root, CheckMATE, ...
 - 3. Which kind of work can MC simulation do?
 - Constant, Prediction, Improve, Explanation
- II Example: Explanation of the ATLAS Z-peaked excess in the NMSSM
 - 1. Brief introduction of SUSY
 - MSSM and NMSSM
 - 2. The ATLAS experimental report for Z-peaked excess
 - Process, identification, selection, results
 - 3. The explanation of the Z-peaked excess
 - $m_{\tilde{g}} = 650\text{GeV}$, $m_{\tilde{\chi}_2^0} = 565\text{GeV}$, $m_{\tilde{\chi}_1^0} = 465\text{GeV}$ can explain at 1σ

WE ARE JUST C
Thank you.

THE WAY