

# Higgs mass prediction in BSM models with FlexibleSUSY using an automatized EFT approach

Tom Steudtner, Dominik Stöckinger, Alexander Voigt

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# Contents

## ① What is FlexibleSUSY?

## ② Approaches to predict $M_h$

Full model approach

EFT approach

Automatic EFT approach

Equivalence of EFT and automatic EFT approach

## ③ Numerical comparison of the approaches

## ④ Summary

# Contents

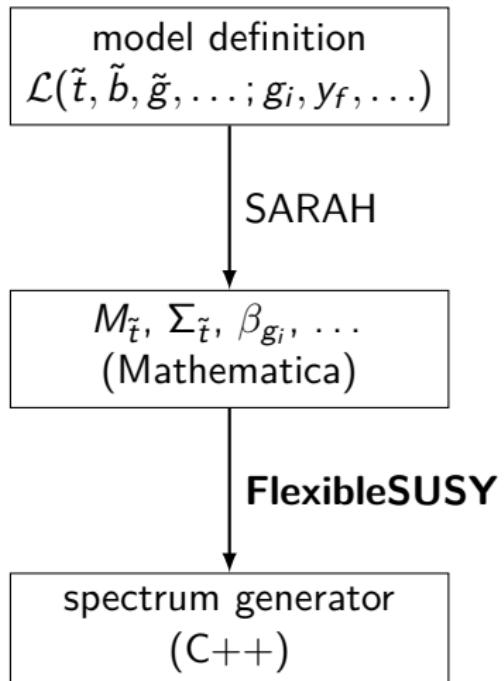
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  - Full model approach
  - EFT approach
  - Automatic EFT approach
  - Equivalence of EFT and automatic EFT approach
- ③ Numerical comparison of the approaches
- ④ Summary

# FlexibleSUSY = spectrum generator generator

FlexibleSUSY



# Generating a spectrum generator



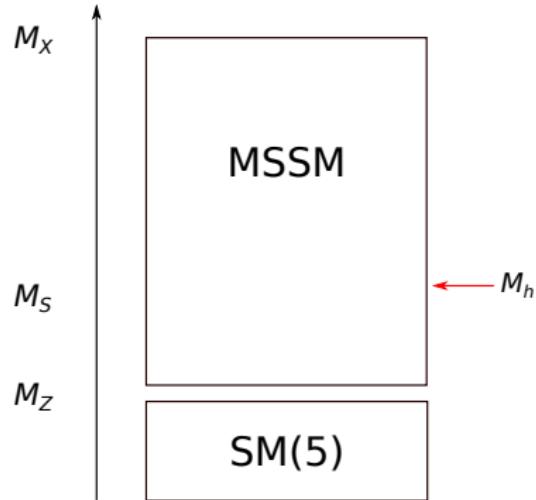
# Available models with MSSM high-scale origin

Model	RGEs	$h$ self-energy contributions	matching conditions to the MSSM
<b>MSSM ("full model")</b>	3L	1L + 2L $O((\alpha_t + \alpha_b)\alpha_s)$ + 2L $O((\alpha_t + \alpha_b)^2)$	-
THDM	2L	1L	1L $\lambda_i O((\alpha_t + \alpha_b + \alpha_\tau)\alpha_i)$ + 2L $\lambda_i O(\alpha_t^2\alpha_s)$ [1508.00576]
THDM + $\tilde{h}_i$	2L	1L	1L $\lambda_i O((\alpha_t + \alpha_b + \alpha_\tau)\alpha_i)$ + 2L $\lambda_i O(\alpha_t^2\alpha_s)$ [1508.00576]
THDM + split	2L	1L	1L $\lambda_i O((\alpha_t + \alpha_b + \alpha_\tau)\alpha_i)$ + 2L $\lambda_i O(\alpha_t^2\alpha_s)$ [1508.00576]
SM + split	2L	1L + 2L $O(\alpha_t(\alpha_s + \alpha_t))$ + 3L gluino $O(\alpha_t\alpha_s^2)$	1L $\tilde{g}_{ij} O(\alpha_t + \alpha_i)$ + 1L $\lambda O((\alpha_t + \alpha_i)^2)$ + 2L $\lambda O(\alpha_s\alpha_t^2)$ [1407.4081]
<b>SM ("EFT")</b>	3L	1L + 2L $O(\alpha_t(\alpha_s + \alpha_t))$	1L $\lambda O((\alpha_t + \alpha_i)^2 + \alpha_b^2 + \alpha_\tau^2)$ + 2L $\lambda O((\alpha_s + \alpha_t)\alpha_t^2)$ [1407.4081, 1504.05200]
<b>SM ("automatic EFT")</b>	3L	1L	1L $\lambda + O(p^2/M_S^2)$ terms

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- ① What is FlexibleSUSY?
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  - EFT approach
  - Automatic EFT approach
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## Full model approach



**Idea:** Calculate  $M_h$  in the MSSM as a function of the  $\overline{\text{DR}}$  parameters:

$$g_i, y_{ij}^f, v_i, \mu, B\mu, m_{H_i}^2, m_{\tilde{f},ij}^2, M_i, T_{ij}^f$$

# Full model approach

Fixed by observables:

Input		Output
$\alpha_{\text{em}}^{\text{SM}(5)}(M_Z)$	$\rightarrow \alpha_{\text{em}}^{\text{MSSM}}(M_Z)$	$\rightarrow g_1^{\text{MSSM}}(M_Z)$
$G_F$	$\rightarrow \sin \theta_W^{\text{MSSM}}(M_Z)$	$\rightarrow g_2^{\text{MSSM}}(M_Z)$
$\alpha_s^{\text{SM}(5)}(M_Z)$		$\rightarrow g_3^{\text{MSSM}}(M_Z)$
$M_Z$	$\rightarrow m_Z^{\text{MSSM}}(M_Z)$	$\rightarrow v^{\text{MSSM}}(M_Z)$
$M_t$	$\rightarrow m_t^{\text{MSSM}}(M_Z)$	$\rightarrow y_t^{\text{MSSM}}(M_Z)$
$m_b^{\text{SM}(5)}(m_b)$	$\rightarrow m_b^{\text{MSSM}}(M_Z)$	$\rightarrow y_b^{\text{MSSM}}(M_Z)$
$M_\tau$	$\rightarrow m_\tau^{\text{MSSM}}(M_Z)$	$\rightarrow y_\tau^{\text{MSSM}}(M_Z)$

Fixed by 2 EWSB conditions:  $m_{H_u}^2, m_{H_d}^2$

Free parameters:  $\tan \beta, \mu, B\mu, m_{f,ij}^2, M_i, T_{ij}^f$

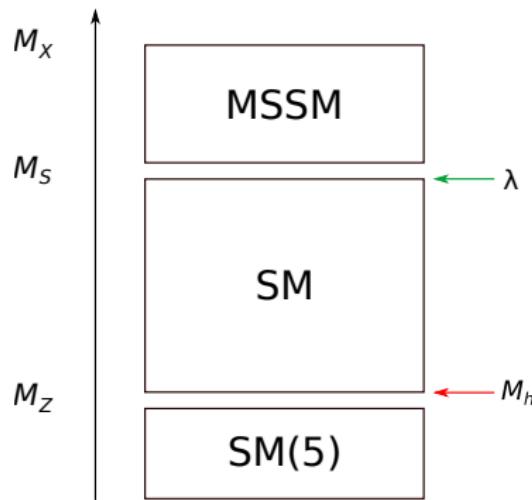
## Full model approach

$$(M_h^{\text{MSSM}})^2 = \text{smallest eigenvalue of} \\ \left[ (m_h^{\text{MSSM}})^2 - \Sigma_h^{\text{MSSM}}(p^2 = (M_h^{\text{MSSM}})^2, Q = M_S) \right]_{ij}$$

**Advantage:** includes terms  $O(p^2/M_S^2)$

**Disadvantage:** suffers from large logarithms  $\propto \log(M_Z/M_S)$  if  $M_S \gg M_Z$

## EFT approach



**Idea:** Calculate  $M_h$  in the SM as a function of the  $\overline{MS}$  parameters:

$$g_i, y_{ij}^f, v, \mu^2, \lambda$$

# EFT approach

Fixed by observables:

Input			Output
$\alpha_{\text{em}}^{\text{SM}(5)}(M_Z)$	$\rightarrow$	$\alpha_{\text{em}}^{\text{SM}}(M_Z)$	$\rightarrow g_1^{\text{SM}}(M_Z)$
$G_F$	$\rightarrow$	$\sin \theta_W^{\text{SM}}(M_Z)$	$\rightarrow g_2^{\text{SM}}(M_Z)$
$\alpha_s^{\text{SM}(5)}(M_Z)$			$\rightarrow g_3^{\text{SM}}(M_Z)$
$M_Z$	$\rightarrow$	$m_Z^{\text{SM}}(M_Z)$	$\rightarrow v^{\text{SM}}(M_Z)$
$M_t$	$\rightarrow$	$m_t^{\text{SM}}(M_Z)$	$\rightarrow y_t^{\text{SM}}(M_Z)$
$m_b^{\text{SM}(5)}(m_b)$	$\rightarrow$	$m_b^{\text{SM}}(M_Z)$	$\rightarrow y_b^{\text{SM}}(M_Z)$
$M_\tau$	$\rightarrow$	$m_\tau^{\text{SM}}(M_Z)$	$\rightarrow y_\tau^{\text{SM}}(M_Z)$

Fixed by 1 EWSB condition:  $\mu^2$

Free parameter:  $\lambda$

EFT approach: matching all  $\Gamma_{\phi_1, \dots, \phi_n}(Q = M_S)$

Determine  $\lambda(M_S)$  from matching all  $\Gamma_{\phi_1, \dots, \phi_n}(Q = M_S)$ :

$$\frac{\partial}{\partial p^2} \Gamma_{hh}^{\text{SM}}(p^2 = 0, Q = M_S) = \frac{\partial}{\partial p^2} \Gamma_{hh}^{\text{MSSM}}(p^2 = 0, Q = M_S)$$
$$\Gamma_{hhh}^{\text{SM}}(p^2 = 0, Q = M_S) = \Gamma_{hhh}^{\text{MSSM}}(p^2 = 0, Q = M_S)$$

$\Rightarrow$

$$\lambda(M_S) = \frac{1}{4} [g_Y^2 + g_2^2] \cos^2 2\beta + \Delta\lambda$$

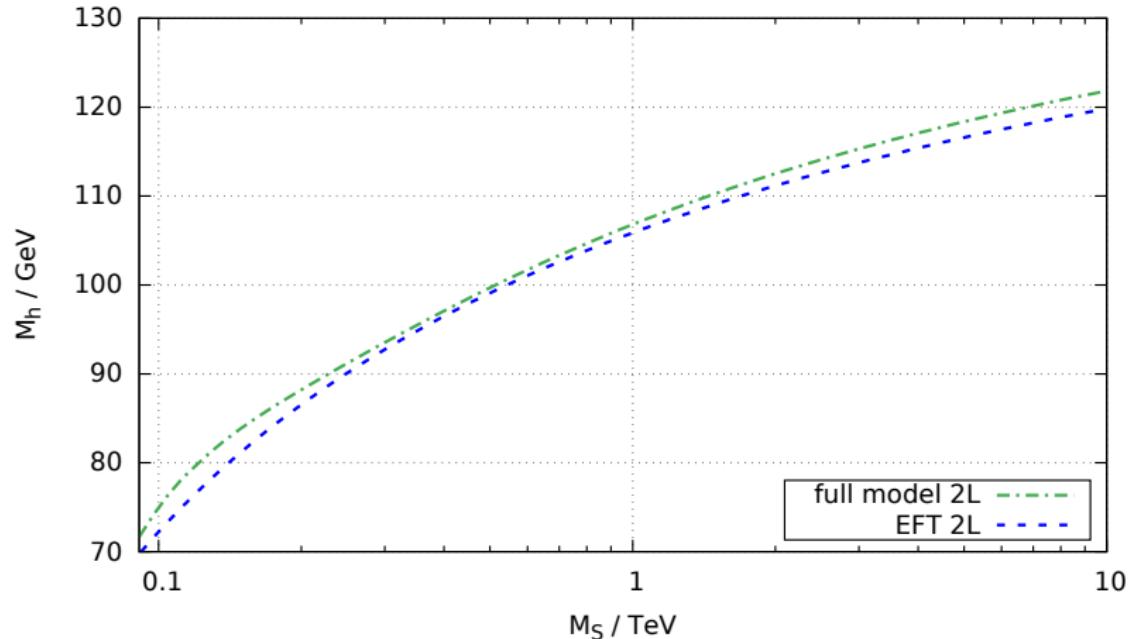
RG running to  $Q = M_Z \Rightarrow$

$$(M_h^{\text{SM}})^2 = \lambda(M_Z)v^2(M_Z) - \Sigma_h^{\text{SM}}(p^2 = (M_h^{\text{SM}})^2, Q = M_Z)$$

**Advantage:** resums large logarithms  $\propto \log(M_Z/M_S)$

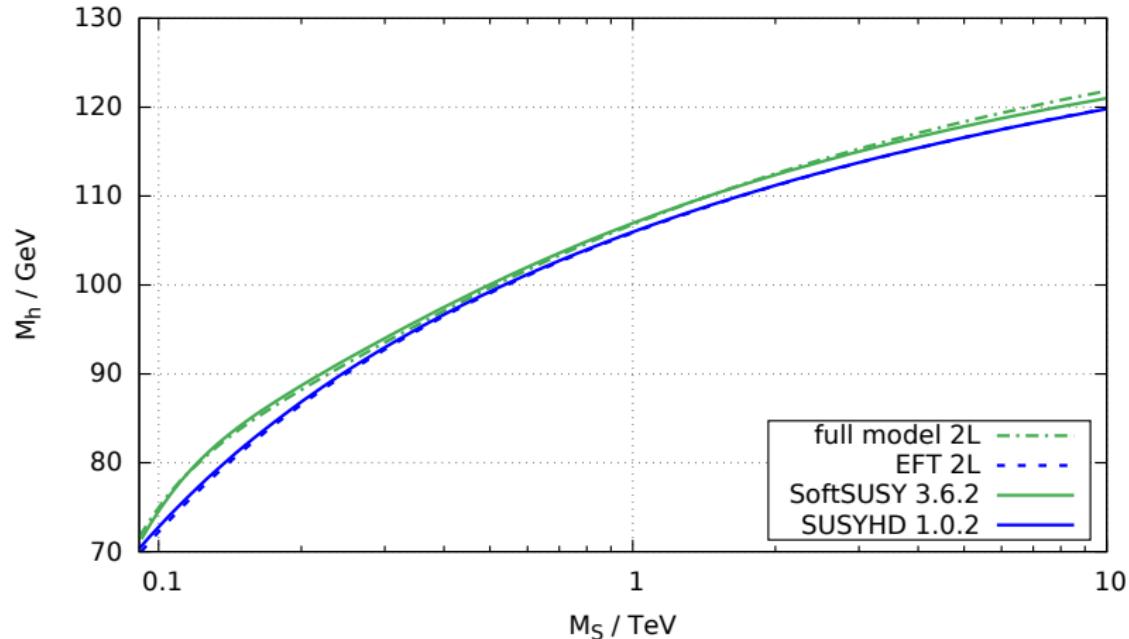
**Disadvantage:** difficult to automatize; neglects terms  $O(p^2/M_S^2)$  from SUSY particles if performed at 1L-level only

# Comparison full model vs. EFT approach



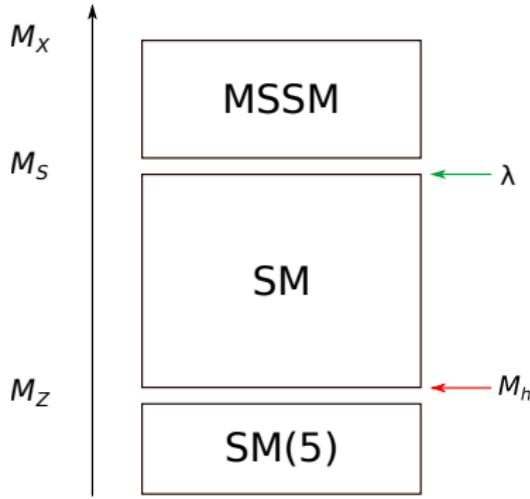
$$\tan \beta = 5, X_{t,b,\tau} = 0$$

# Comparison full model vs. EFT approach



$$\tan \beta = 5, X_{t,b,\tau} = 0$$

# Automatic EFT approach



**Idea:** Calculate  $M_h$  in the SM as a function of the  $\overline{\text{MS}}$  parameters:

$$g_i, y_{ij}^f, v, \mu^2, \lambda$$

where  $\lambda(M_S)$  is determined numerically by the 1L matching condition  $M_h^{\text{SM}} \stackrel{!}{=} M_h^{\text{MSSM}}$

## Automatic EFT approach: matching $M_h$

$$M_h^{\text{SM}} \stackrel{!}{=} M_h^{\text{MSSM}} \quad \text{at} \quad Q = M_S, 1L$$

$\Rightarrow$

$$\lambda(M_S) = \frac{1}{v^2(M_S)} \left[ (M_h^{\text{MSSM}})^2 + \Sigma_h^{\text{SM},1L}(p^2 = (M_h^{\text{SM}})^2, Q = M_S) \right]$$

where

$$(M_h^{\text{SM}})^2 = \lambda(M_S)v^2(M_S) - \Sigma_h^{\text{SM},1L}(p^2 = (M_h^{\text{SM}})^2, Q = M_S)$$

$(M_h^{\text{MSSM}})^2$  = smallest eigenvalue of

$$\left[ (m_h^{\text{MSSM}})^2 - \Sigma_h^{\text{MSSM},1L}(p^2 = (M_h^{\text{MSSM}})^2, Q = M_S) \right]_{ij}$$

## Automatic EFT approach: matching $M_h$

Next step: RG running to  $Q = M_Z$

$\Rightarrow$

$$(M_h^{\text{SM}})^2 = \lambda(M_Z)v^2(M_Z) - \Sigma_h^{\text{SM},1L}(p^2 = (M_h^{\text{SM}})^2, Q = M_Z)$$

# Equivalence EFT and automatic EFT approach $O(\hbar y_t^4)$

$$M_h^{\text{SM}} \stackrel{!}{=} M_h^{\text{MSSM}} \quad \text{at} \quad Q = M_S, 1L$$

where

$$(M_h^{\text{SM}})^2 = \lambda v^2 - \Sigma_h^{\text{SM}} + t_h^{\text{SM}}$$

$$t_h^{\text{SM}} = -6(y_t^{\text{SM}})^2 A_0(m_t) / (4\pi)^2$$

and [neglecting stop mass mixing  $O(m_t X_t / M_S^2)$ ]

$$(M_h^{\text{MSSM}})^2 = \frac{1}{4}(g_Y^2 + g_2^2)(v_u^2 + v_d^2)c_{2\beta}^2 - \Sigma_h^{\text{MSSM}} + t_h^{\text{MSSM}}$$

$$\begin{aligned} \Sigma_h^{\text{MSSM}} &= \Sigma_h^{\text{SM}} \frac{c_\alpha^2}{s_\beta^2} + 3 \frac{(y_t^{\text{SM}})^2}{(4\pi)^2} \frac{c_\alpha^2}{s_\beta^2} \left\{ A_0(m_{Q_3}) + A_0(m_{U_3}) \right. \\ &\quad \left. + 2m_t [B_0(m_{Q_3}, m_{Q_3}) + B_0(m_{U_3}, m_{U_3})] \right\} \end{aligned}$$

$$t_h^{\text{MSSM}} = -3 \frac{(y_t^{\text{SM}})^2}{(4\pi)^2} \frac{c_\alpha^2}{s_\beta^2} [2A_0(m_t) - A_0(m_{Q_3}) - A_0(m_{U_3})]$$

## Equivalence EFT and automatic EFT approach $O(\hbar y_t^4)$

in SM limit  $\frac{c_\alpha^2}{s_\beta^2} \rightarrow 1$   
 $\Rightarrow$

$$\begin{aligned}\lambda &= \frac{1}{4}(g_Y^2 + g_2^2)c_{2\beta}^2 \\ &\quad - 3\frac{(y_t^{\text{SM}})^4}{(4\pi)^2} \left[ B_0(p^2, m_{Q_3}, m_{Q_3}) + B_0(p^2, m_{U_3}, m_{U_3}) \right] \\ &= \frac{1}{4}(g_Y^2 + g_2^2)c_{2\beta}^2 \\ &\quad - 3\frac{(y_t^{\text{SM}})^4}{(4\pi)^2} \left[ -\log \frac{m_{Q_3}^2}{Q^2} + \frac{p^2}{6m_{Q_3}^2} + O\left(\frac{p^4}{m_{Q_3}^4}\right) \right. \\ &\quad \left. - \log \frac{m_{U_3}^2}{Q^2} + \frac{p^2}{6m_{U_3}^2} + O\left(\frac{p^4}{m_{U_3}^4}\right) \right] \\ &= [\text{Bagnaschi et. al. 2014}] + O\left(\frac{p^2}{m_{Q_3}^2}\right) + O\left(\frac{p^2}{m_{U_3}^2}\right)\end{aligned}$$

# Summary of automatic EFT approach

## Advantages:

- easily automatizable
- resums large logarithms  $\log(M_Z/M_S)$  due to RG running
- includes 1L terms  $O(p^2/M_S^2)$  from SUSY particles
- incorporates SUSY particle mixings

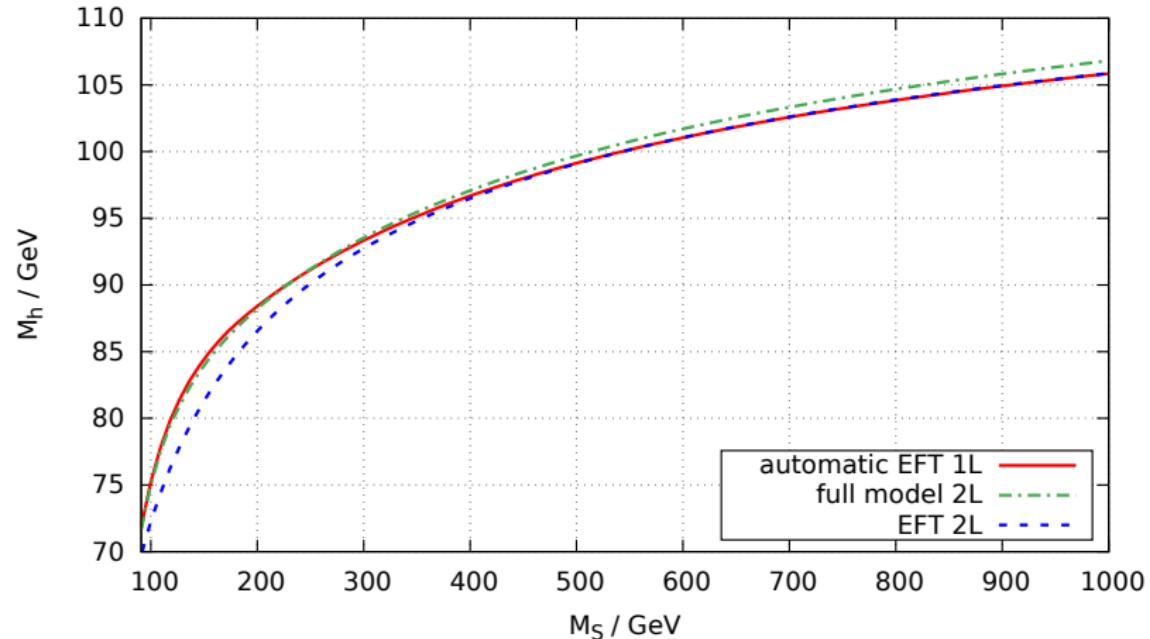
## Disadvantage:

- large logs in matching cancel only in SM limit  $\cos(\beta - \alpha) \rightarrow 0$
- only 1L calculation so far

# Contents

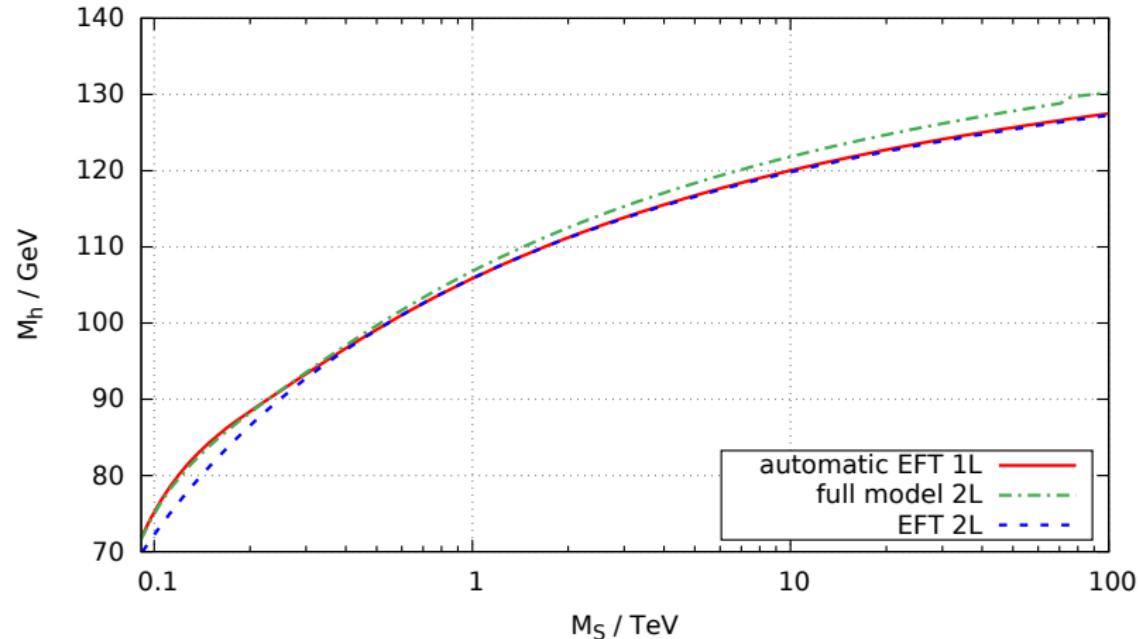
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# Comparison full vs. EFT vs. automatic EFT approach



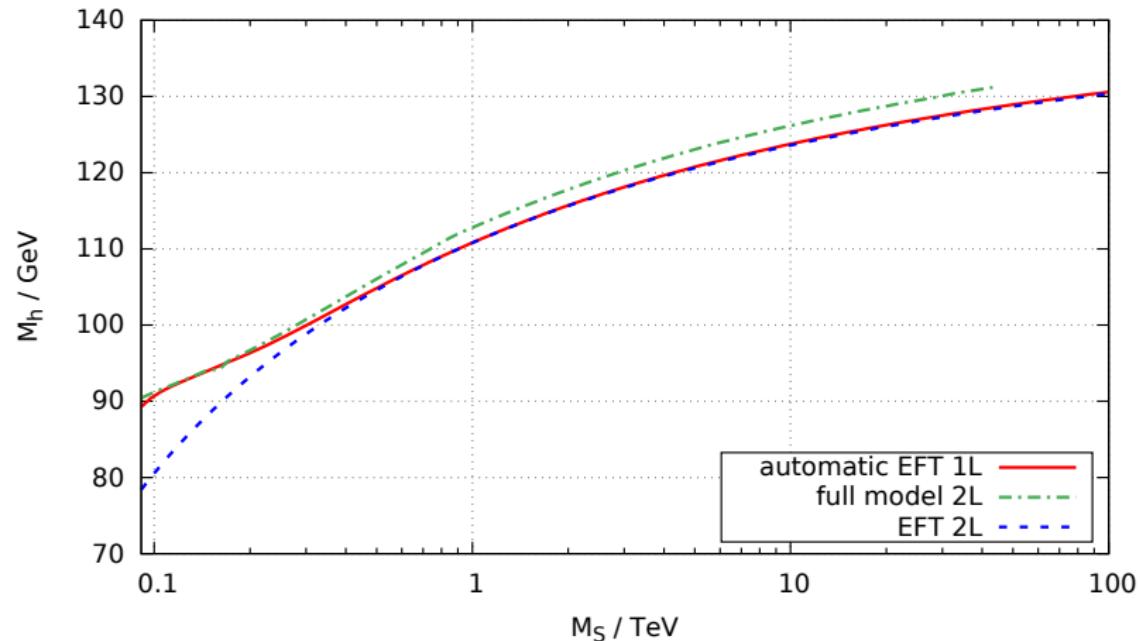
$$\tan \beta = 5, X_{t,b,\tau} = 0$$

# Comparison full vs. EFT vs. automatic EFT approach



$$\tan \beta = 5, X_{t,b,\tau} = 0$$

# Comparison full vs. EFT vs. automatic EFT approach



$$\tan \beta = 20, X_{t,b,\tau} = 0$$

# Summary

## Summary:

- full model approach suffers from large logs  $\propto \log(M_Z/M_S)$  if  $M_S \gg M_Z$
- EFT approach resums large logs, but difficult to automatize (also misses 1L terms  $O(p^2/M_S^2)$  if performed at 1L-level only)
- “automatic EFT” approach
  - can be automatized easily → incorporation into FlexibleSUSY
  - resums large logs  $\propto \log(M_Z/M_S)$
  - includes terms  $O(p^2/M_S^2)$  at 1L
  - handles SUSY particle mixings
  - large logs cancel only in SM limit

## Todo:

- perform matching and calculation of  $M_h$  at 2-loop level
- test non-SUSY and non-minimal SUSY models

# Backup

## Determination of MSSM parameters

# Determination of $g_3^{\text{MSSM}}(M_S)$

$$\alpha_s^{\text{MSSM}}(M_S) = \frac{\alpha_s^{\text{SM}}(M_S)}{1 - \Delta\alpha_s(M_S)}$$

with

$$\Delta\alpha_s(Q) = \frac{\alpha_s}{2\pi} \left[ \frac{1}{2} - \sum_{\text{SUSY particle } f} T_f \log \frac{m_f}{Q} \right]$$

$\Rightarrow$

$$g_3^{\text{MSSM}}(M_S) = \sqrt{4\pi\alpha_s^{\text{MSSM}}(M_S)}$$

## Determination of $v_i^{\text{MSSM}}(M_S)$

$$M_Z^{\text{SM}} = M_Z^{\text{MSSM}}$$

$\Rightarrow$

$$(m_Z^{\text{MSSM}}(M_S))^2 = (M_Z^{\text{SM}})^2 + \Pi_Z^{\text{MSSM},1L}(Q = M_S)$$

$$(M_Z^{\text{SM}})^2 = \frac{1}{4} \left[ (g_Y^{\text{SM}})^2 + (g_2^{\text{SM}})^2 \right] (v^{\text{SM}})^2 - \Pi_Z^{\text{SM},1L}(Q = M_S)$$

$\Rightarrow$

$$v^{\text{MSSM}}(M_S) = \frac{2m_Z^{\text{MSSM}}(M_S)}{\sqrt{(g_Y^{\text{MSSM}})^2 + (g_2^{\text{MSSM}})^2}}$$

$\Rightarrow$

$$v_u^{\text{MSSM}}(M_S) = v^{\text{MSSM}}(M_S) \sin \beta(M_S)$$

$$v_d^{\text{MSSM}}(M_S) = v^{\text{MSSM}}(M_S) \cos \beta(M_S)$$

# Determination of $y_i^{\text{MSSM}}(M_S)$

$$M_f^{\text{SM}} = M_f^{\text{MSSM}}$$

$\Rightarrow$

$$m_f^{\text{MSSM}}(M_S) = M_f^{\text{SM}} + \Sigma_f^{\text{MSSM},1L}(Q = M_S)$$

$$M_f^{\text{SM}} = \frac{\sqrt{2}m_f^{\text{SM}}}{v_i^{\text{SM}}} - \Sigma_f^{\text{SM},1L}(Q = M_S)$$

$\Rightarrow$

$$y_f^{\text{MSSM}}(M_S) = \frac{\sqrt{2}m_f^{\text{MSSM}}(M_S)}{v_i^{\text{MSSM}}(M_S)}$$

## Determination of SM parameters

# Determination of $g_3^{\text{SM}}(M_Z)$

**Input:**  $\alpha_s^{\text{SM}(5)}(M_Z) = 0.1185$

$\rightarrow$

$$\alpha_s^{\text{SM}}(M_Z) = \frac{\alpha_s^{\text{SM}(5)}(M_Z)}{1 - \Delta\alpha_s(M_Z)}$$

with

$$\Delta\alpha_s(Q) = \frac{\alpha_s}{2\pi} \left[ -\frac{2}{3} \log \frac{m_t}{Q} \right]$$

$\Rightarrow$

$$g_3^{\text{SM}}(M_Z) = \sqrt{4\pi\alpha_s^{\text{SM}}(M_Z)}$$

# Determination of $y_t^{\text{SM}}(M_Z)$

$$y_t^{\text{SM}}(M_Z) = \frac{\sqrt{2} m_t^{\text{SM}}(M_Z)}{v(M_Z)}$$

where

$$\begin{aligned} m_t^{\text{SM}}(Q) &= M_t + \text{Re } \Sigma_t^S(M_Z) + M_t \left[ \text{Re } \Sigma_t^L(M_Z) \right. \\ &\quad \left. + \text{Re } \Sigma_t^R(M_Z) + \Delta m_t^{1L, \text{gluon}} + \Delta m_t^{2L, \text{gluon}} \right] \end{aligned}$$

$$\Delta m_t^{1L, \text{gluon}} = -\frac{g_3^2}{12\pi^2} \left[ 4 - 3 \log \left( \frac{m_t^2}{Q^2} \right) \right]$$

$$\begin{aligned} \Delta m_t^{2L, \text{gluon}} &= \left( \Delta m_t^{1L, \text{gluon}} \right)^2 \\ &\quad - \frac{g_3^4}{4608\pi^4} \left[ 396 \log^2 \left( \frac{m_t^2}{Q^2} \right) - 2028 \log \left( \frac{m_t^2}{Q^2} \right) \right. \\ &\quad \left. - 48\zeta(3) + 2821 + 16\pi^2(1 + \log 4) \right] \end{aligned}$$

## Determination of $v^{\text{SM}}$

The VEV  $v^{\text{SM}}$  is calculated from the running  $Z$  mass at  $Q = M_Z$ :

$$v^{\text{SM}}(M_Z) = \frac{2m_Z^{\text{SM}}(M_Z)}{\sqrt{g_Y^2 + g_2^2}}$$

$$m_Z^{\text{SM}}(M_Z) = \sqrt{M_Z^2 + \Pi_Z^{1L}(p^2 = M_Z^2, Q = M_Z)}$$

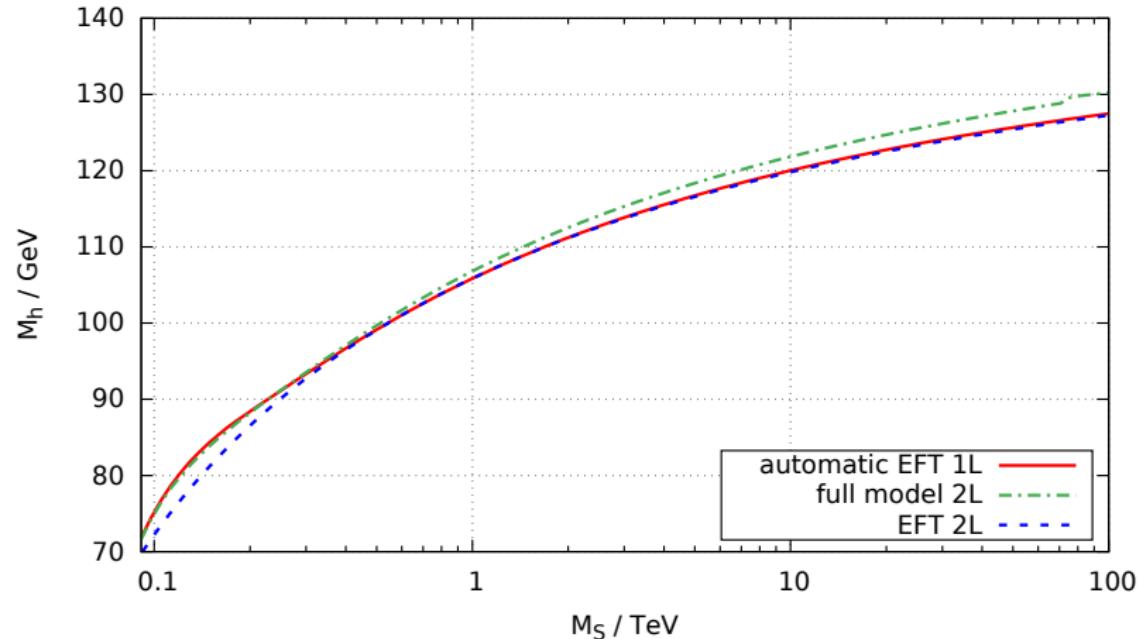
$v^{\text{SM}}$  evolves under RG running according to  
[Sperling, Stöckinger, AV, 2013, 2014]

$X_t$  dependence

(default algorithm:  $1L$  matching SM  $\rightarrow$  MSSM)

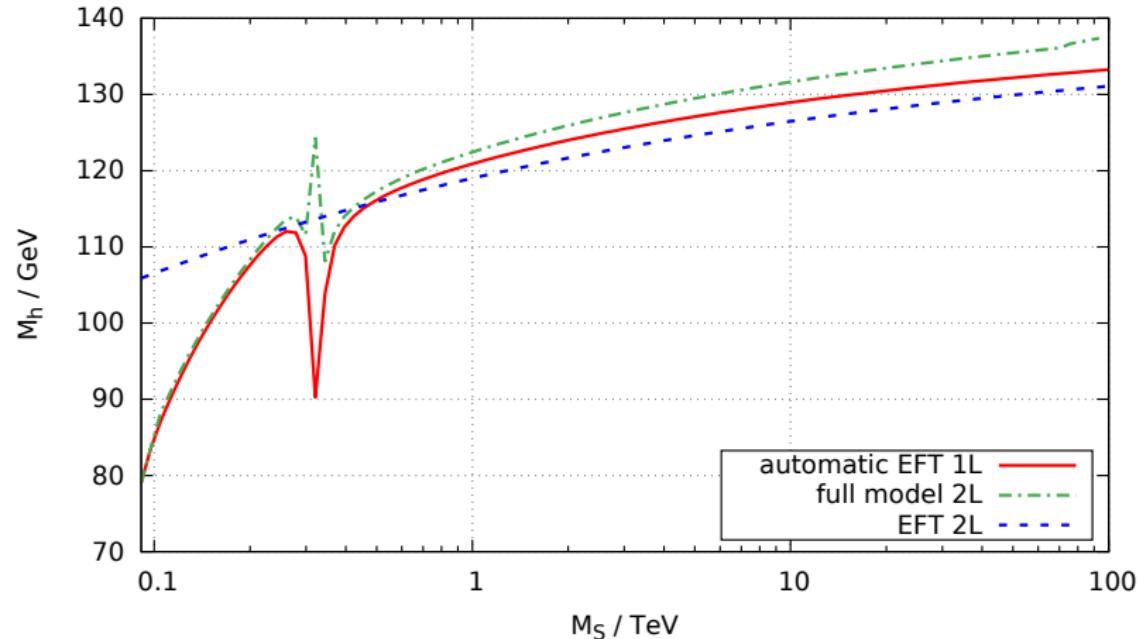
$$\tan \beta = 5$$

# Comparison full vs. EFT vs. automatic EFT approach



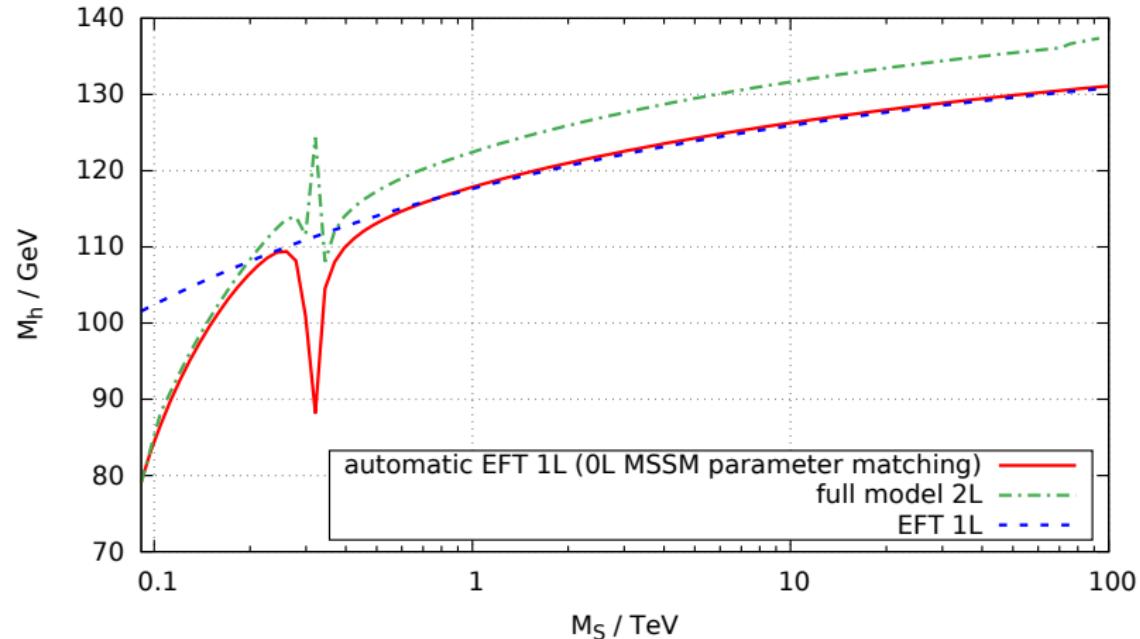
$$\tan \beta = 5, X_{t,b,\tau} = 0$$

# Comparison full vs. EFT vs. automatic EFT approach



$$\tan \beta = 5, X_t = \sqrt{6}, X_{b,\tau} = 0$$

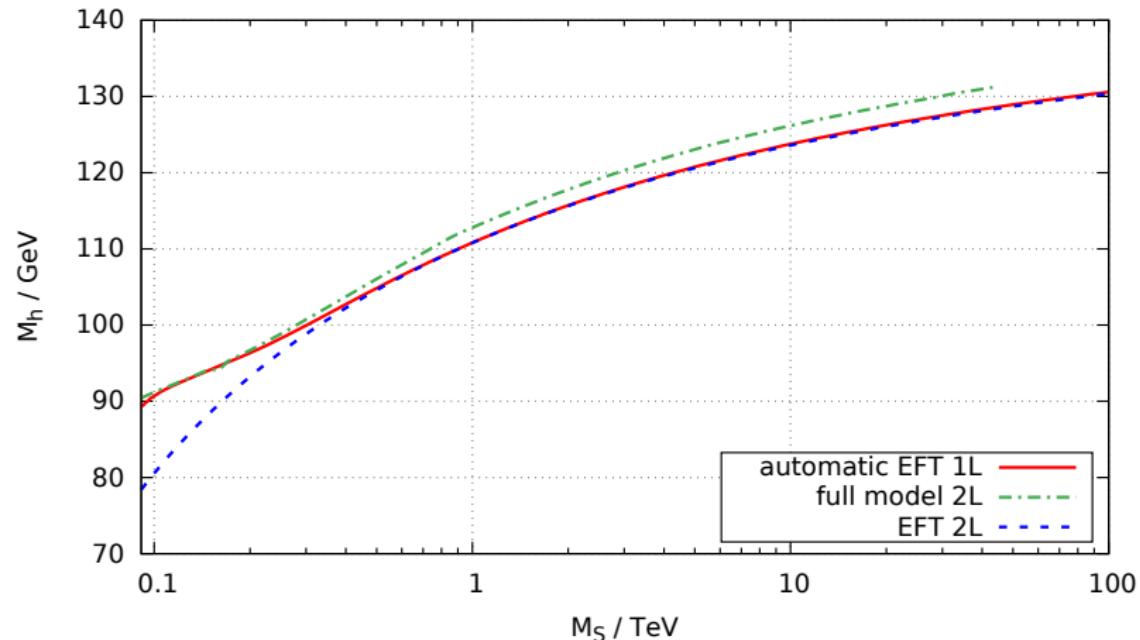
# Comparison full vs. EFT vs. automatic EFT approach



$\tan \beta = 5, X_t = \sqrt{6}, X_{b,\tau} = 0$  (0L matching SM  $\rightarrow$  MSSM)

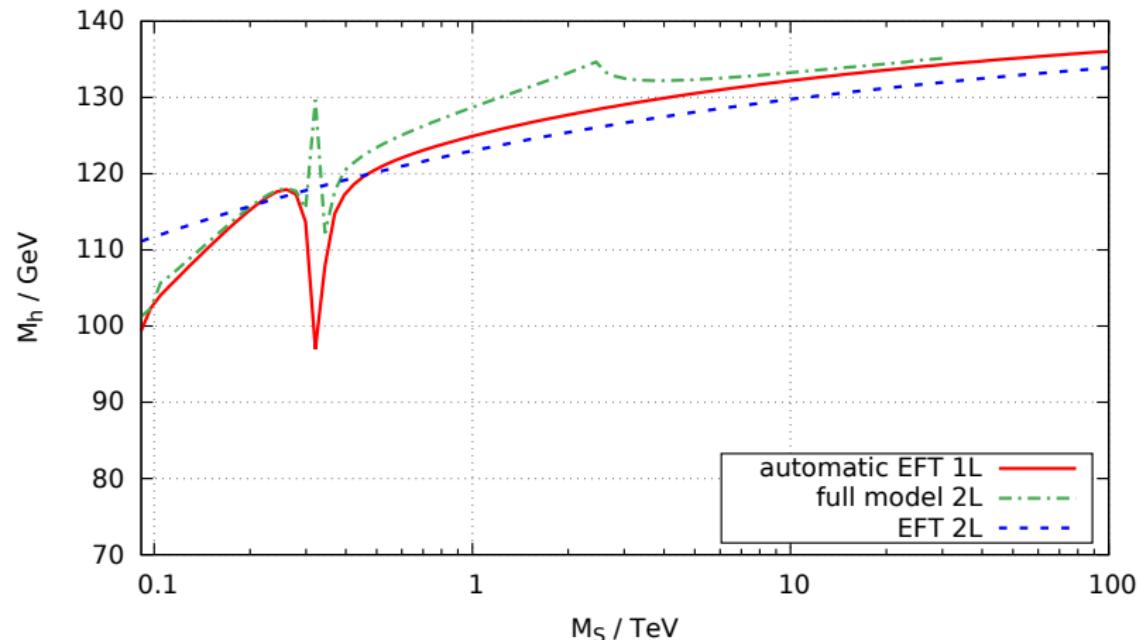
$$\tan \beta = 20$$

# Comparison full vs. EFT vs. automatic EFT approach



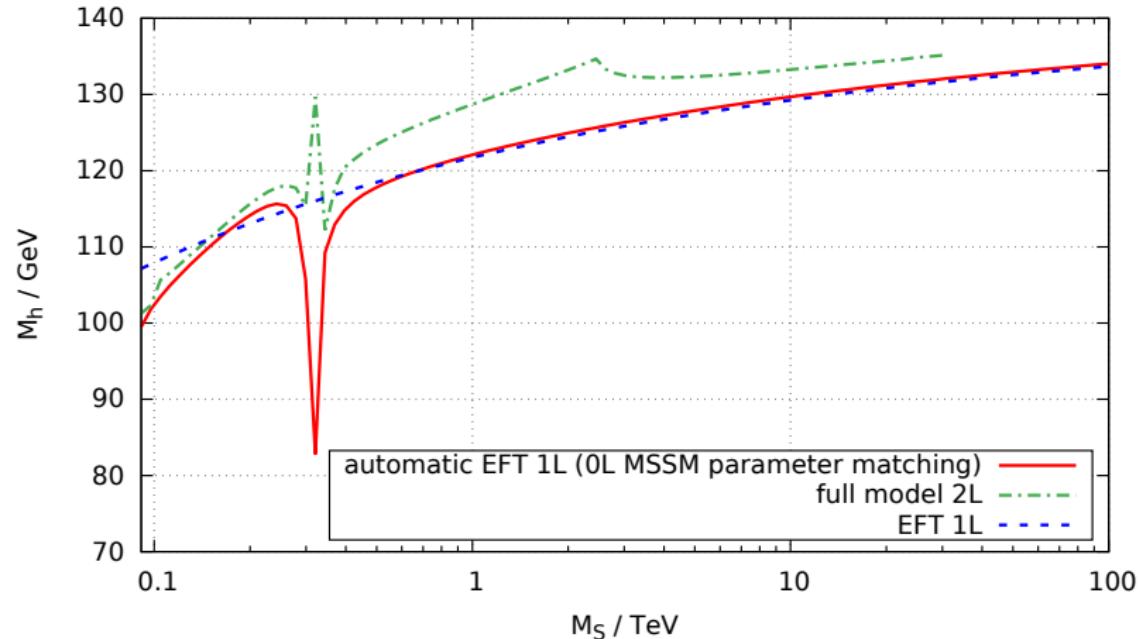
$$\tan \beta = 20, X_{t,b,\tau} = 0$$

# Comparison full vs. EFT vs. automatic EFT approach



$$\tan \beta = 20, X_t = \sqrt{6}, X_{b,\tau} = 0$$

# Comparison full vs. EFT vs. automatic EFT approach

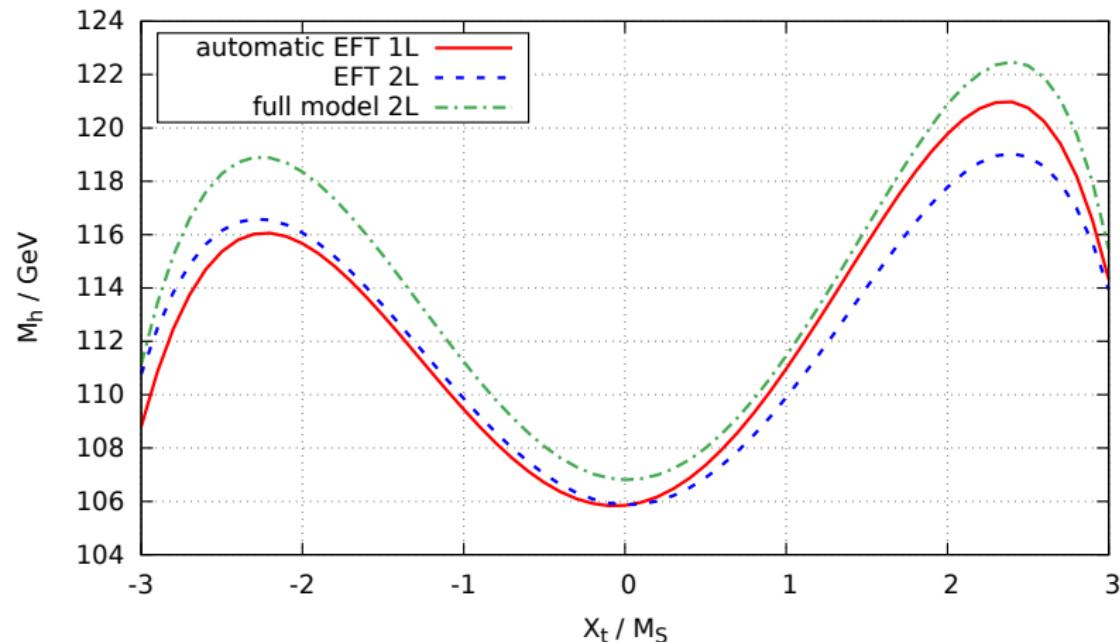


$\tan \beta = 20, X_t = \sqrt{6}, X_{b,\tau} = 0$  (0L matching SM  $\rightarrow$  MSSM)

$X_t$  dependence

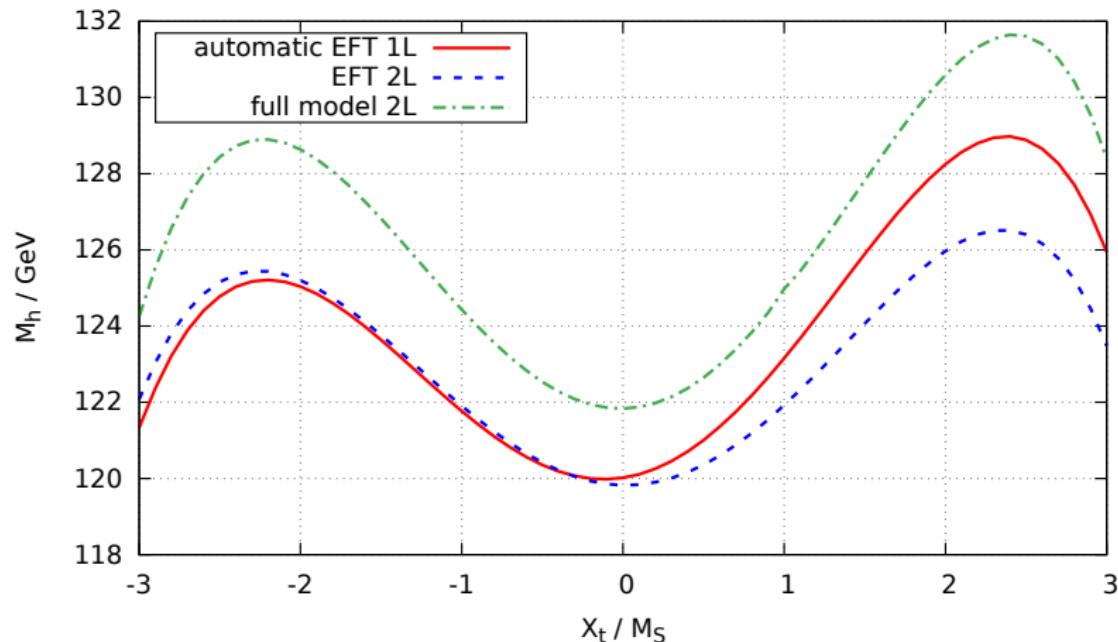
(default algorithm:  $1L$  matching SM  $\rightarrow$  MSSM)

# Comparison full vs. EFT vs. automatic EFT approach



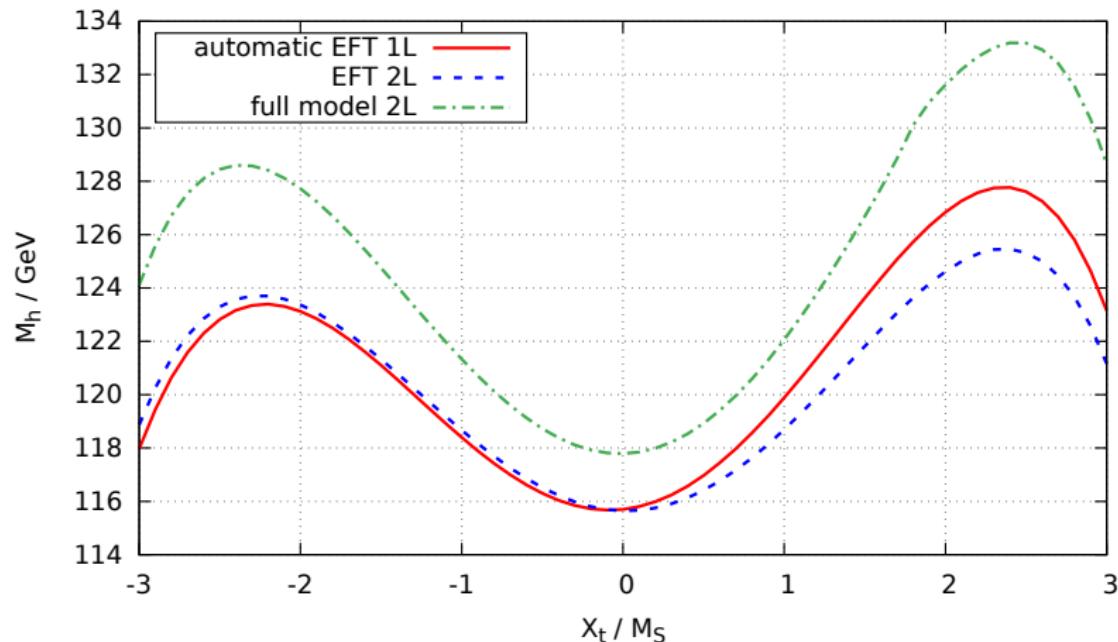
$$\tan \beta = 5, X_{b,\tau} = 0, M_S = 1 \text{ TeV}$$

# Comparison full vs. EFT vs. automatic EFT approach



$$\tan \beta = 5, X_{b,\tau} = 0, M_S = 10 \text{ TeV}$$

# Comparison full vs. EFT vs. automatic EFT approach

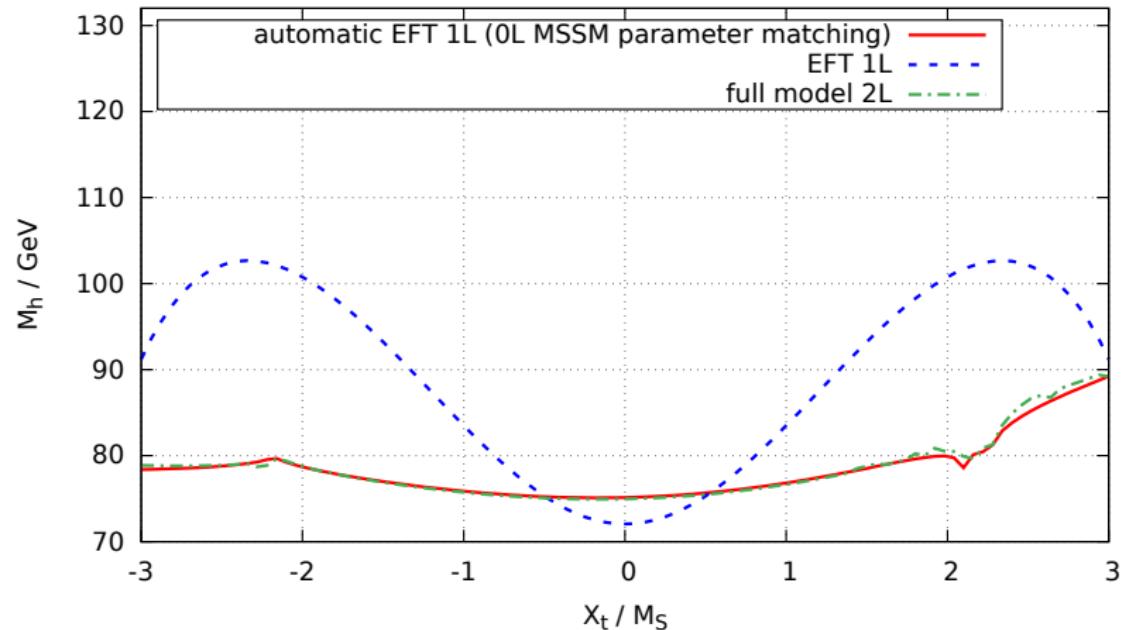


$$\tan \beta = 20, X_{b,\tau} = 0, M_S = 2 \text{ TeV}$$

$X_t$  dependence

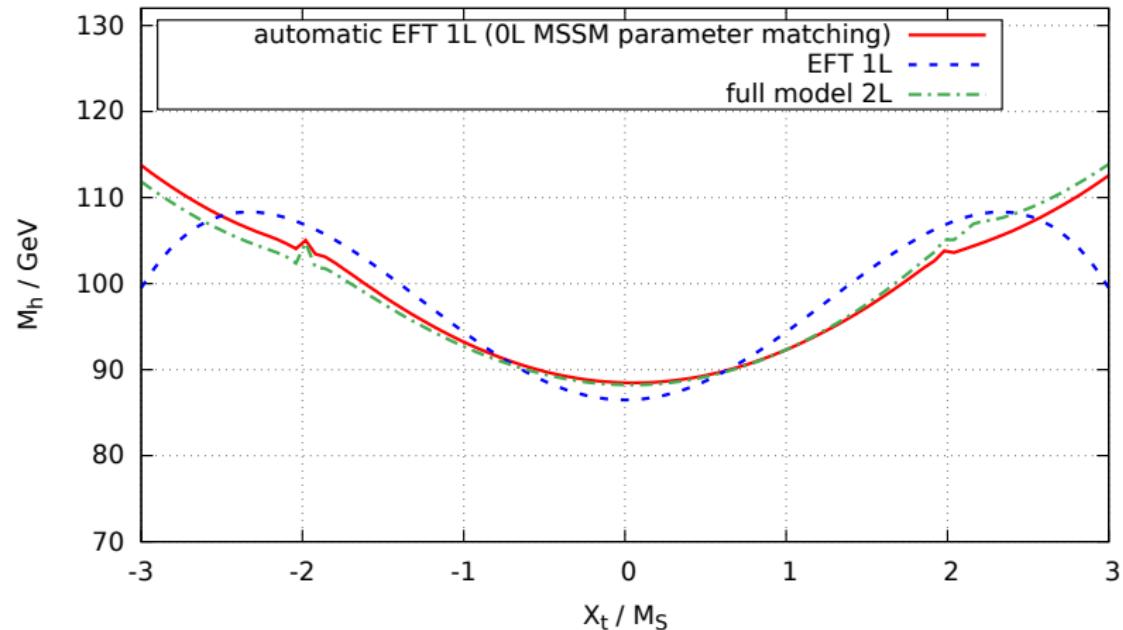
( $0L$  matching SM  $\rightarrow$  MSSM)

# Comparison full vs. EFT vs. automatic EFT approach



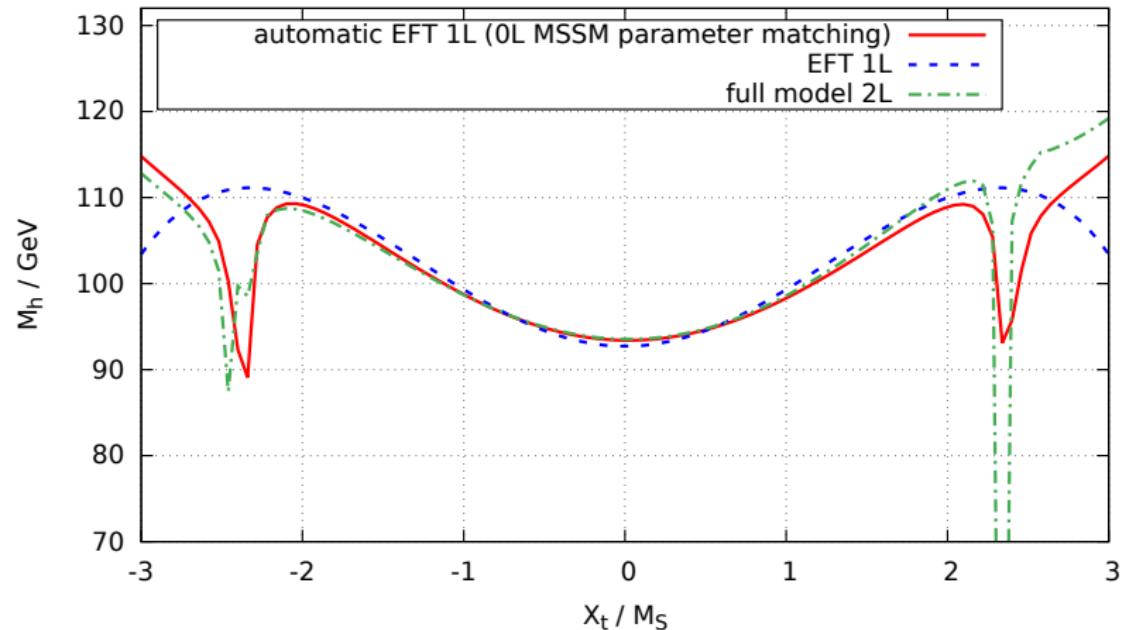
$\tan \beta = 5, X_{b,\tau} = 0, M_S = 100 \text{ GeV}$  (0L matching SM  $\rightarrow$  MSSM)

# Comparison full vs. EFT vs. automatic EFT approach



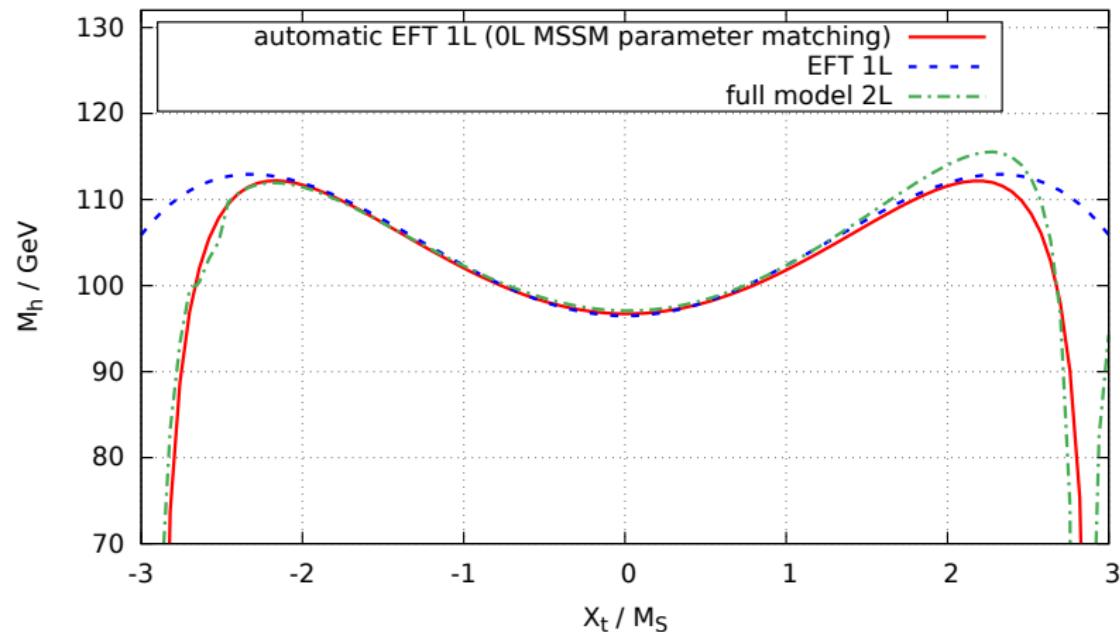
$\tan \beta = 5, X_{b,\tau} = 0, M_S = 200 \text{ GeV}$  (0L matching SM  $\rightarrow$  MSSM)

# Comparison full vs. EFT vs. automatic EFT approach



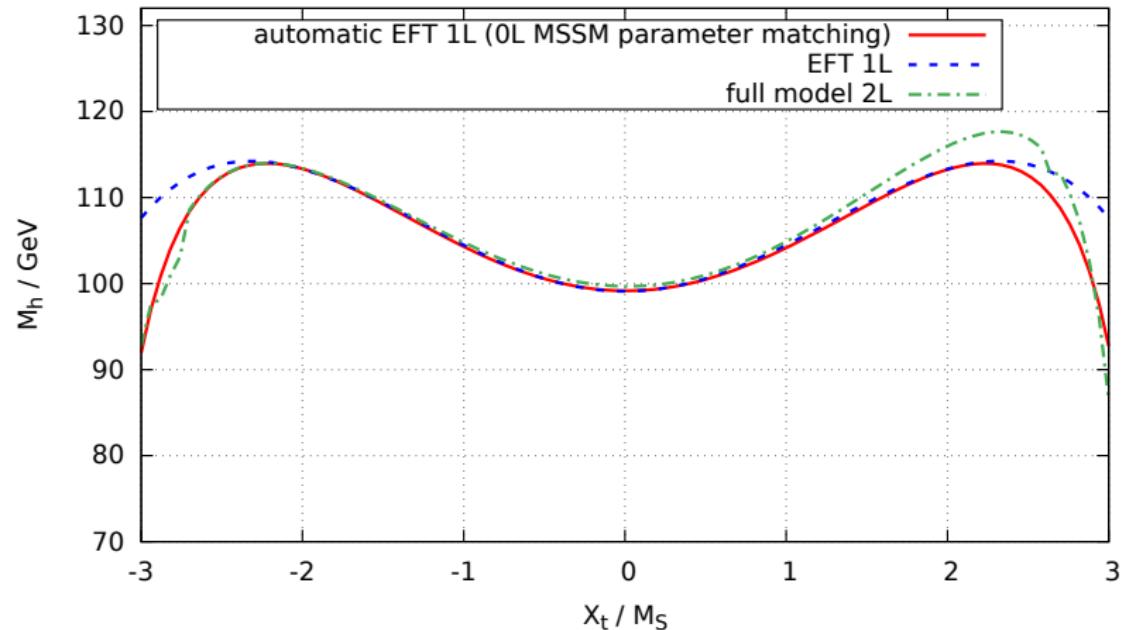
$\tan \beta = 5, X_{b,\tau} = 0, M_S = 300 \text{ GeV}$  (0L matching SM  $\rightarrow$  MSSM)

# Comparison full vs. EFT vs. automatic EFT approach



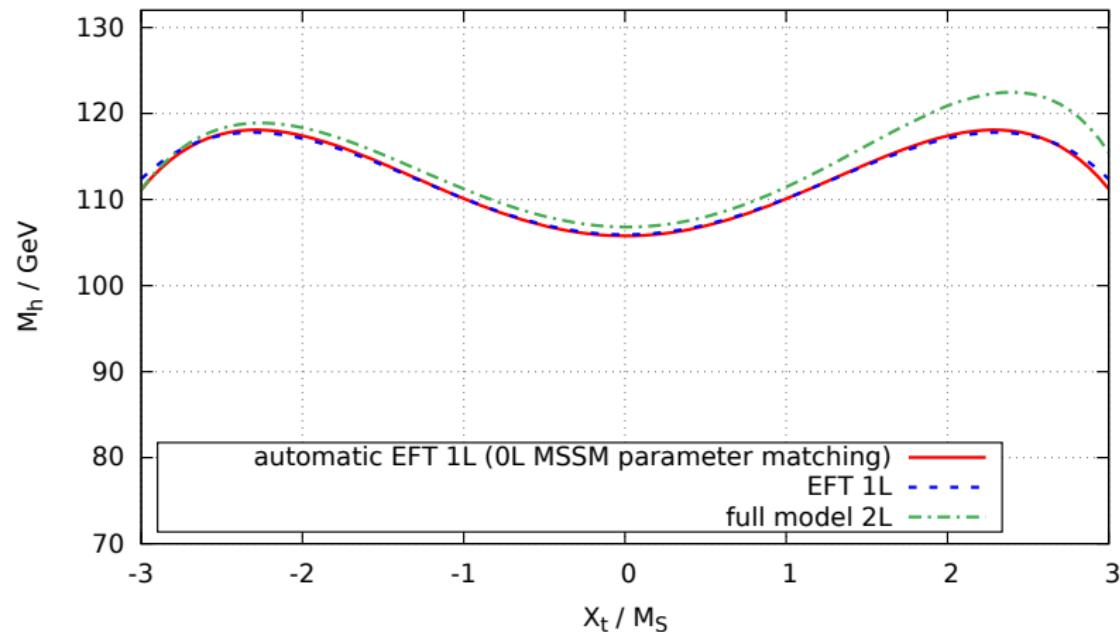
$\tan \beta = 5, X_{b,\tau} = 0, M_S = 400 \text{ GeV}$  (0L matching SM  $\rightarrow$  MSSM)

# Comparison full vs. EFT vs. automatic EFT approach



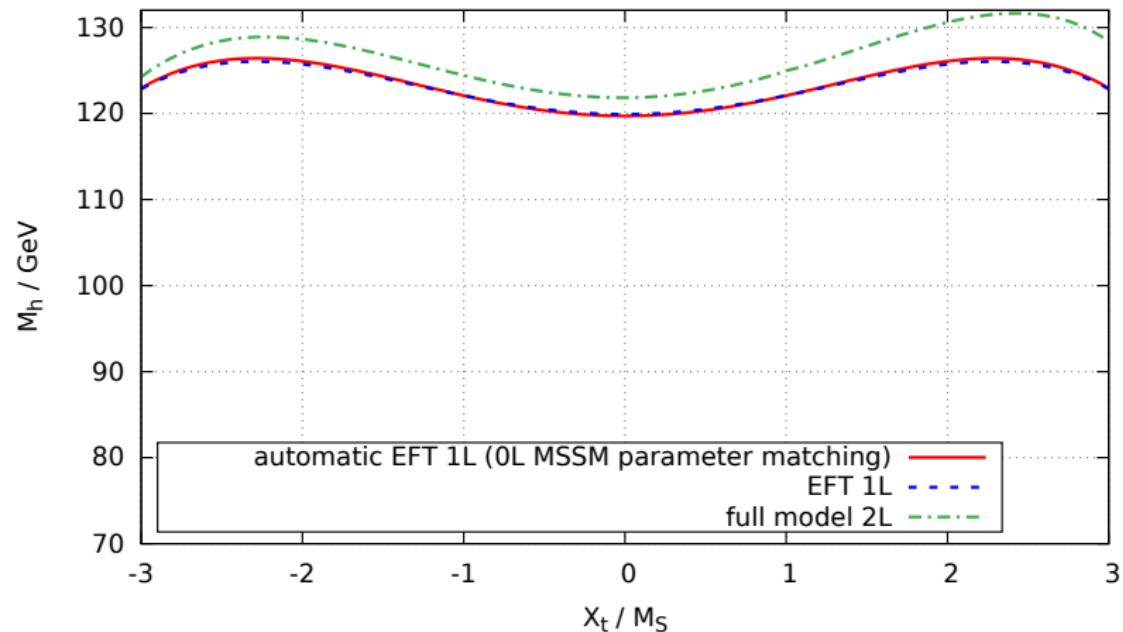
$\tan \beta = 5, X_{b,\tau} = 0, M_S = 500 \text{ GeV}$  (0L matching SM  $\rightarrow$  MSSM)

# Comparison full vs. EFT vs. automatic EFT approach



$\tan \beta = 5, X_{b,\tau} = 0, M_S = 1 \text{ TeV}$  (0L matching SM  $\rightarrow$  MSSM)

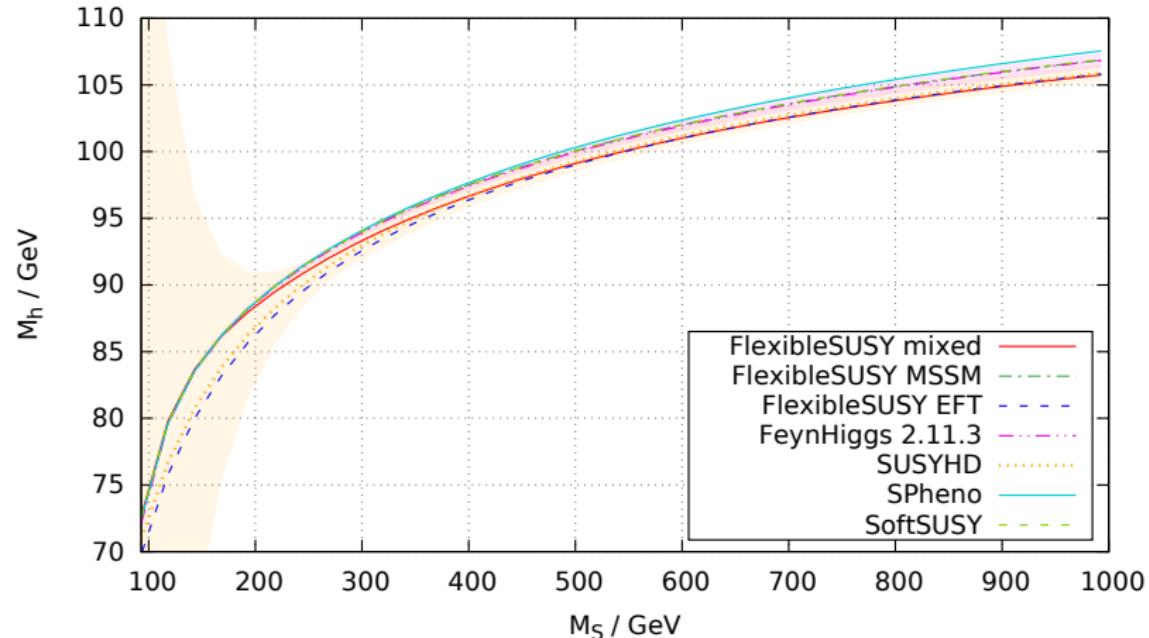
# Comparison full vs. EFT vs. automatic EFT approach



$\tan \beta = 5, X_{b,\tau} = 0, M_S = 10 \text{ TeV}$  (0L matching SM  $\rightarrow$  MSSM)

Comparison with other spectrum generators

# Comparison full vs. EFT vs. automatic EFT approach

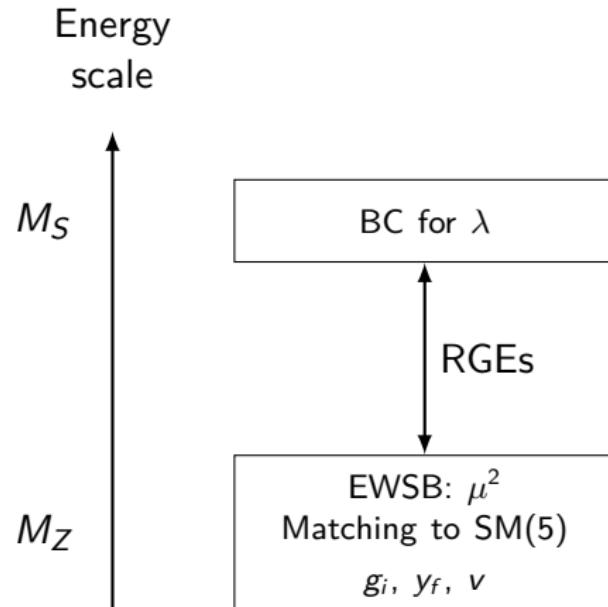


$$\tan \beta = 5, X_{t,b,\tau} = 0$$

# FlexibleSUSY's Weltanschauung

- Model is defined in terms of Lagrangian parameters:  
 $g_i, y_{ij}, v_i, \dots$  in the  $\overline{\text{MS}}/\overline{\text{DR}}$  scheme
- Input parameters:  
 $\alpha_{\text{em,SM}}^{(5),\overline{\text{MS}}}(M_Z), \alpha_{\text{s,SM}}^{(5),\overline{\text{MS}}}(M_Z), M_Z, M_t, G_F, \dots$
- Output parameters:  
 $m_h, M_h, \dots$

# Physical problem statement for the SM



# Algorithm to calculate the model parameters consistent with all BCs

