

SUSY and Higgs mass predictions

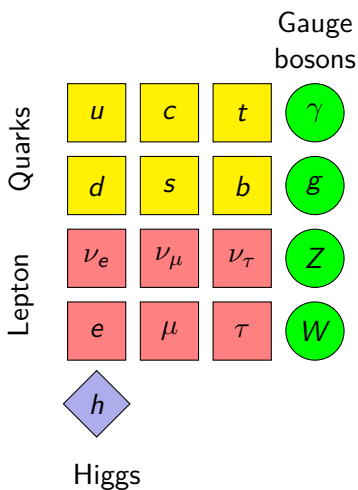
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DESY fellow meeting
DESY Hamburg

02.02.2016



The Standard Model of particle physics



Describes

- Quarks, Leptons, Higgs
- electromagnetic, strong, weak interactions

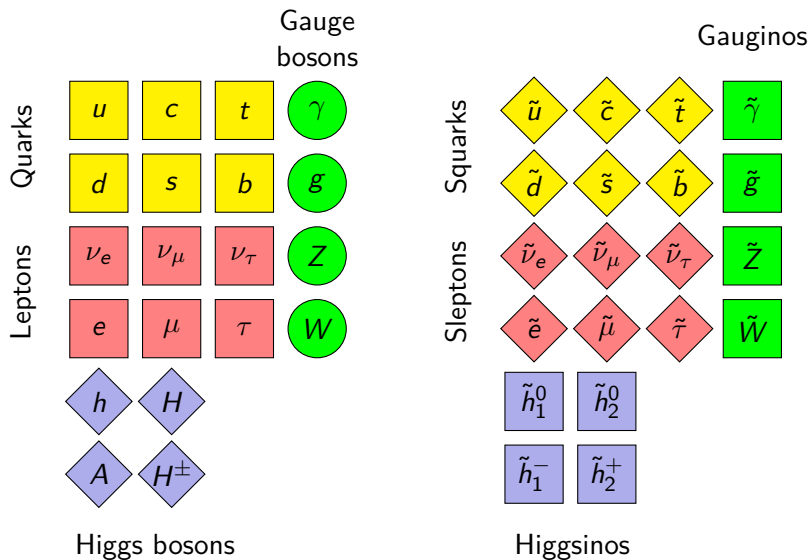
Problems:

- no gravitation
- no dark matter

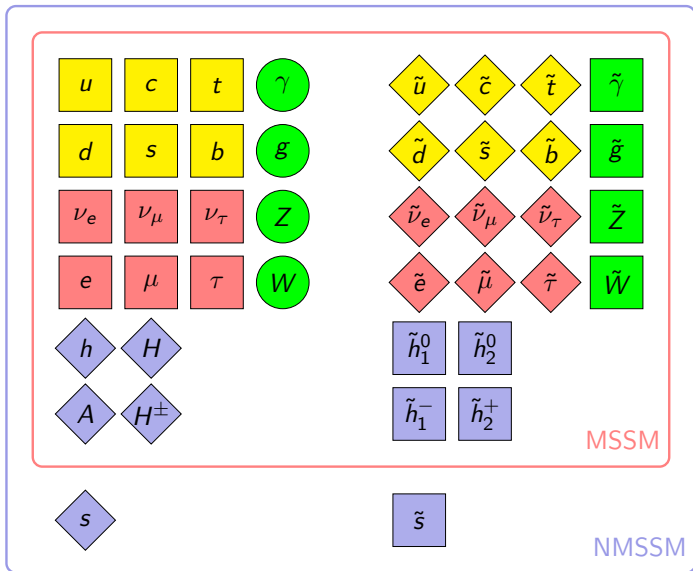
Weaknesses:

- no unification of gauge couplings
- hierarchy problem
- prediction of $(g - 2)_\mu$

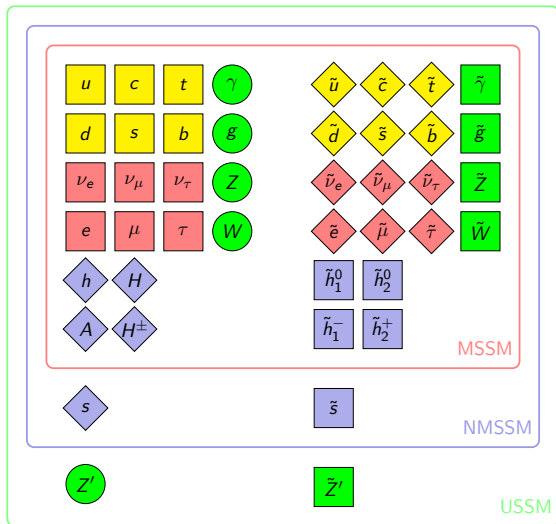
Minimal Supersymmetric Standard Model (MSSM)



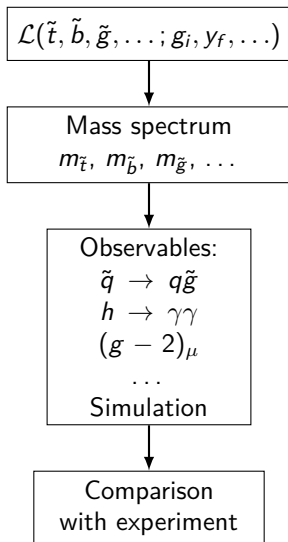
Next-to-MSSM (NMSSM)



$U(1)'$ -extended Supersymmetric Standard Modell (USSM)



Comparison with experiment



Publicly available SUSY spectrum generators

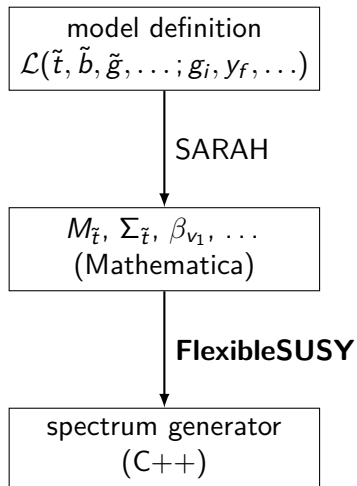
Model	spectrum generator
SM	SMH
THDM	2HDMC
MSSM	ISASUSY, FeynHiggs, SOFTSUSY, SPheno, SuSeFlav, SuSpect, SUSYHD
NMSSM	NMSPEC, SOFTSUSY
USSM	–
CE ₆ SSM	CE6SSMSpecGen
any SUSY-Modell	SARAH, FlexibleSUSY

FlexibleSUSY: a SUSY spectrum generator generator

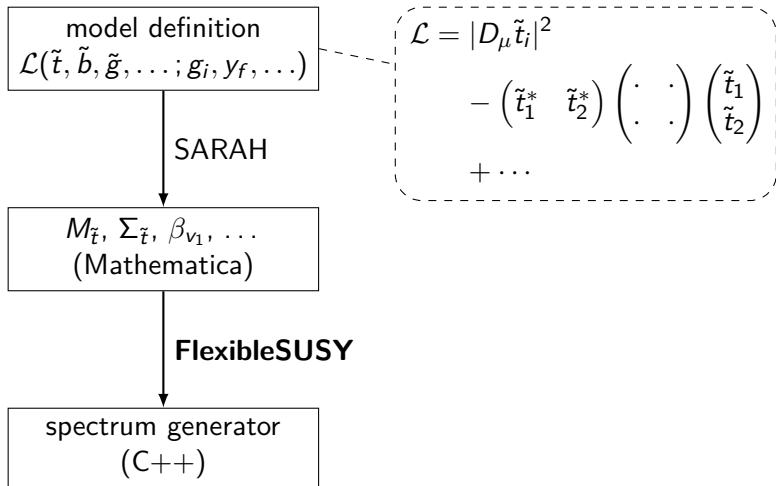
FlexibleSUSY



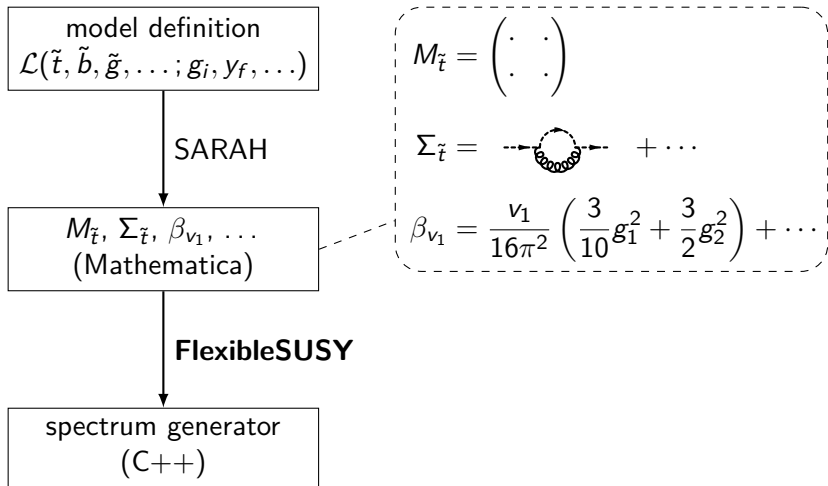
How a spectrum generator is generated



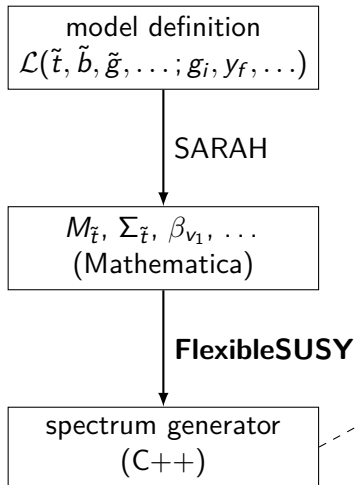
How a spectrum generator is generated



How a spectrum generator is generated



How a spectrum generator is generated



```
Matrix<2,2> get_mass_matrix_St() {  
    Matrix<2,2> mass_matrix;  
  
    mass_matrix(0,0) = ...;  
    mass_matrix(0,1) = ...;  
    mass_matrix(1,0) = ...;  
    mass_matrix(1,1) = ...;  
  
    return mass_matrix;  
}  
  
complex<double> self_energy_St() {  
    complex<double> self_energy;  
  
    self_energy += ...;  
    self_energy += ...;  
    self_energy += ...;  
  
    return self_energy;  
}  
  
double beta_v1() {  
    double beta_v1;  
  
    beta_v1 = v1*(0.3*Sqr(g1)  
        + 1.5*Sqrt(g2))/(16.*Sqr(Pi) + ...;  
  
    return beta_v1;  
}
```

Available models in FlexibleSUSY

Supersymmetric models:

MSSM, NMSSM, SMSSM, USSM, E_6 SSM, TMSSM, MRSSM, NE_6 SSM, $\mu\nu$ SSM, ...

Non-supersymmetric models:

SM, SM + split, THDM, THDM + Higgsinos, THDM + split, SM + singlet, TSM, ...

<https://flexiblesusy.hepforge.org>

Thank you!



Backup

Available models with MSSM high-scale origin

Model	RGEs	h self-energy contributions	matching conditions to the MSSM
MSSM ("full model")	3L	1L + 2L $O((\alpha_t + \alpha_b)\alpha_s)$ + 2L $O((\alpha_t + \alpha_b)^2)$	–
THDM	2L	1L	1L $\lambda_i O((\alpha_t + \alpha_b + \alpha_\tau)\alpha_i)$ + 2L $\lambda_i O(\alpha_t^2\alpha_s)$ [1508.00576]
THDM + \tilde{h}_i	2L	1L	1L $\lambda_i O((\alpha_t + \alpha_b + \alpha_\tau)\alpha_i)$ + 2L $\lambda_i O(\alpha_t^2\alpha_s)$ [1508.00576]
THDM + split	2L	1L	1L $\lambda_i O((\alpha_t + \alpha_b + \alpha_\tau)\alpha_i)$ + 2L $\lambda_i O(\alpha_t^2\alpha_s)$ [1508.00576]
SM + split	2L	1L + 2L $O(\alpha_t(\alpha_s + \alpha_t))$ + 3L gluino $O(\alpha_t\alpha_s^2)$	1L $\tilde{g}_{ij} O(\alpha_t + \alpha_i)$ + 1L $\lambda O((\alpha_t + \alpha_i)^2)$ + 2L $\lambda O(\alpha_s\alpha_t^2)$ [1407.4081]
SM ("EFT")	3L	1L + 2L $O(\alpha_t(\alpha_s + \alpha_t))$	1L $\lambda O((\alpha_t + \alpha_i)^2 + \alpha_b^2 + \alpha_\tau^2)$ + 2L $\lambda O((\alpha_s + \alpha_t)\alpha_t^2)$ [1407.4081, 1504.05200]
SM ("automatic EFT")	3L	1L	1L $\lambda + O(p^2/M_S^2)$ terms

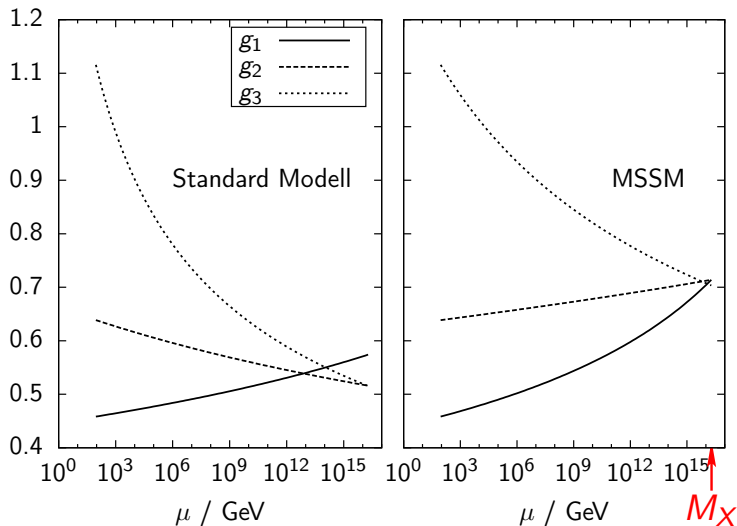
Minimal Supersymmetric Standard Modell (MSSM)

Field	$SU(3)_c \times SU(2)_L \times U(1)_Y$	$SU(5)$
$Q_i = (Q_{u_i} \quad Q_{d_i})$	$(\mathbf{3}, \mathbf{2}, \frac{1}{6})_i$	} $(\mathbf{10})_i$
\bar{U}_i	$(\bar{\mathbf{3}}, \mathbf{1}, -\frac{2}{3})_i$	
\bar{E}_i	$(\mathbf{1}, \mathbf{1}, 1)_i$	
\bar{D}_i	$(\bar{\mathbf{3}}, \mathbf{1}, \frac{1}{3})_i$	} $(\bar{\mathbf{5}})_i$
$L_i = (L_{\nu_i} \quad L_{e_i})$	$(\mathbf{1}, \mathbf{2}, -\frac{1}{2})_i$	
$H_1 = (H_1^0 \quad H_1^-)$	$(\bar{\mathbf{3}}, \mathbf{1}, \frac{1}{3})$	} $(\bar{\mathbf{5}})$
	$(\mathbf{1}, \mathbf{2}, -\frac{1}{2})$	
$H_2 = (H_2^+ \quad H_2^0)$	$(\mathbf{3}, \mathbf{1}, -\frac{1}{3})$	} $(\mathbf{5})$
	$(\mathbf{1}, \mathbf{2}, \frac{1}{2})$	
V_g^a	$(\mathbf{8}, \mathbf{1}, 0)$	$\ni (\mathbf{24})$
V_W^i	$(\mathbf{1}, \mathbf{3}, 0)$	$\ni (\mathbf{24})$
V_Y	$(\mathbf{1}, \mathbf{1}, 0)$	$\ni (\mathbf{24})$

Minimal Supersymmetric Standard Modell (MSSM)

$$\mathcal{W}_{\text{MSSM}} = \mu(H_1 H_2) - y_{ij}^e(H_1 L_i)\bar{E}_j - y_{ij}^d(H_1 Q_i)\bar{D}_j - y_{ij}^u(Q_i H_2)\bar{U}_j$$

Advantage of the MSSM: Gauge Coupling Unification



Weakness of the MSSM: Fine-tuning problem

$$(m_h^{\text{pole}})^2 \approx m_Z^2 \cos^2 2\beta + \Delta m_h^2,$$

$$m_h^{\text{pole}} \approx 125 \text{ GeV}$$

$$m_Z \approx 91 \text{ GeV}$$

$$\Rightarrow \Delta m_h \gtrsim 87 \text{ GeV}$$

\Rightarrow large splitting between m_t and $m_{\tilde{t}_i}$

\Rightarrow large SUSY breaking terms

\Rightarrow spoils SUSY's solution to hierarchy problem

Weakness of the MSSM: μ -Problem

$$\mathcal{W}_{\text{MSSM}} = \mu(H_1 H_2) + \dots$$

On the one hand:

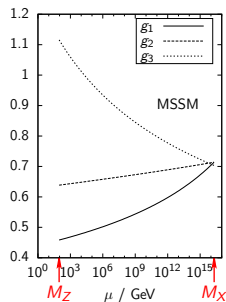
μ has its origin at the GUT scale M_X

$$\Rightarrow \mu \sim M_X \sim 10^{16} \text{ GeV}$$

On the other hand:

μ fixed by EWSB at M_Z

$$\Rightarrow \mu \sim M_Z \sim 10^2 \text{ GeV}$$



Solution: introduce new Higgs singlet S with VEV v_S

\Rightarrow Next-to Minimal Supersymmetric Standard Model
(NMSSM)

Next-to-MSSM (NMSSM)

Field	$SU(3)_c \times SU(2)_L \times U(1)_Y$	$SU(5)$
$Q_i = (Q_{u_i} \quad Q_{d_i})$	$(\mathbf{3}, \mathbf{2}, \frac{1}{6})_i$	} $(\mathbf{10})_i$
\bar{U}_i	$(\bar{\mathbf{3}}, \mathbf{1}, -\frac{2}{3})_i$	
\bar{E}_i	$(\mathbf{1}, \mathbf{1}, 1)_i$	
\bar{D}_i	$(\bar{\mathbf{3}}, \mathbf{1}, \frac{1}{3})_i$	} $(\bar{\mathbf{5}})_i$
$L_i = (L_{\nu_i} \quad L_{e_i})$	$(\mathbf{1}, \mathbf{2}, -\frac{1}{2})_i$	
	$(\bar{\mathbf{3}}, \mathbf{1}, \frac{1}{3})$	} $(\bar{\mathbf{5}})$
$H_1 = (H_1^0 \quad H_1^-)$	$(\mathbf{1}, \mathbf{2}, -\frac{1}{2})$	
	$(\mathbf{3}, \mathbf{1}, -\frac{1}{3})$	} $(\mathbf{5})$
$H_2 = (H_2^+ \quad H_2^0)$	$(\mathbf{1}, \mathbf{2}, \frac{1}{2})$	
S	$(\mathbf{1}, \mathbf{1}, 0)$	$(\mathbf{1})$
V_g^a	$(\mathbf{8}, \mathbf{1}, 0)$	$\ni (\mathbf{24})$
V_W^i	$(\mathbf{1}, \mathbf{3}, 0)$	$\ni (\mathbf{24})$
V_Y	$(\mathbf{1}, \mathbf{1}, 0)$	$\ni (\mathbf{24})$

Next-to-MSSM (NMSSM)

$$\mathcal{W}_{\text{MSSM}+S} = \lambda S(H_1 H_2) + \frac{\kappa}{3} S^3 + \frac{\mu'}{2} S^2 + \xi_F S \\ + \mu(H_1 H_2) - y_{ij}^e(H_1 L_i)\bar{E}_j - y_{ij}^d(H_1 Q_i)\bar{D}_j - y_{ij}^u(Q_i H_2)\bar{U}_j$$

Next-to-MSSM (NMSSM)

$$\mathcal{W}_{\text{MSSM}+S} = \lambda S(H_1 H_2) + \frac{\kappa}{3} S^3 + \frac{\mu'}{2} S^2 + \xi_F S \\ + \mu(H_1 H_2) - y_{ij}^e(H_1 L_i)\bar{E}_j - y_{ij}^d(H_1 Q_i)\bar{D}_j - y_{ij}^u(Q_i H_2)\bar{U}_j$$

Impose Z_3 symmetry to forbid dimensionful couplings and solve the μ -problem

\Rightarrow

$$\mathcal{W}_{\text{NMSSM}} = \lambda S(H_1 H_2) + \frac{\kappa}{3} S^3 \\ - y_{ij}^e(H_1 L_i)\bar{E}_j - y_{ij}^d(H_1 Q_i)\bar{D}_j - y_{ij}^u(Q_i H_2)\bar{U}_j$$

Advantage of the NMSSM: Reduced m_h Fine-tuning

$$(m_h^{\text{pole}})^2 \approx \underbrace{m_Z^2 \cos^2 2\beta}_{\text{MSSM}} + \frac{\lambda^2 v^2}{2} \sin^2 2\beta + \Delta m_h^2$$

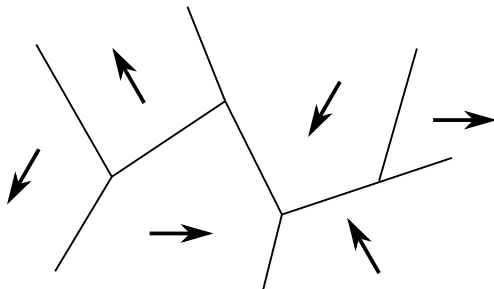
$\underbrace{\hspace{10em}}_{\text{NMSSM}}$

$$\Rightarrow \Delta m_h \gtrsim 55 \text{ GeV}$$

$$\text{MSSM: } \Delta m_h \gtrsim 87 \text{ GeV}$$

Problem of the NMSSM: Domain Walls

Problem: $\mathcal{W}_{\text{NMSSM}}$ has a *discrete* Z_3 symmetry.
 \Rightarrow domain walls



Solution: new *continuous* gauge symmetry $U(1)'$

$\Rightarrow U(1)'$ -extended Supersymmetric Standard Modell (USSM)

$U(1)'$ -extended Supersymmetric Standard Modell (USSM)

Field	$G_{\text{SM}} \times U(1)'$	$SU(5) \times U(1)'$
$Q_i = (Q_{u_i} \quad Q_{d_i})$	$(\mathbf{3}, \mathbf{2}, \frac{1}{6}, N_Q)_i$	} $(\mathbf{10}, N_{10})_i$
\bar{U}_i	$(\bar{\mathbf{3}}, \mathbf{1}, -\frac{2}{3}, N_U)_i$	
\bar{E}_i	$(\mathbf{1}, \mathbf{1}, 1, N_E)_i$	
\bar{D}_i	$(\bar{\mathbf{3}}, \mathbf{1}, \frac{1}{3}, N_D)_i$	} $(\bar{\mathbf{5}}, N_{\bar{\mathbf{5}}})_i$
$L_i = (L_{\nu_i} \quad L_{e_i})$	$(\mathbf{1}, \mathbf{2}, -\frac{1}{2}, N_L)_i$	
	$(\bar{\mathbf{3}}, \mathbf{1}, \frac{1}{3}, N_X)$	} $(\bar{\mathbf{5}}, N_{\bar{\mathbf{5}}})$
$H_1 = (H_1^0 \quad H_1^-)$	$(\mathbf{1}, \mathbf{2}, -\frac{1}{2}, N_{H_1})$	
	$(\mathbf{3}, \mathbf{1}, -\frac{1}{3}, N_X)$	} $(\mathbf{5}, N_{\mathbf{5}})$
$H_2 = (H_2^+ \quad H_2^0)$	$(\mathbf{1}, \mathbf{2}, \frac{1}{2}, N_{H_2})$	
S	$(\mathbf{1}, \mathbf{1}, 0, N_S)$	$(\mathbf{1}, N_1)$
V_g^a	$(\mathbf{8}, \mathbf{1}, 0, 0)$	$\ni (\mathbf{24}, 0)$
V_W^i	$(\mathbf{1}, \mathbf{3}, 0, 0)$	$\ni (\mathbf{24}, 0)$
V_Y	$(\mathbf{1}, \mathbf{1}, 0, 0)$	$\ni (\mathbf{24}, 0)$
V_N	$(\mathbf{1}, \mathbf{1}, 0, 0)$	$\ni (\mathbf{1}, 0)$

$U(1)'$ -extended Supersymmetric Standard Modell (USSM)

If $N_S \neq 0$ and $N_S + N_{H_1} + N_{H_2} = 0$
 \Rightarrow all S^n terms forbidden

$$\mathcal{W}_{\text{USSM}} = \lambda S(H_1 H_2) - y_{ij}^e(H_1 L_i) \bar{E}_j - y_{ij}^d(H_1 Q_i) \bar{D}_j - y_{ij}^u(Q_i H_2) \bar{U}_j$$

Advantage of the USSM: Reduced m_h Fine-tuning

$$(m_h^{\text{pole}})^2 = (m_h^{\text{tree}})^2 + \Delta m_h^2$$

$$(m_h^{\text{tree}})^2 \approx \underbrace{m_Z^2 \cos^2 2\beta}_{\text{MSSM}} + \underbrace{\frac{\lambda^2 v^2}{2} \sin^2 2\beta}_{\text{NMSSM}} + \frac{m_Z^2}{4} \left(1 + \frac{1}{4} \cos 2\beta\right)^2$$

USSM

$$\Rightarrow \Delta m_h \gtrsim 32 \text{ GeV}$$

MSSM: $\Delta m_h \gtrsim 87 \text{ GeV}$

NMSSM: $\Delta m_h \gtrsim 55 \text{ GeV}$

Problem of the USSM: Anomalies

Problem: $U(1)'$ -charges are arbitrary
(as long as $\mathcal{W}_{\text{USSM}}$ is gauge invariant)
 \Rightarrow unsuitable choice can lead to gauge anomalies:

$$\sum_f Z' \text{ (triangle diagram) } = \infty$$

Solution: anomaly-free gauge group, e.g. $SO(10)$ or E_6

\Rightarrow Exceptional Supersymmetric Standard Modell ($E_6\text{SSM}$)

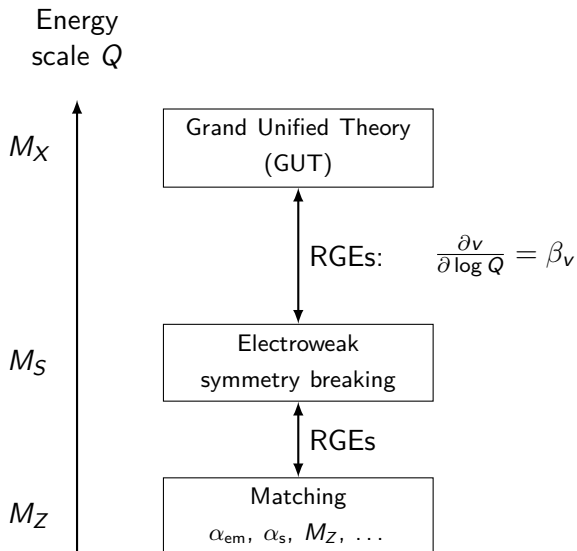
Exceptional Supersymmetric Standard Modell (E_6 SSM)

Field	$G_{SM} \times U(1)_N$	$SU(5) \times U(1)_N$	E_6
$Q_i = (Q_{u_i} \quad Q_{d_i})$	$(\mathbf{3}, \mathbf{2}, \frac{1}{6}, 1)_i$	$(\mathbf{10}, 1)_i$	$(\mathbf{27})_i$
\bar{U}_i	$(\bar{\mathbf{3}}, \mathbf{1}, -\frac{2}{3}, 1)_i$		
\bar{E}_i	$(\mathbf{1}, \mathbf{1}, 1, 1)_i$		
\bar{D}_i	$(\bar{\mathbf{3}}, \mathbf{1}, \frac{1}{3}, 2)_i$	$(\bar{\mathbf{5}}, 2)_i$	
$L_i = (L_{\nu_i} \quad L_{e_i})$	$(\mathbf{1}, \mathbf{2}, -\frac{1}{2}, 2)_i$		
\bar{X}_i	$(\bar{\mathbf{3}}, \mathbf{1}, \frac{1}{3}, -3)_i$	$(\bar{\mathbf{5}}, -3)_i$	
$H_{1i} = (H_{1i}^0 \quad H_{1i}^-)$	$(\mathbf{1}, \mathbf{2}, -\frac{1}{2}, -3)_i$		
X_i	$(\mathbf{3}, \mathbf{1}, -\frac{1}{3}, -2)_i$	$(\mathbf{5}, -2)_i$	
$H_{2i} = (H_{2i}^+ \quad H_{2i}^0)$	$(\mathbf{1}, \mathbf{2}, \frac{1}{2}, -2)_i$		
S_i	$(\mathbf{1}, \mathbf{1}, 0, 5)_i$	$(\mathbf{1}, 5)_i$	
\bar{N}_i	$(\mathbf{1}, \mathbf{1}, 0, 0)_i$	$(\mathbf{1}, 0)_i$	
$H' = (H'^0 \quad H'^-)$	$(\mathbf{1}, \mathbf{2}, -\frac{1}{2}, 2)$	$\ni (\bar{\mathbf{5}}, 2)'$	$\ni (\mathbf{27})'$
$\bar{H}' = (\bar{H}'^+ \quad \bar{H}'^0)$	$(\mathbf{1}, \mathbf{2}, \frac{1}{2}, -2)$	$\ni (\mathbf{5}, -2)'$	$\ni (\bar{\mathbf{27}})'$
V_g^a	$(\mathbf{8}, \mathbf{1}, 0, 0)$	$\ni (\mathbf{24}, 0)$	$\ni (\mathbf{78})$
V_W^i	$(\mathbf{1}, \mathbf{3}, 0, 0)$	$\ni (\mathbf{24}, 0)$	$\ni (\mathbf{78})$
V_Y	$(\mathbf{1}, \mathbf{1}, 0, 0)$	$\ni (\mathbf{24}, 0)$	$\ni (\mathbf{78})$
V_N	$(\mathbf{1}, \mathbf{1}, 0, 0)$	$\ni (\mathbf{1}, 0)$	$\ni (\mathbf{78})$

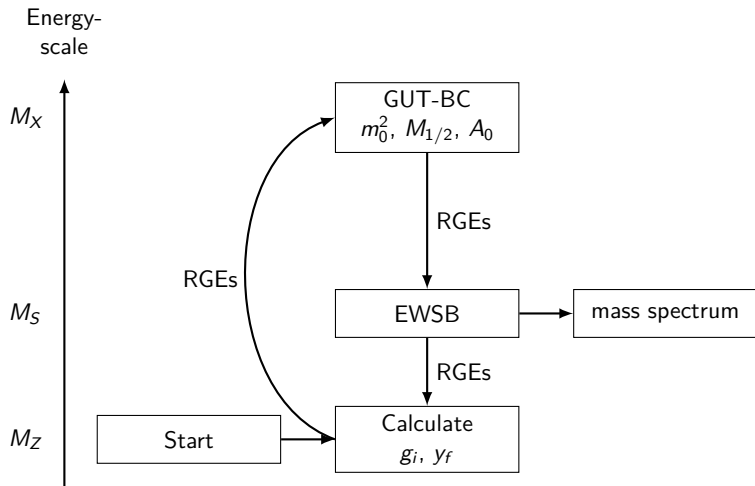
Exceptional Supersymmetric Standard Modell (E₆SSM)

$$\begin{aligned}\mathcal{W}_{E_6SSM} = & \lambda_3 S_3(H_{13}H_{23}) \\ & - y_{ij}^e(H_{13}L_i)\bar{E}_j - y_{ij}^d(H_{13}Q_i)\bar{D}_j - y_{ij}^u(Q_iH_{23})\bar{U}_j \\ & + \kappa_{ij} S_3(X_i\bar{X}_j) + \lambda_{\alpha\beta} S_3(H_{1\alpha}H_{2\beta}) + \mu'(H'\bar{H}')\end{aligned}$$

Physical problem statement

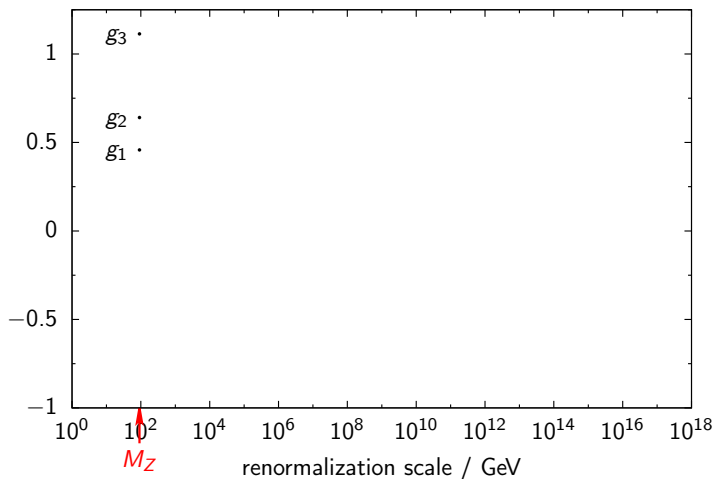


Algorithm to calculate the mass spectrum



Algorithm to calculate the mass spectrum

Iteration step 1: calculate gauge couplings $g_i^{\overline{\text{DR}}}(M_Z)$



Calculate gauge coupling $g_3^{\overline{\text{DR}}}(M_Z)$

Input: $\alpha_{s,\text{SM}}^{(5),\overline{\text{MS}}}(M_Z) = 0.1185$

→

$$\alpha_s^{\overline{\text{DR}}}(M_Z) = \frac{\alpha_{s,\text{SM}}^{(5),\overline{\text{MS}}}(M_Z)}{1 - \Delta\alpha_{s,\text{SM}}(M_Z) - \Delta\alpha_s(M_Z)}$$

with

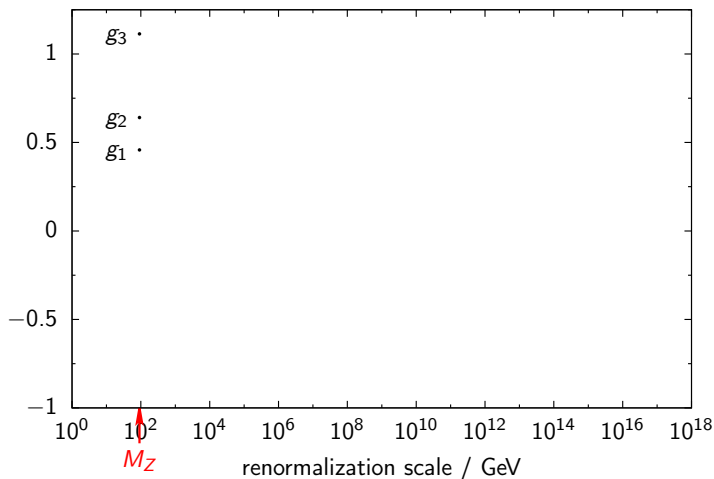
$$\Delta\alpha_{s,\text{SM}}(\mu) = \frac{\alpha_s}{2\pi} \left[-\frac{2}{3} \log \frac{m_t}{\mu} \right]$$
$$\Delta\alpha_s(\mu) = \frac{\alpha_s}{2\pi} \left[\frac{1}{2} - \sum_{\text{SUSY particle } f} T_f \log \frac{m_f}{\mu} \right]$$

⇒

$$g_3^{\overline{\text{DR}}}(M_Z) = \sqrt{4\pi\alpha_s^{\overline{\text{DR}}}(M_Z)}$$

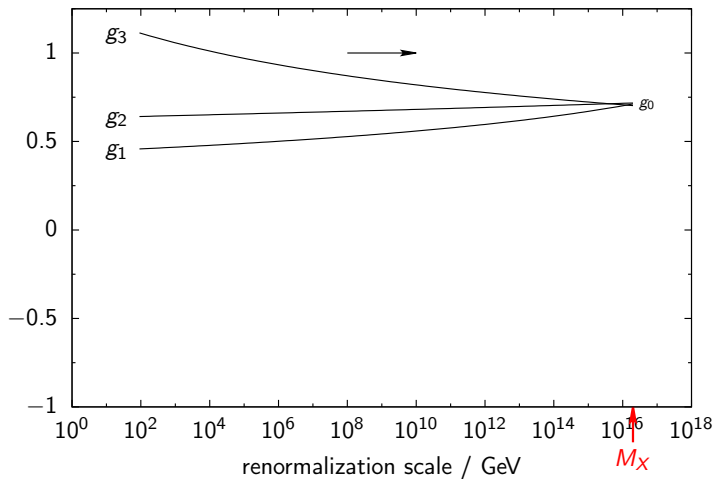
Algorithm to calculate the mass spectrum

Iteration step 1: calculate gauge couplings $g_i^{\overline{\text{DR}}}(M_Z)$



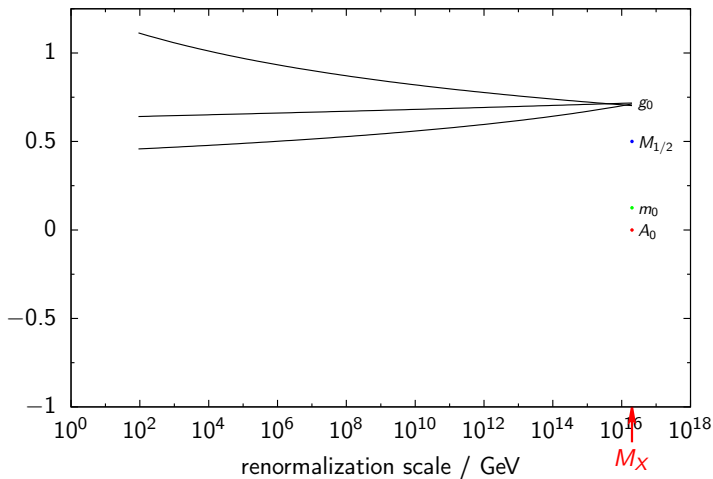
Algorithm to calculate the mass spectrum

Iteration step 1: RG running of gauge couplings to M_X



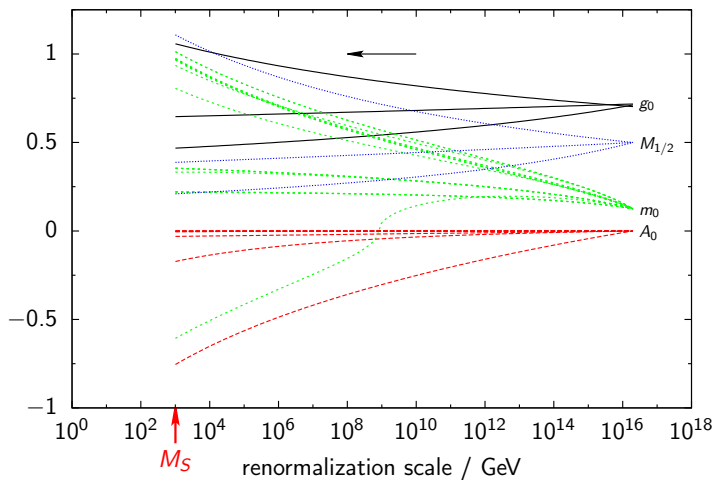
Algorithm to calculate the mass spectrum

Iteration step 1: impose boundary conditions at M_X



Algorithm to calculate the mass spectrum

Iteration step 1: RG running to M_S , impose EWSB conditions



Solve EWSB conditions

Solve EWSB conditions by fixing N model parameters (p_1, \dots, p_N) :

$$0 = \frac{\partial V^{\text{tree}}}{\partial v_i} - t_i(m_f) \quad (i = 1, \dots, N)$$

For example: CMSSM solve for $(\mu, B\mu)$

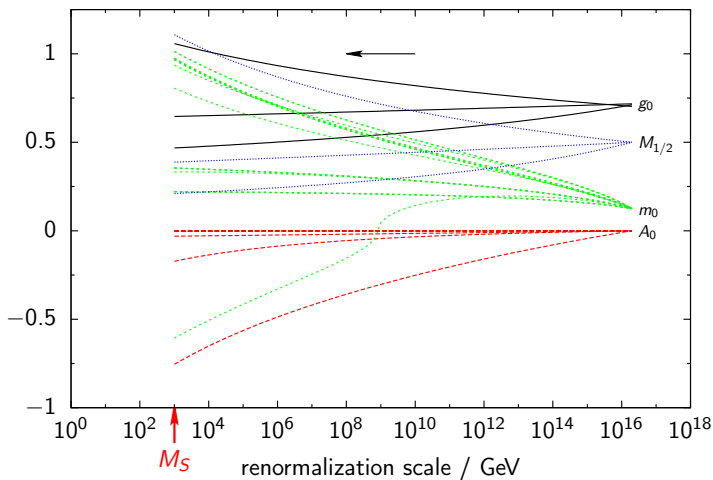
$$0 = m_{h_1}^2 v_1 + |\mu|^2 v_1 - B\mu v_2 + \frac{\bar{g}^2}{8} (v_1^2 - v_2^2) v_1 - t_1(m_f)$$

$$0 = m_{h_2}^2 v_2 + |\mu|^2 v_2 - B\mu v_1 - \frac{\bar{g}^2}{8} (v_1^2 - v_2^2) v_2 - t_2(m_f)$$

where $\bar{g}^2 = g_Y^2 + g_2^2$

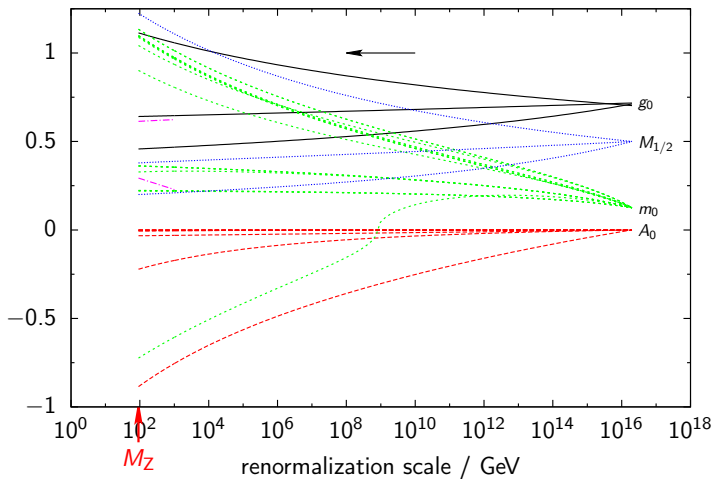
Algorithm to calculate the mass spectrum

Iteration step 1: RG running to M_S , impose EWSB on $(\mu, B\mu)$



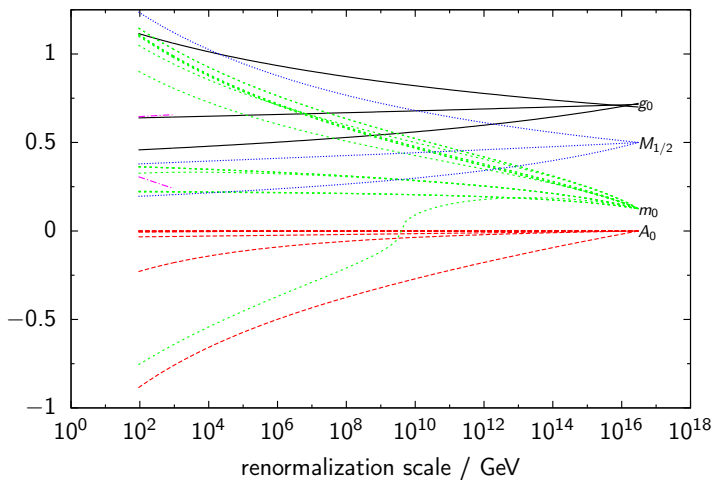
Algorithm to calculate the mass spectrum

Iteration step 1: RG running to M_Z



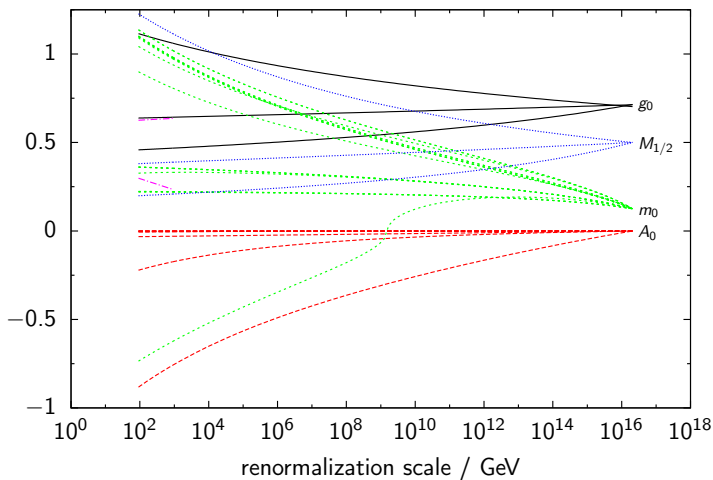
Algorithm to calculate the mass spectrum

Iteration step 2:



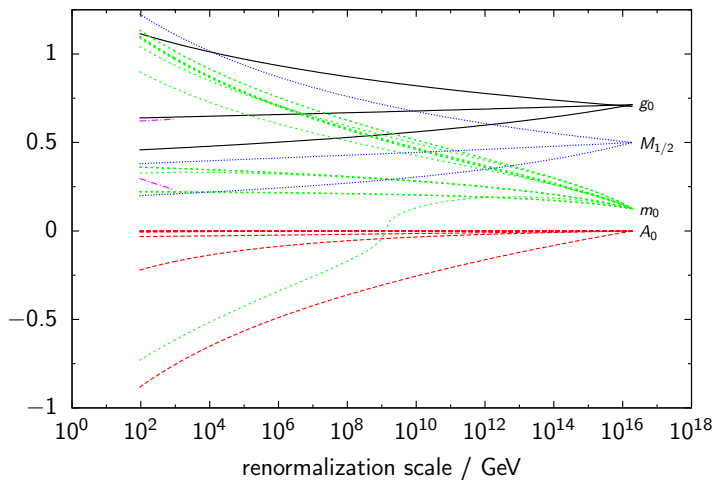
Algorithm to calculate the mass spectrum

Iteration step 3:

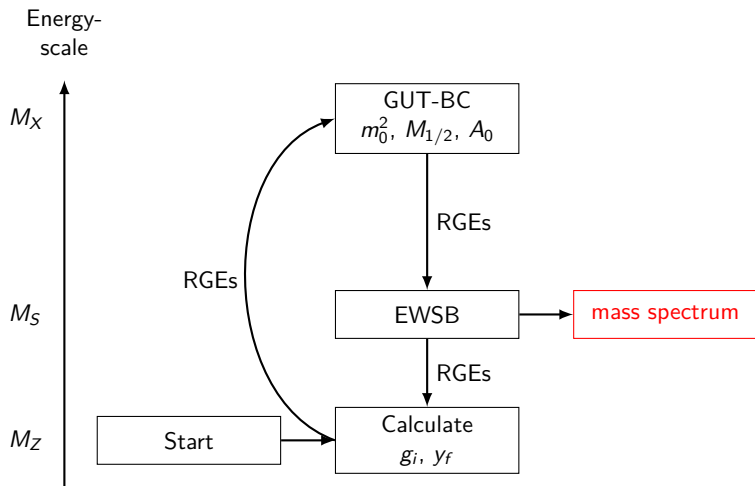


Algorithm to calculate the mass spectrum

Iteration step 8: convergence



Algorithm to calculate the mass spectrum



Calculate the pole mass spectrum

Solve for each particle f :

$$0 = \det \left[p^2 \mathbf{1} - M_f + \hat{\Sigma}_f(p^2) \right]$$

For example $f = \text{Higgs}$:

$$M_h = \begin{pmatrix} (M_h)_{11} & (M_h)_{12} \\ (M_h)_{12} & (M_h)_{22} \end{pmatrix}$$

$$(M_h)_{11} = m_{h_1}^2 + |\mu|^2 + \frac{1}{8}(g_Y^2 + g_2^2)(3v_1^2 - v_2^2)$$

$$(M_h)_{12} = -\frac{1}{2}(B\mu + B\mu^*) - \frac{1}{4}v_2 v_1 (g_Y^2 + g_2^2)$$

$$(M_h)_{22} = m_{h_2}^2 + |\mu|^2 + \frac{1}{8}(g_Y^2 + g_2^2)(3v_2^2 - v_1^2)$$

$$\hat{\Sigma}_h(p^2) = \Sigma_h(p^2) - \delta M_h^2 + (p^2 \mathbf{1} - M_h^2) \delta Z_h,$$

$$\delta M_h^2 = \Sigma_h(p^2) \Big|_{\Delta}, \quad \delta Z_h = -\Sigma'_h(p^2) \Big|_{\Delta}$$

Parameters scans in the CMSSM, NMSSM, USSM, E₆SSM

Models: CMSSM, NMSSM, USSM, E₆SSM

mSUGRA-inspired GUT scale BCs:

$$(m_{\tilde{f}}^2)_{ij}(M_X) = m_0^2 \delta_{ij}$$

$$A_{ij}^f(M_X) = A_0$$

$$M_i(M_X) = M_{1/2}$$

parameters fixed by EWSB conditions at M_S :

CMSSM: $\mu, B\mu$

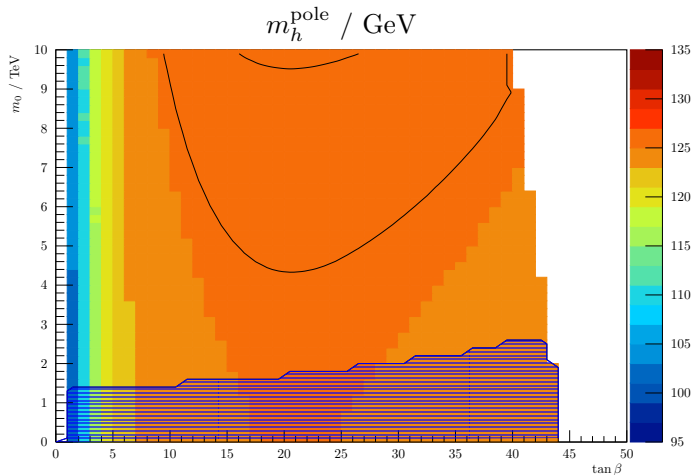
NMSSM: κ, v_S, m_S^2

USSM: $m_{h_1}^2, m_{h_2}^2, m_S^2$

E₆SSM: $m_{h_1}^2, m_{h_2}^2, m_S^2$

2-dimensional scan over: $m_0, \tan \beta = v_2/v_1$

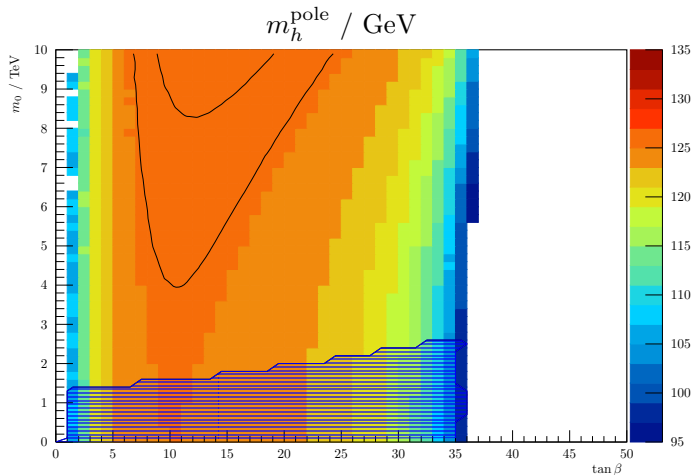
CMSSM parameter scan



$M_{1/2} = A_0 = 5 \text{ TeV}$, $\text{sign } \mu = +1$

Higgs mass contours at $m_h^{\text{pole}} = (125.7 \pm 0.4) \text{ GeV}$

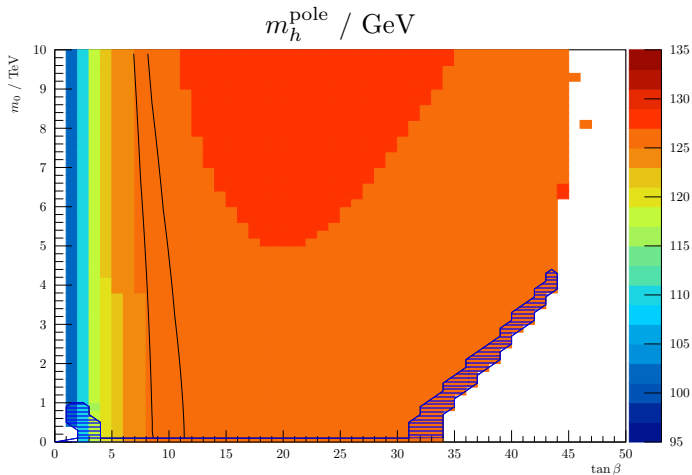
NMSSM parameter scan



$M_{1/2} = -A_0 = 5 \text{ TeV}$, $\lambda(M_X) = 0.1$, $\text{sign } v_s = +1$

Higgs mass contours at $m_h^{\text{pole}} = (125.7 \pm 0.4) \text{ GeV}$

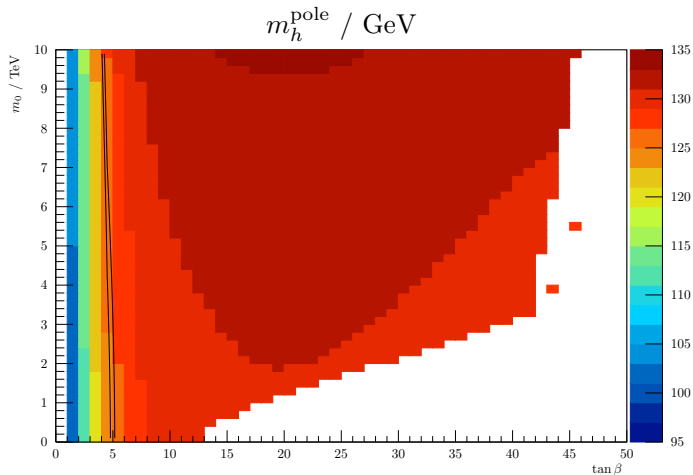
USSM parameter scan



$M_{1/2} = A_0 = 5 \text{ TeV}, \lambda(M_X) = 0.1, v_s = 10 \text{ TeV}$

Higgs mass contours at $m_h^{\text{pole}} = (125.7 \pm 0.4) \text{ GeV}$

E_6 SSM parameter scan



$M_{1/2} = A_0 = 5 \text{ TeV}, \lambda(M_X) = \kappa(M_X) = 0.1, v_s = 10 \text{ TeV}$
Higgs mass contours at $m_h^{\text{pole}} = (125.7 \pm 0.4) \text{ GeV}$