

The space  $V$  is complete: (assuming appropriate sequences / series converge).  
 A complete inner product space like this is a Hilbert space - often  $V$  is called that in QM  
 ( $V$  can be finite or  $\infty$ -dimensional)  
 An operator  $Q$  is a linear map on states ( $V \rightarrow V$ )  $| \psi \rangle \mapsto Q | \psi \rangle$   
 ie  $Q(\alpha | \psi \rangle + \beta | \phi \rangle) = \alpha Q | \psi \rangle + \beta Q | \phi \rangle$   
 Some op. can be regarded as acting 'to the left' and dual states ( $V^* \rightarrow V^*$ ).

$\langle \phi | \mapsto \langle \phi | Q$ ,  
 where the RHS is defined by knowing  $\langle \phi | Q | \psi \rangle = \langle \phi | (Q | \psi \rangle)$  or just write  $\langle \phi | Q | \psi \rangle$ .

$\forall Q$  hermitian conjugate or adjoint is an op  $Q^\dagger$  defined by  $\langle \phi | Q^\dagger = (Q | \phi \rangle)^\dagger$  or  
 equivalently  $\langle \psi | Q^\dagger | \phi \rangle = (Q | \psi \rangle)^\dagger | \phi \rangle = \langle \psi | Q | \phi \rangle^*$   
 $\Rightarrow$  for any ops  $A, B$ :  $(\alpha A + \beta B)^\dagger = \alpha^* A^\dagger + \beta^* B^\dagger$ ,  $(AB)^\dagger = B^\dagger A^\dagger$ .

For any op.  $Q$ ,  $| \psi \rangle$  is an eigenstate with e' value  $\lambda$  if  $Q | \psi \rangle = \lambda | \psi \rangle$  (take  $\psi \Rightarrow \langle \psi | Q = \lambda^* \langle \psi |$ )  
 For general  $Q$ ,  $\lambda$  can be complex. Commutator take on explicit particular importance:  
 $[A, B] = AB - BA = - [B, A]$ .

1b Observables + Measurements

An op.  $Q$  is hermitian or self-adjoint if  $Q^\dagger = Q$ : called "observables" - correspond to physical, measurable quantities of position, momentum, energy, angular momentum.

- Key results
- (i) All eigenvalues are  $\mathbb{R}$
  - (ii) E' states with distinct e' vals are orthogonal
  - (iii) Any state can be expanded in terms of (ie written as a linear combination of) e' states.

Prove (i), (ii) + assume (iii):

(i)  $Q | \psi \rangle = \lambda | \psi \rangle \Rightarrow \langle \psi | Q^\dagger = \lambda^* \langle \psi | \Rightarrow \langle \psi | Q | \psi \rangle = \lambda \langle \psi | \psi \rangle = \lambda^* \langle \psi | \psi \rangle$   
 $= \langle \psi | Q | \psi \rangle$  if  $Q$  hermitian. and  $\langle \psi | \psi \rangle \neq 0$  so  $\lambda = \lambda^*$ .

(ii) Let  $| n \rangle$  be e' states of  $Q$ , e' vals  $\lambda = q_n$  real ( $n$  discrete label, pos.  $\infty$ )  
 $Q | n \rangle = q_n | n \rangle$  and  $Q | m \rangle = q_m | m \rangle$  or  $\langle m | Q = q_m \langle m |$   
 $\Rightarrow \langle m | Q | n \rangle = q_n \langle m | n \rangle = q_m \langle m | n \rangle$ . so  $q_n \neq q_m \Rightarrow \langle m | n \rangle = 0$ .

Combining these 3 properties we have:

\* For any observable  $Q \exists$  an orthonormal basis of e' states  $\{ | n \rangle \}$  for the space of states  $V$   
 with  $Q | n \rangle = q_n | n \rangle$ ,  $\langle m | n \rangle = \delta_{mn}$   
 This means that a general state can be expanded  $| \psi \rangle = \sum_n \alpha_n | n \rangle$  where  $\alpha_n = \langle n | \psi \rangle$

For the state to be properly normalised  $\| | \psi \rangle \|^2 = \langle \psi | \psi \rangle = 1 \Leftrightarrow (\sum_n \alpha_n^* \langle m |) (\sum_n \alpha_n | n \rangle) = \sum_n |\alpha_n|^2 = 1$ .

There may be several states with the same e' val.  $\lambda$ . Define the e'space for a given e' val  
 $V_\lambda = \{ | \psi \rangle : Q | \psi \rangle = \lambda | \psi \rangle \}$  which has basis  $\{ | n \rangle : q_n = \lambda \}$

The "degeneracy" of  $\lambda$  is the # of states in this basis, or  $\dim V_\lambda$ .  
 $\lambda$  is "non-degenerate" if  $\dim V_\lambda = 1$ .

[Passing from our three key results to the conclusion (\*) is achieved by choosing an orthonormal basis for each  $V_\lambda$ . (ii) ensures that these spaces are mutually orthogonal, (iii)  $\Rightarrow$  sum of all eigenspaces is  $V$ , the entire space of states.]

A projection operator  $P_\lambda$  onto any e'space ( $V + V_\lambda$ ) can be defined by its action on basis states  
 $P_\lambda | n \rangle = \begin{cases} | n \rangle & \text{if } q_n = \lambda \\ 0 & \text{otherwise} \end{cases}$  ;  $\sum_\lambda P_\lambda = 1$  ;  $\lambda \neq \mu \Rightarrow P_\lambda P_\mu = P_\mu P_\lambda = 0$ .

Consider a measurement of  $Q$  when the system is in a state  $| \psi \rangle$  immediately before.  
 Then,  
 • The result is an e' value  $\lambda$   
 • This is obtained with probability  $p(\lambda) = \| P_\lambda | \psi \rangle \|^2 = \sum_{n: q_n = \lambda} |\alpha_n|^2$ .